

Supersymmetric Dark Matter candidates in light of constraints from collider and astroparticle observables

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Outline

1 Motivations

2 Neutralino DM in the (N)MSSM

- C. Boehm, JDS, A. Mazumdar and E. Pukartas, Phys. Rev. D87 (2013) 023529, arXiv :1205.2815
- G. Bélanger, C. Boehm, M. Cirelli, JDS and A. Pukhov, JCAP 1211 (2012) 028, arXiv :1208.5009
- D. A. Vasquez, G. Bélanger, C. Boehm, JDS, P. Richardson and C. Wymant, Phys. Rev. D86 (2012) 035023, arXiv :1203.3446

3 U(1) extensions of the MSSM

- G. Bélanger, JDS and A. Pukhov, JCAP 1112 (2011) 014, arXiv :1110.2414
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4 Conclusions

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3 U(1) extensions of the MSSM

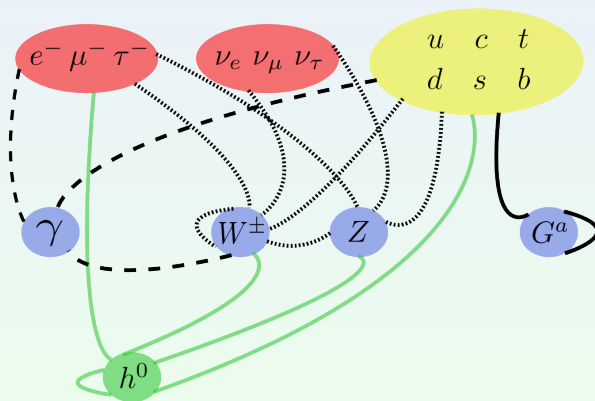
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4 Conclusions

Success of the Standard Models

* Particle Physics (SM)

- * All particles discovered since 1 year - 1 day

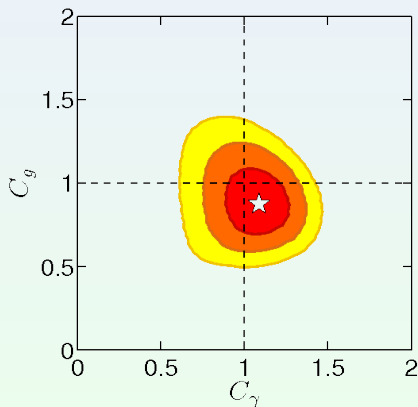
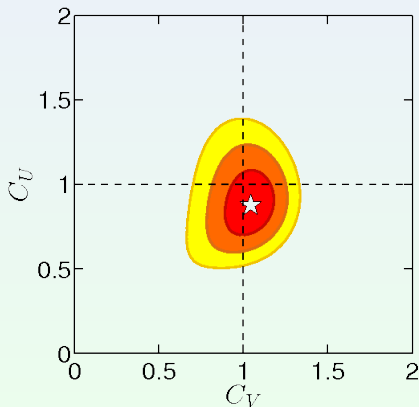


SM interactions, at tree-level

Success of the Standard Models

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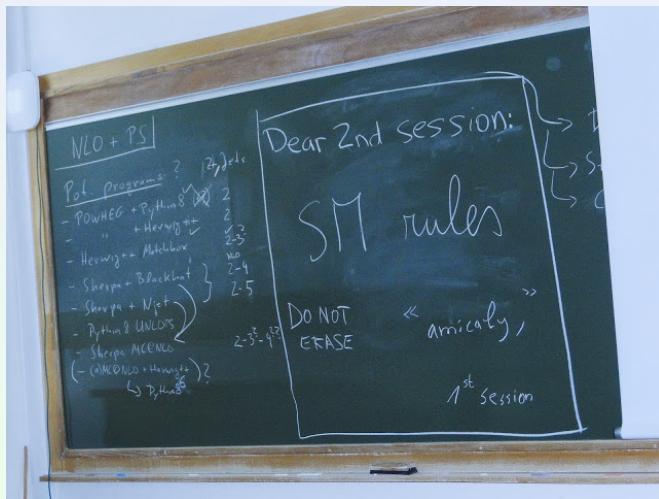
- * All particles discovered since 1 year - 1 day
- * With the expected properties, ...



Reduced Higgs couplings, G. Bélanger, B. Dumont, U. Ellwanger, J. F. Gunion and S. Kraml, arXiv :1306.2941

Success of the Standard Models

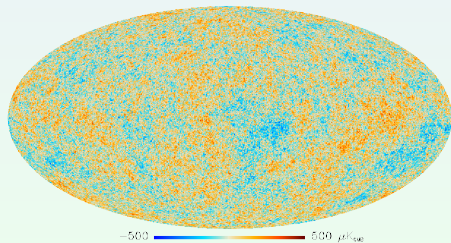
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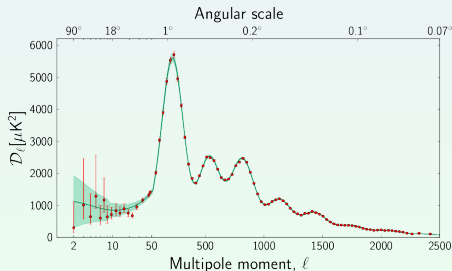
Les Houches Workshop 2013, courtesy of F. Boudjema

Success of the Standard Models

- * Particle Physics (SM)
- * Cosmology (Λ CDM)
 - * Simple cosmological model which fits even the most accurate measurements (Planck satellite)



Temperature fluctuations of the CMB



Temperature angular power spectrum

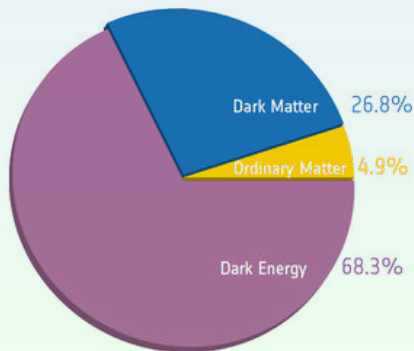
P. A. R. Ade et al., arXiv :1303.5062

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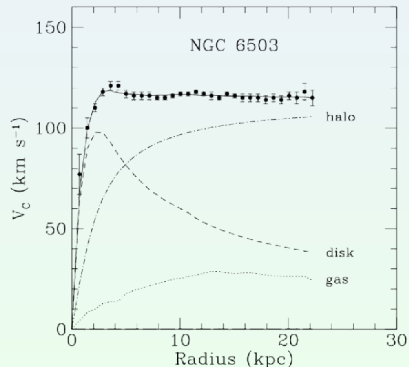
* Particle Physics (SM)

* Cosmology (Λ CDM)

- * Needs Dark Energy and Dark Matter (DM, other evidence : rotation curves of galaxies, bullet cluster, ...)



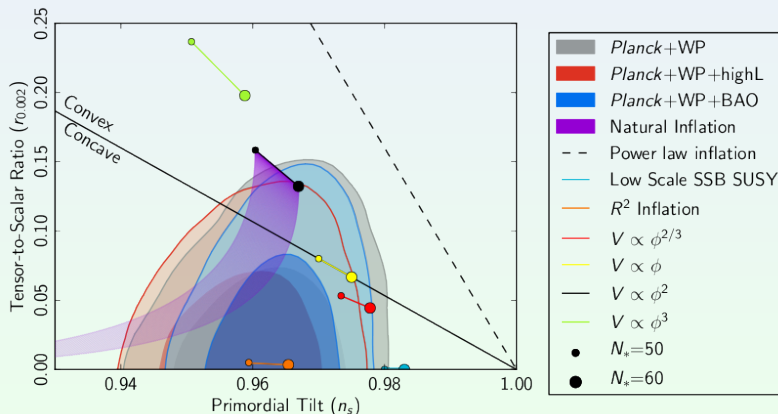
Λ CDM after Planck measurements



K. G. Begeman, A. H. Broeils and R. H. Sanders, 1991, MNRAS, 249, 523

Success of the Standard Models

- * Particle Physics (SM)
- * Cosmology (Λ CDM)
 - * Simple models of inflation are still valid

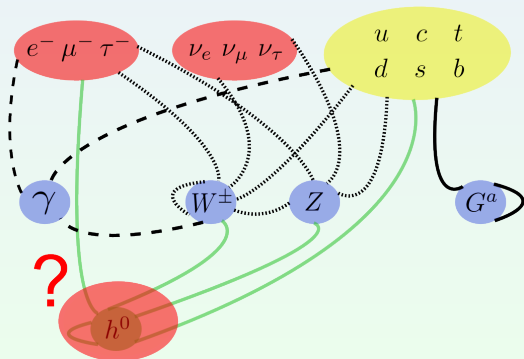


Constraints on inflationary models from Planck satellite

Drawbacks of the Standard Models

* Particle Physics (SM)

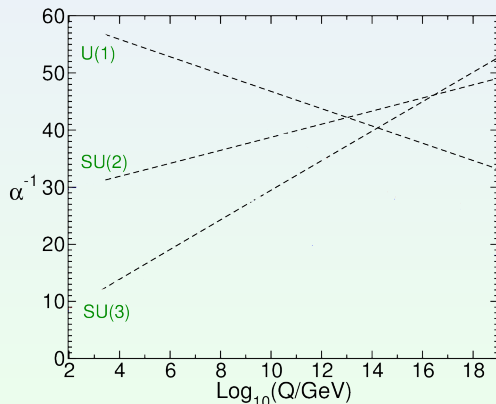
- * Hierarchy problem between EW (~ 100 GeV) and Planck ($\sim 10^{19}$ GeV) scales
Quadratic divergences to the Higgs boson mass squared



Drawbacks of the Standard Models

* Particle Physics (SM)

- * Hierarchy problem between EW (~ 100 GeV) and Planck ($\sim 10^{19}$ GeV) scales
Quadratic divergences to the Higgs boson mass squared
- * Grand Unification (GUT)

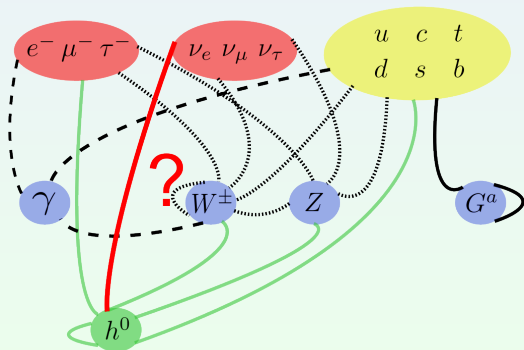


Evolution of SM gauge couplings, Stephen P. Martin, arXiv :hep-ph/9709356

Drawbacks of the Standard Models

* Particle Physics (SM)

- * Hierarchy problem between EW (~ 100 GeV) and Planck ($\sim 10^{19}$ GeV) scales
Quadratic divergences to the Higgs boson mass squared
- * Grand Unification
- * Neutrino sector (Dirac, Majorana ??), ...



Drawbacks of the Standard Models

- * Particle Physics (SM)

- * Cosmology (Λ CDM)

- * DM made of particles \neq SM particles :

- ✗ baryons : BBN, CMB, ...

- ✗ **charged leptons** : we would have seen DM (overproduction of γ , ...)

- ✗ **neutrinos** : too light \Rightarrow low relic density + HDM

\Rightarrow Example of DM candidate which gives the right abundance :
Weakly Interacting Massive Particle (WIMP)

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\Rightarrow Example of DM candidate which gives the right abundance :
Weakly Interacting Massive Particle (WIMP)

✓ Candidates can be found beyond the Standard Model
Here : Supersymmetry (SUSY)

Supersymmetry

* Fermions \Leftrightarrow bosons \Rightarrow solution to the Hierarchy problem

New particles : cancellation of the quadratic term in

$$\Delta m_{h^0}^2 \Big|_{\text{SM}} = \frac{y_f^2}{16\pi^2} \left(-2\Lambda^2 + 6m_f^2 \ln \frac{\Lambda}{m_f} + \dots \right) \text{ with that of}$$

$$\Delta m_{h^0}^2 \Big|_{\text{SUSY}} = \frac{\lambda_s}{16\pi^2} \left(\Lambda^2 - 2m_s^2 \ln \frac{\Lambda}{m_s} + \dots \right)$$

Supersymmetry

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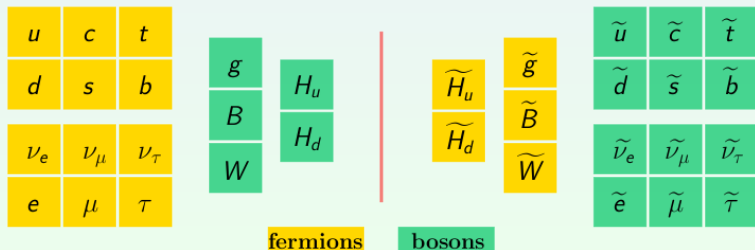
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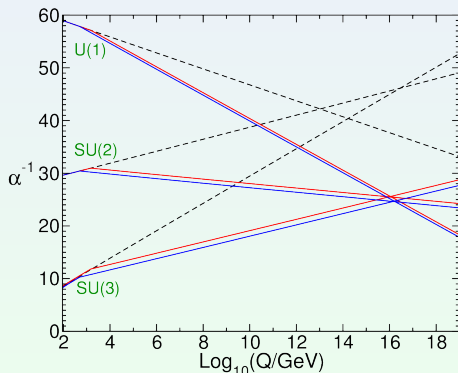
Not yet observed \Rightarrow SUSY breaking

Minimal Supersymmetric Standard Model (MSSM) :



Supersymmetry

- * Fermions \Leftrightarrow bosons \Rightarrow solution to the Hierarchy problem
- * Unification at GUT scale



Gauge coupling unification, Stephen P. Martin, arXiv :hep-ph/9709356

Supersymmetry

- * Fermions \leftrightarrow bosons \Rightarrow solution to the Hierarchy problem
- * Unification at GUT scale
- * LSP/DM (R-Parity)

The lightest supersymmetric particle (LSP) is stable, at the GeV-TeV scale, and can be weakly charged under the SM gauge group

\Rightarrow **DM candidates in supersymmetric models**

Supersymmetry

- * Fermions \Leftrightarrow bosons \Rightarrow solution to the Hierarchy problem
- * Unification at GUT scale
- * LSP/DM (supersymmetry breaking, R-Parity)
- * Examples :

u	c	t
d	s	b
ν_{eL}	$\nu_{\mu L}$	$\nu_{\tau L}$
e	μ	τ

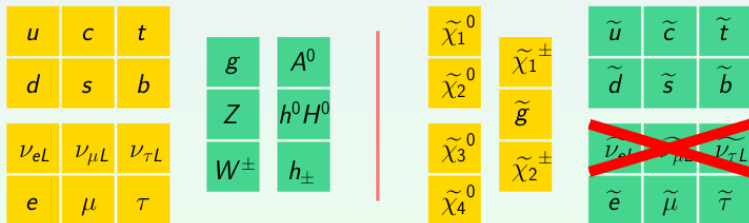
g	A^0
Z	$h^0 H^0$
W^\pm	h_\pm

$\tilde{\nu}_1^0$	$\tilde{\chi}_1^\pm$
$\tilde{\chi}_2^0$	\tilde{g}
$\tilde{\chi}_3^0$	$\tilde{\chi}_2^\pm$
$\tilde{\chi}_4^0$	

\tilde{u}	\tilde{c}	\tilde{t}
\tilde{d}	\tilde{s}	\tilde{b}
$\tilde{\nu}_{eL}$	$\tilde{\nu}_{\mu L}$	$\tilde{\nu}_{\tau L}$
\tilde{e}	$\tilde{\mu}$	$\tilde{\tau}$

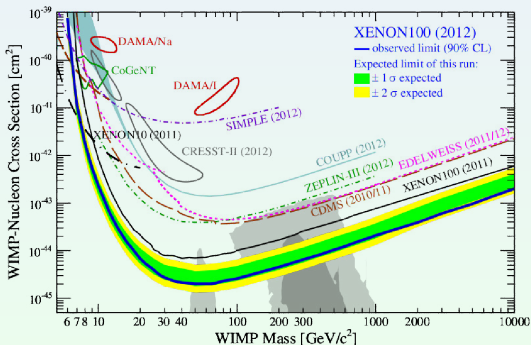
Supersymmetry

- * Fermions \Leftrightarrow bosons \Rightarrow solution to the Hierarchy problem
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Supersymmetry

- * Constraints on SUSY/DM
 - * DM relic abundance
 - * Direct detection of DM



E. Aprile et al., XENON100 Collaboration, Phys. Rev. Lett. 109 :181301, arXiv :1207.5988

Supersymmetry

* Constraints on SUSY/DM

- * DM relic abundance

- * Direct detection of DM

- * Indirect detection of DM (search for anomalous features in cosmic rays like $\gamma, \nu, e^+, \bar{p}$)

- ✗ “Background drawback” : ID depends on the current knowledge of astrophysical sources

- ✓ Remove carefully known (modelled) background

- ✓ Clear features not mimicked by astrophysical sources

- * A huge number of data validates the modelling of astrophysical background sources in the GeV-TeV range :

- absence of anomalies in the \bar{p} spectrum less exploited ⇒ **Set constraints**

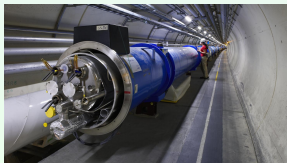
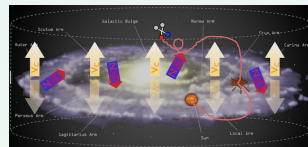
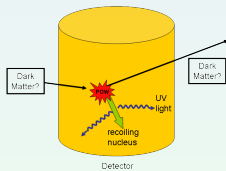
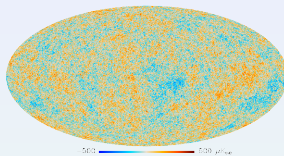
Supersymmetry

* Constraints on SUSY/DM

- * DM relic abundance
- * Direct detection of DM
- * Indirect detection of DM
- * Collider constraints
 - * LEP \Rightarrow charged sparticles
 - * LHC \Rightarrow coloured sparticles
 - * Low energy observables

$$\mathcal{B}(\bar{\mathbf{B}}^0 \rightarrow \mathbf{X}_s \gamma), \mathcal{B}(\mathbf{B}_s^0 \rightarrow \mu^+ \mu^-), \mathcal{B}(\mathbf{B}^\pm \rightarrow \tau^\pm \nu_\tau), \\ \Delta M_{d,s}, \delta a_\mu, \Delta \rho, \dots$$

Supersymmetry



Neutralino DM in the (N)MSSM

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Supersymmetric inflaton

* NUHM2 (Non-Universal Higgs Masses type 2)

- * Supersymmetric model with gravity-mediated supersymmetry breaking based on the MSSM
- * Most popular : mSUGRA/CMSSM, universal scalar masses is assumed, free parameters :

$$m_0, m_{1/2}, A_0, \tan\beta \text{ and } \text{sign}(\mu)$$

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- * We considered a non-universal scalar masses model, with $m_0^2 \neq m_{H_u}^2 \neq m_{H_d}^2$ (H. Baer et al [hep-ph/0504001], J. R. Ellis et al [hep-ph/0210205])
 - ✓ Easier to reach $m_{h^0} = 125$ GeV, increase DM annihilation rates with higgsino LSP
- * NUHM2 free parameter :

$$m_0, m_{1/2}, A_0, \tan \beta, \mu \text{ and } m_{A^0}$$

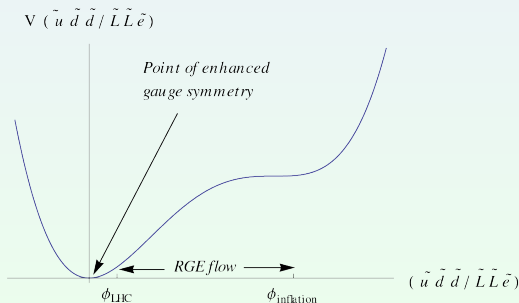
Supersymmetric inflaton

- ★ NUHM2

- ★ $\tilde{L}\tilde{L}\tilde{e}$ and $\tilde{u}\tilde{d}\tilde{d}$

- ★ Inflaton, scalar field whose flat direction potential (with a non-negligible slope) leads to the end of the inflation phase
- ★ Charged under the visible sector of the particle physics model considered, i.e. NUHM2

- ★ $\Rightarrow \tilde{L}\tilde{L}\tilde{e}$ and $\tilde{u}\tilde{d}\tilde{d}$ D-terms can be such candidates
(R. Allahverdi et al, [hep-ph/0610134], [hep-ph/0605035])



Method and constraints

Constraints imposed on a scan made using Markov Chain Monte Carlo method :

- On $\tilde{L}\tilde{L}e$ and $\tilde{u}\tilde{d}\tilde{d}$, explain the observed temperature anisotropy in the CMB with :
 - The amplitude of density perturbations $\delta_H = \frac{8}{\sqrt{5}\pi} \frac{m_\phi M_P}{\phi_0^2} \frac{1}{\Delta^2} \sin^2[\mathcal{N}_{\text{COBE}}\sqrt{\Delta^2}]$,
 $\Delta^2 \equiv 900\alpha^2 \mathcal{N}_{\text{COBE}}^{-2} \left(\frac{M_P}{\phi_0}\right)^4$, $\mathcal{N}_{\text{COBE}} \sim 50$
 - The scalar spectral index n_s of the corresponding power spectrum
 $n_s = 1 - 4\sqrt{\Delta^2} \cot[\mathcal{N}_{\text{COBE}}\sqrt{\Delta^2}]$,

Constraint	Value/Range
m_{h^0} (GeV)	[115.5, 127]
$\Omega_{\chi_1^0} h^2$	[0.1088, 0.1158]
$\mathcal{B}(\bar{B}^0 \rightarrow X_s \gamma) \times 10^4$	3.55
$\delta a_\mu \times 10^{10}$	28.7
$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) \times 10^9$	4.5
$\Delta\rho$	0.002
$R_{B^\pm \rightarrow \tau^\pm \nu_\tau} (\frac{\text{NUHM2}}{\text{SM}})$	2.219
$Z \rightarrow \chi_1^0 \chi_1^0$ (MeV)	1.7
$\sigma_{e^+e^- \rightarrow \chi_1^0 \chi_{2,3}^0}$ $\times \mathcal{B}(\chi_{2,3}^0 \rightarrow Z \chi_1^0)$ (pb)	1

- On NUHM2 model in general :

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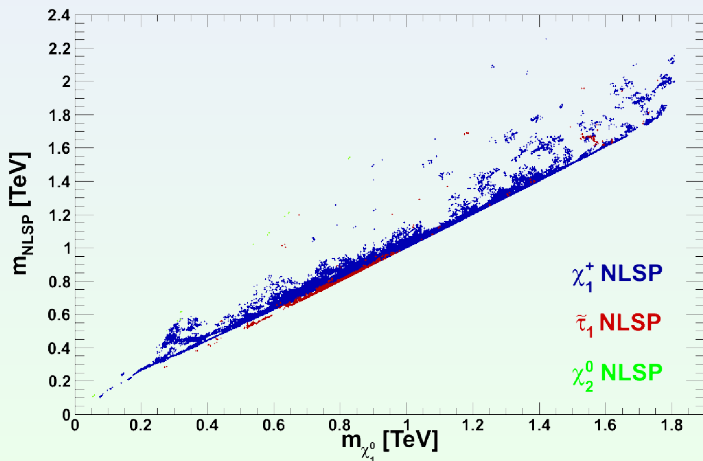
* On NUHM2 model in general :

* Compute total likelihood of a point M in NUHM2 parameter space :

if $\mathcal{L}_{\text{tot}}^M > \frac{\mathcal{L}_{\text{tot}}^m}{p}$, with $p \in [1, 1 - \ln \mathcal{L}_{\text{tot}}^m] \Rightarrow$ keep M

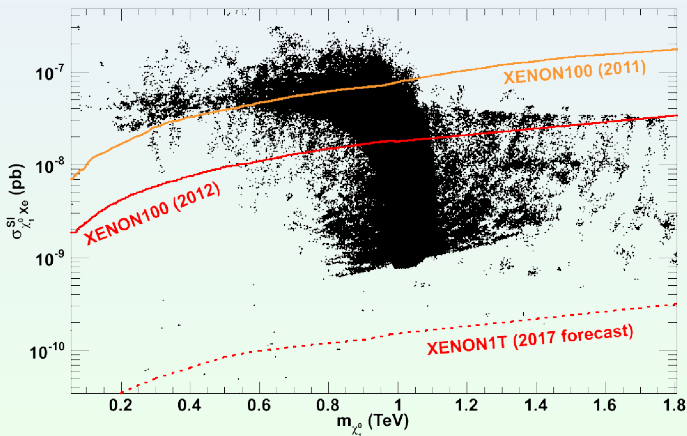
Results

- * Hard to find bino-like LSP + correct LSP relic density (mass mainly close to $m_{A^0}/2$)
- * Get mainly higgsino-like LSP, degeneracy between $\chi_{1,2}^0$ and χ_1^\pm



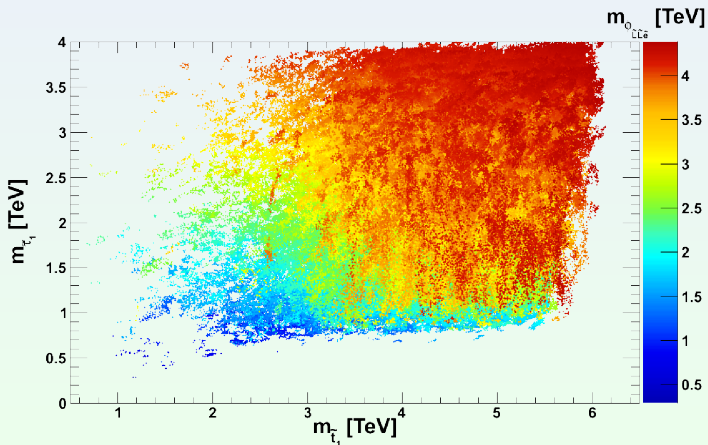
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- * Informations on inflaton mass if we discover lightest stop/stau at LHC



DM ID limits on the LSP-NLSP mass degeneracy

Possibility to set stringent constraints on DM properties by looking at DM annihilation into W^\pm , when LSP and NLSP are mass degenerate (difficult at the LHC), using FERMILAT AND PAMELA data

- * From γ -rays : FERMILAT analysis of the diffuse γ -ray emission from dwarf spheroidal galaxies (Ackermann et al, *Phys. Rev. Lett.* **107** (2011) 241302)
- * From \bar{p} : derived bounds from PAMELA antiprotons data using several approaches

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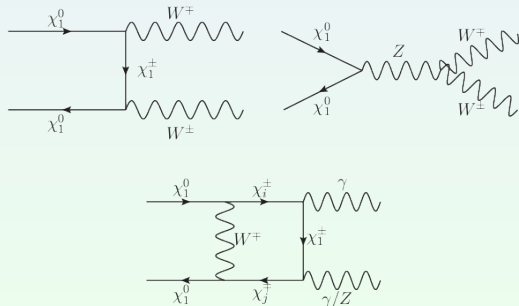
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⇒ A “simplified” version of the pMSSM (phenomenological MSSM)

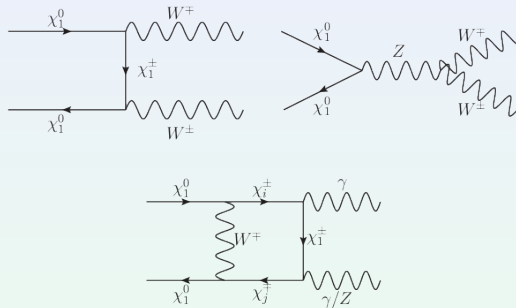
Aim : dominant neutralino DM annihilation channels into gauge bosons



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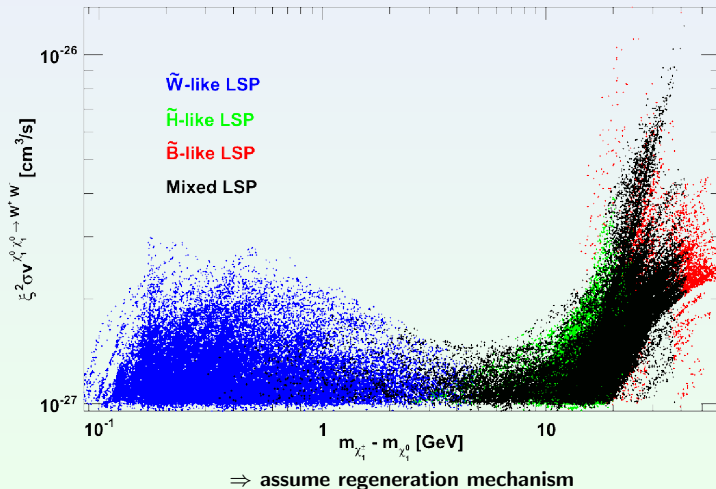
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- * All sfermion masses + CP-odd Higgs boson are set to 2 TeV (except for the third generation of squarks, to get $m_{h0} \sim 125$ GeV), light chargino/neutralino ($m_{\chi_1^0} < 500$ GeV) such that the mass splitting $\Delta m = m_{\chi_1^\pm} - m_{\chi_1^0}$ is small
- * MCMC scan
- * **How powerful are the \bar{p}/γ -ray limits on excluding parts of pMSSM parameter space and Δm values?**

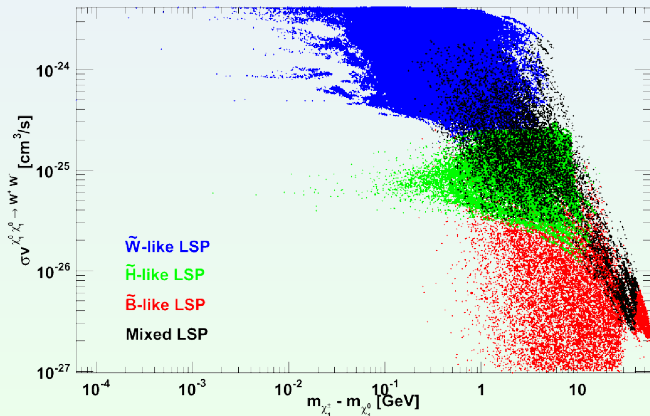
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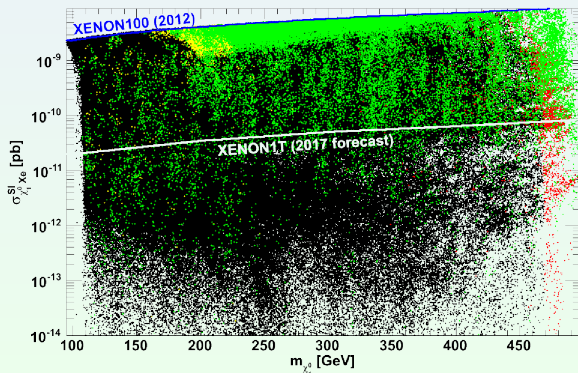
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- * ID constrains scenarios with $\Delta m \lesssim 20$ GeV, DM relic density being regenerated at 100%
- * If $m_{\chi_1^0} < 500$ GeV and $\Delta m < 0.25$ GeV wino DM ruled out



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- * If $m_{\chi_1^0} < 500$ GeV and $\Delta m < 0.25$ GeV wino DM ruled out
- * No explanation of the “130 GeV line” in this simplified pMSSM
- * ID constraints really competitive with direct detection experiments



OK
 Fermi-LAT
 Pamela
 both

NMSSM and SUSY searches @ LHC

- * Adding a singlet of SM gauge symmetry to solve the μ problem of the MSSM
- * $\mathcal{W}_{\text{MSSM}} = \tilde{u}_R^* y_u \tilde{Q} H_u - \tilde{d}_R^* y_d \tilde{Q} H_d - \tilde{e}_R^* y_e \tilde{L} H_d + \mu H_u H_d$
 $\Rightarrow \mathcal{W}_{\text{NMSSM}} = \mathcal{W}_{\text{MSSM}}|_{\mu=0} + \lambda S H_u H_d + \frac{1}{3} \kappa S^3$
- * 2 CP-odd Higgs boson (a_1, a_2), 3 CP-even Higgs boson (h_1, h_2, h_3)
 $m_{h_1}^2|_{\text{tree}} \lesssim M_Z^2 \cos^2 2\beta + \frac{\lambda^2}{2} v^2$
 \Rightarrow less fine tuned $m_{h_1} \sim 125$ GeV
- * 5 neutralinos χ_i^0 in the basis $(\tilde{B}, \tilde{W}^3, \tilde{H}_d^0, \tilde{H}_u^0, \tilde{S})$

NMSSM and SUSY searches @ LHC

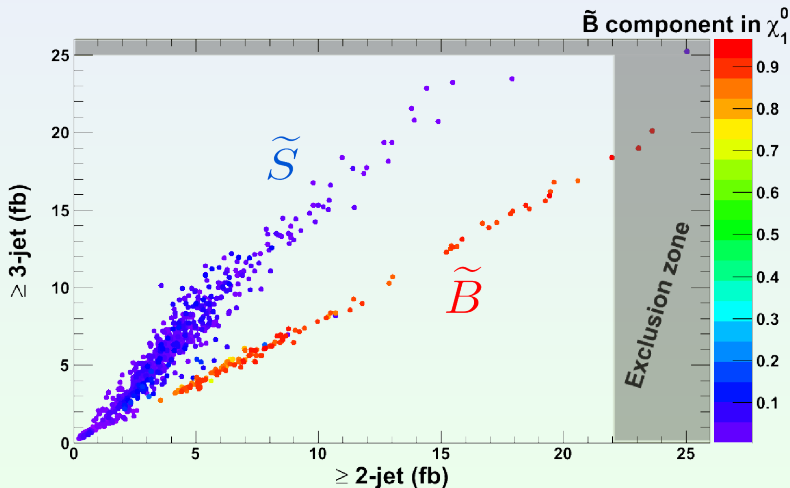
- * Adding a singlet of SM gauge symmetry to solve the μ problem of the MSSM
- * $\mathcal{W}_{\text{MSSM}} = \tilde{u}_R^* y_u \tilde{Q} H_u - \tilde{d}_R^* y_d \tilde{Q} H_d - \tilde{e}_R^* y_e \tilde{L} H_d + \mu H_u H_d$
 $\Rightarrow \mathcal{W}_{\text{NMSSM}} = \mathcal{W}_{\text{MSSM}}|_{\mu=0} + \lambda S H_u H_d + \frac{1}{3} \kappa S^3$
- * 2 CP-odd Higgs boson (a_1, a_2), 3 CP-even Higgs boson (h_1, h_2, h_3)
 $m_{h_1}^2|_{\text{tree}} \lesssim M_Z^2 \cos^2 2\beta + \frac{\lambda^2}{2} v^2$
 \Rightarrow less fine tuned $m_{h_1} \sim 125$ GeV
- * 5 neutralinos χ_i^0 in the basis $(\tilde{B}, \tilde{W}^3, \tilde{H}_d^0, \tilde{H}_u^0, \tilde{S})$
- * Using results of a previous work (D. Albornoz Vasquez et al., arXiv :1107.1614, arXiv :1201.6150) with constraints on DM, B and Higgs physics to define the relevant NMSSM parameter space
- * Motivated by hints of a signal in direct detection experiments (DAMA/Libra, arXiv :1002.1028; CoGeNT, arXiv :1201.6150)
 \Rightarrow light DM (\tilde{B} or \tilde{S}) scenarios (mostly light a_1 and/or h_1)

NMSSM and SUSY searches @ LHC

- * Searches for exotic particles are now reaching a high level of exclusion that allow to reject a wide class of models
but limits obtained assuming simplified models of New Physics
⇒ what about the NMSSM?
- * Example of the exclusion limit coming from the ATLAS 1.04 fb^{-1} search for squarks and gluinos via jets and missing E_T
- * In general exclude squarks lighter than 0.6 - 1 TeV and gluinos below 0.5 TeV in the constrained MSSM via $\tilde{q} \rightarrow q\chi_1^0$ and $\tilde{g} \rightarrow q\bar{q}\chi_1^0$ decays
 - * Applying SUSY searches@LHC with ATLAS's 1.04 fb^{-1} 0-lepton jets + missing E_T search using Herwig++ 2.5.1 and RIVET 1.5.2
⇒ Are ATLAS limits so constraining?

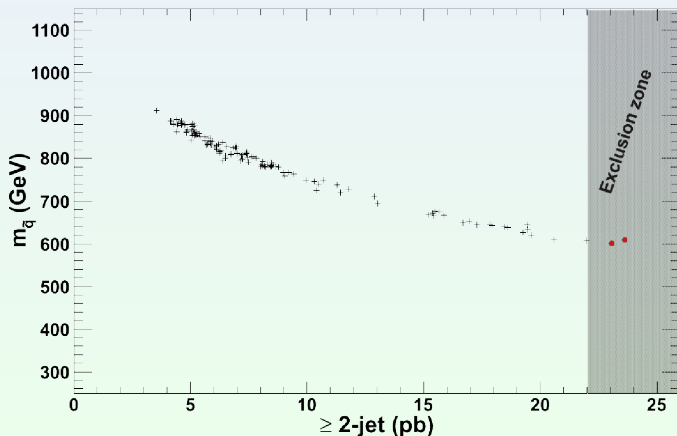
Results

- * Reduced acceptance into jets + missing E_T search channels and more jets for \tilde{S} LSP
- * $\tilde{q} \rightarrow q + (\chi_2^0 \rightarrow \chi_1^0 + (f\bar{f} \text{ or } a_1 \text{ or } h_1))$



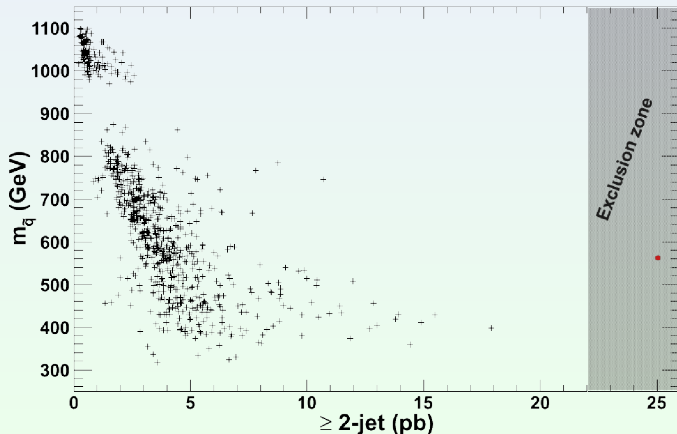
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- * Usual exclusion (\tilde{B} -like LSP) :



Results

- * Reduced acceptance into jets + missing E_T search channels and more jets for \tilde{S} LSP
- * $\tilde{q} \rightarrow q + (\chi_2^0 \rightarrow \chi_1^0 + (f\bar{f} \text{ or } a_1 \text{ or } h_1))$
- * 300 GeV squarks allowed when (\tilde{S} -like LSP) :



U(1) extensions of the MSSM

1 Motivations

2 Neutralino DM in the (N)MSSM

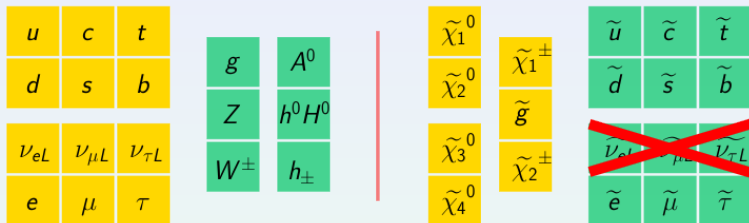
- C. Boehm, JDS, A. Mazumdar and E. Pukartas, Phys. Rev. D87 (2013) 023529, arXiv :1205.2815
- G. Bélanger, C. Boehm, M. Cirelli, JDS and A. Pukhov, JCAP 1211 (2012) 028, arXiv :1208.5009
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3 U(1) extensions of the MSSM

- G. Bélanger, JDS and A. Pukhov, JCAP 1112 (2011) 014, arXiv :1110.2414
- G. Bélanger, JDS et al., in preparation

4 Conclusions

Sneutrinos



- * Neutrino oscillations indicative of massive neutrinos \Rightarrow possibility to add right-handed (RH) neutrino fields
 - \Rightarrow Extensions of the MSSM with RH (s)neutrino can provide DM candidate
- * Here RH neutrino mass generated by introducing Dirac mass terms
 - \Rightarrow supersymmetric partner can be at the TeV scale
- * **This candidate couples to new vector, scalar field by adding a new abelian gauge symmetry \Rightarrow the UMSSM**

The model

- * **Symmetry group** : $SU(3)_c \times SU(2)_L \times U(1)_Y \times U'(1)$

Coupling constants : g_3, g_2, g_Y and $g'_1 = \sqrt{\frac{5}{3}} g_Y$

- * $U'(1)$ stems from string-inspired E_6 :

$E_6 \rightarrow SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_\chi \times U(1)_\psi \Rightarrow U'(1)$ charge :

$$Q' = \cos \theta_{E_6} Q'_\chi + \sin \theta_{E_6} Q'_\psi, \quad \theta_{E_6} \in [-\pi/2, \pi/2]$$

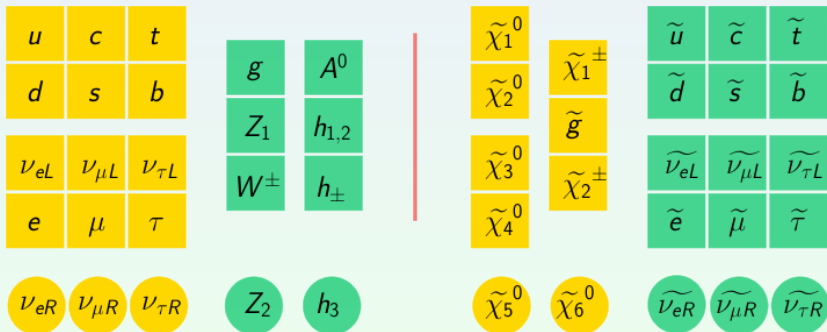
- * **Superpotential** :

$$\mathcal{W}_{UMSSM} = \mathcal{W}_{MSSM}|_{\mu=0} + \lambda \mathbf{S} \mathbf{H}_u \mathbf{H}_d + \tilde{\nu}_{R\nu}^* \tilde{\mathbf{L}} \mathbf{H}_u + \mathcal{O}(\text{TeV}s)$$

- * **As the NMSSM, this model solves the μ problem** : $\mu = \lambda \frac{v_s}{\sqrt{2}}$
- * **New D-terms for m_{h_1}**

The model

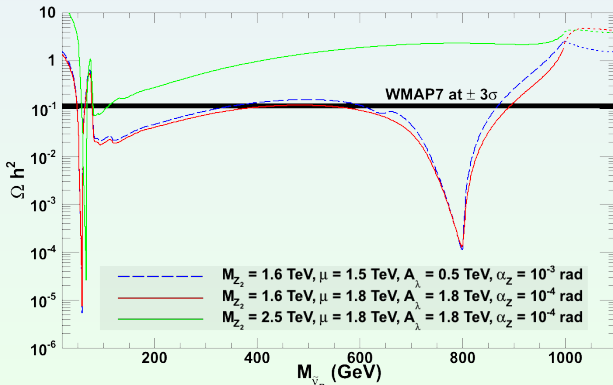
- * Gauge sector : Physical abelian gauge bosons : Z_1 and Z_2 , mixing between the Z of the SM and the Z' , α_Z is the mixing angle $\Rightarrow \tan \beta$ constrained
- * Gauginos sector : 6 neutralinos in the basis $(\tilde{B}, \tilde{W}^3, \tilde{H}_d^0, \tilde{H}_u^0, \tilde{S}, \tilde{B}')$
- * To sum up :



WIMP annihilation

Parameter space regions with $\Omega_{\text{WIMP}} h^2 \approx 0.1 \Rightarrow$ need to increase the annihilation cross section : interesting WIMP mass from 50 GeV to TeV-scale :

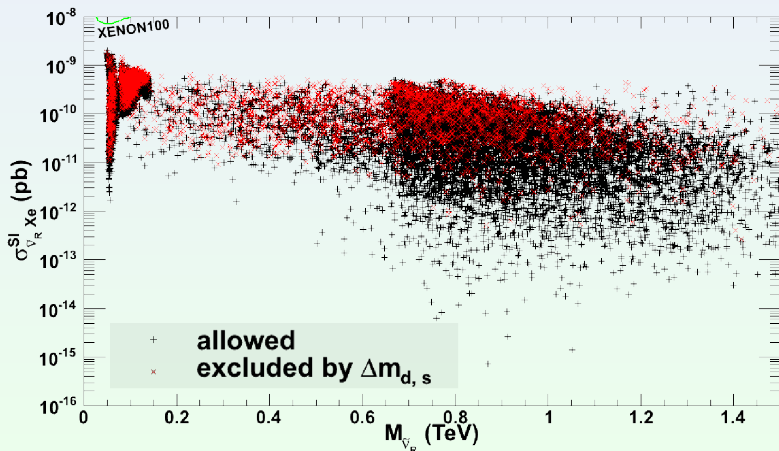
- * WIMP mass near $m_{h_1}/2$
- * WIMP mass near $M_{Z_2}/2$ (also $m_{h_i}/2$)
- * WIMP mass near $m_{h_i}/2$ or above W pair threshold
- * Coannihilation processes (mainly higgsino-like)



Scattering on nucleons

For some $U'(1)$ models we can have a good suppression of the gauge boson or/and Higgs boson contribution

here $U(1)_{\psi} \Rightarrow \theta_{E_6} = \pi/2$

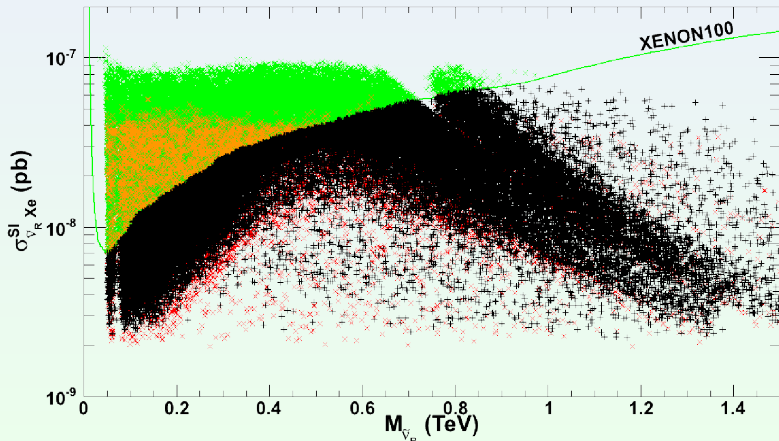


Scattering on nucleons

For other models, huge constraints on the parameter space appear

here $U(1)_\eta \Rightarrow \tan \theta_{E_6} = -\sqrt{5/3}$

OK, $\Delta m_{d,s}$, XENON100, both

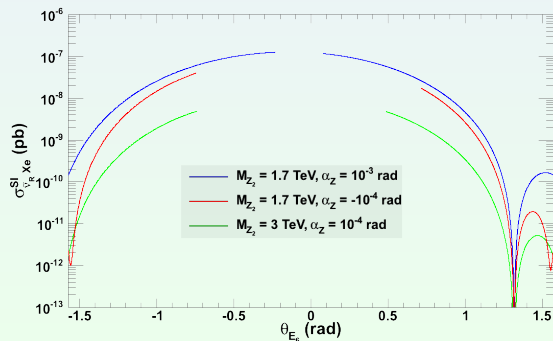


Scattering on nucleons

Abelian boson contribution to direct detection cross section :

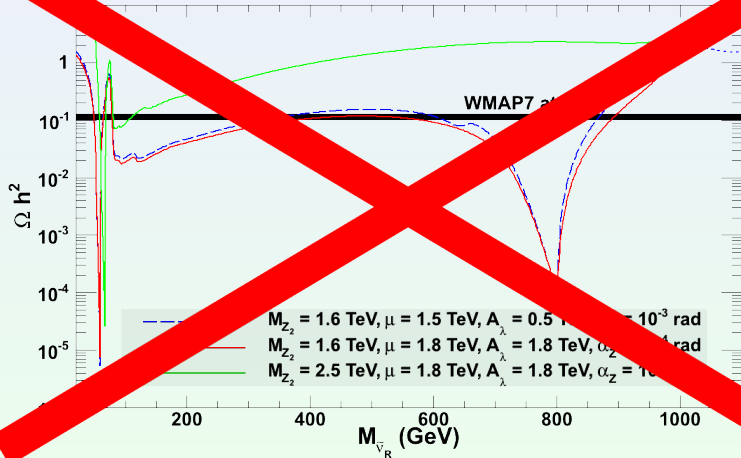
$$\sigma_{\tilde{\nu}_{RN}^{Z_1, Z_2}} = \frac{\mu_{\tilde{\nu}_{RN}}^2}{\pi} (\mathbf{g}'_1 Q'_\nu)^2 [(y(1 - 4s_W^2) + y')Z + (-y + 2y')(A - Z)]^2$$

$$\text{with } y = \frac{\mathbf{g}_Y \sin \alpha_Z \cos \alpha_Z}{4 \sin \theta_W} \left(\frac{1}{M_{Z_2}^2} - \frac{1}{M_{Z_1}^2} \right), \quad y' = -\frac{\mathbf{g}'_1 Q'_d}{2} \left(\frac{\sin^2 \alpha_Z}{M_{Z_1}^2} + \frac{\cos^2 \alpha_Z}{M_{Z_2}^2} \right)$$



\Rightarrow stringent constraints for small $|\theta_{E6}|$ because of Q'_d term

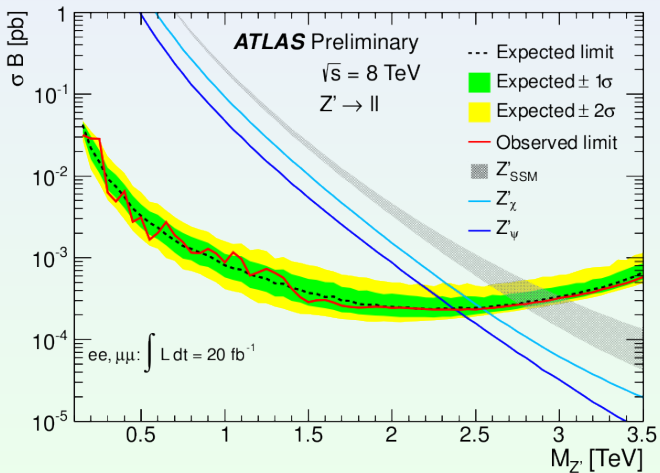
Need for an update



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* Updates :

- * New limits on M_{Z_2}



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- * DM observables (Planck satellite, update on XENON100 results)
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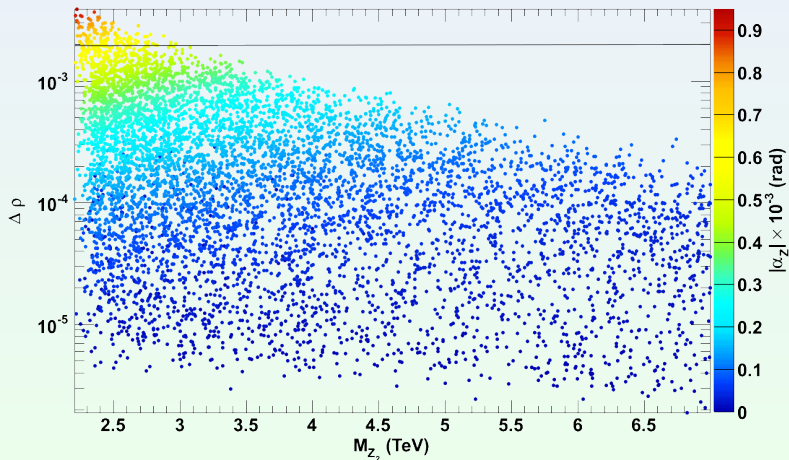
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* New inputs :

- * Higgs boson signal strengths + more low energy observables
⇒ Modification of the `NMSSMTools` code : `UMSSMTools`
- * Also neutralino as DM candidate
- * Relax relic abundance constraint
- * Third generation of sfermions allowed to be light

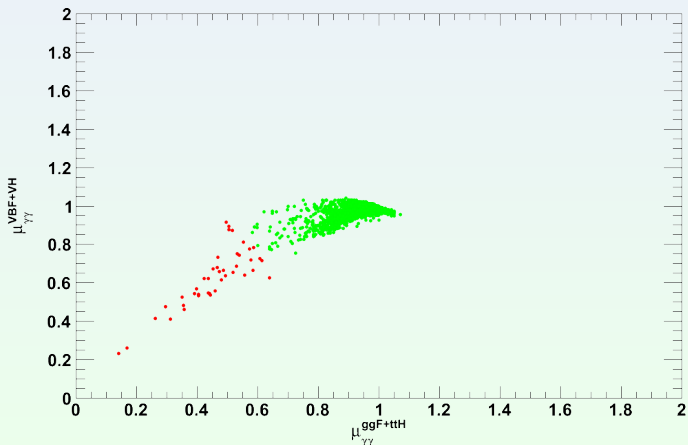
Results

- * Decrease of the upper bound on $|\alpha_Z|$
- * $\Delta\rho$: main new constraint for low energy observables



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- * Decrease of the upper bound on $|\alpha_Z|$
- * $\Delta\rho$: main new constraint for low energy observables
- * Constraints from 2σ signal strength ellipses derived in [G. Bélanger et al, arXiv :1306.2941](#)



Conclusions

- 1 Motivations
- 2 Neutralino DM in the (N)MSSM
 - C. Bøehm, JDS, A. Mazumdar and E. Pukartas, Phys. Rev. D87 (2013) 023529, arXiv :1205.2815
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Conclusions

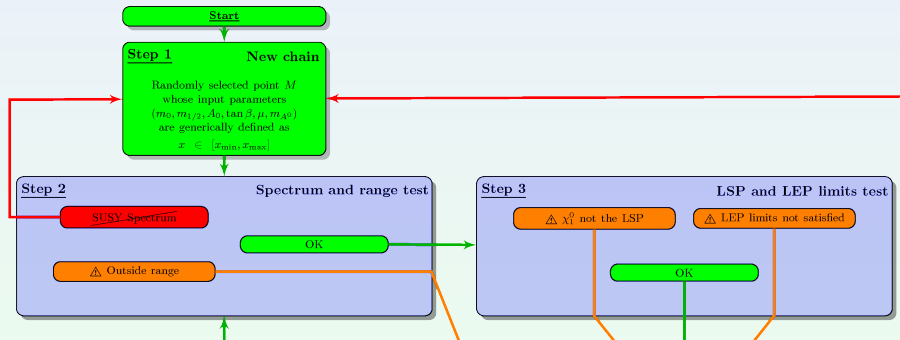
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Thanks !

BACKUP

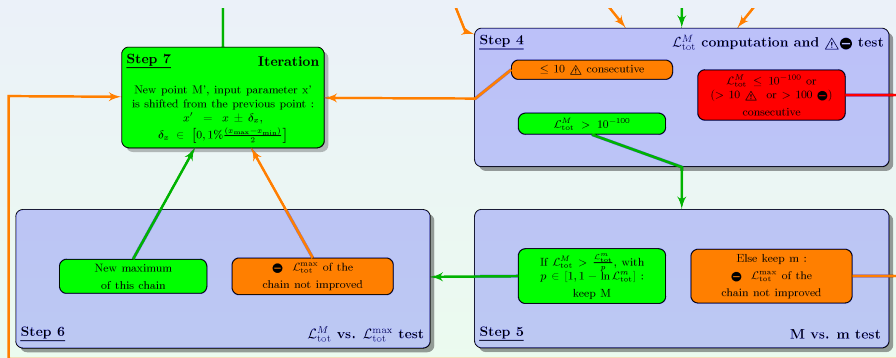
BACKUP

MCMC procedure



BACKUP

MCMC procedure



BACKUP

Models : $\tilde{L}\tilde{L}\tilde{e}$ and $\tilde{u}\tilde{d}\tilde{d}$

Inflaton, scalar field whose flat direction potential (with a non-negligible slope) leads to the end of the inflation phase

Charged under the visible sector of the particle physics model considered, i.e. NUHM2

$$\phi = \frac{\tilde{u} + \tilde{d} + \tilde{d}}{\sqrt{3}}, \quad \phi = \frac{\tilde{L} + \tilde{L} + \tilde{e}}{\sqrt{3}}$$

$$V(\phi) = \frac{1}{2} m_\phi^2 \phi^2 - A \frac{\lambda \phi^6}{6 M_{\text{P}}^3} + \lambda^2 \frac{\phi^{10}}{M_{\text{P}}^6}$$

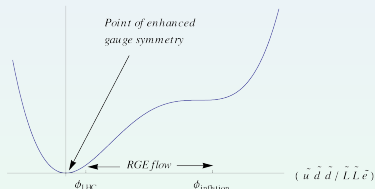
$$\phi_{\text{inflation}}^4 \simeq \frac{m_\phi M_{\text{P}}^3}{\lambda \sqrt{10}}, \quad V''(\phi_{\text{inflation}}) = 0$$

$\tilde{u}\tilde{d}\tilde{d}$ RGEs

$$\hat{\mu} \frac{d m_\phi^2}{d \hat{\mu}} = -\frac{1}{6\pi^2} (4 M_3^2 g_3^2 + \frac{2}{5} M_1^2 g_1^2),$$

$$\hat{\mu} \frac{d A}{d \hat{\mu}} = -\frac{1}{4\pi^2} \left(\frac{16}{3} M_3 g_3^2 + \frac{8}{5} M_1 g_1^2 \right)$$

$V(\tilde{u}\tilde{d}\tilde{d} / \tilde{L}\tilde{L}\tilde{e})$



$\tilde{L}\tilde{L}\tilde{e}$ RGEs

$$\hat{\mu} \frac{d m_\phi^2}{d \hat{\mu}} = -\frac{1}{6\pi^2} \left(\frac{3}{2} M_2^2 g_2^2 + \frac{9}{10} M_1^2 g_1^2 \right),$$

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BACKUP

ID constraints from \bar{p} W^\pm production leads also to abundant \bar{p} production (after hadronization)

⇒ \bar{p} flux produced by DM annihilation determined by :

$$\sigma_{\text{DM DM} \rightarrow \text{W}^+\text{W}^-}$$

$$m_{\text{DM}}$$

DM halo profile (here Einasto profile)

\bar{p} propagation parameters in the galactic halo :

Model	δ	\mathcal{K}_0 [kpc ² /Myr]	V_{conv} [km/s]	L [kpc]
MIN	0.85	0.0016	13.5	1
MED	0.70	0.0112	12	4
MAX	0.46	0.0765	5	15

⇒ We compare the sum of the astrophysical background flux and predicted \bar{p} flux originating from DM with the PAMELA data, 2 methods :

“Aggressive” procedure : fixed background (standard flux from T. Bringmann and P. Salati, *Phys. Rev. D* 75 (2007) 083006)

“Conservative” procedure : marginalized background, namely standard description of the background spectrum multiplied by $A(T/T_0)^p$ with :

$T = \bar{p}$ kinetic energy

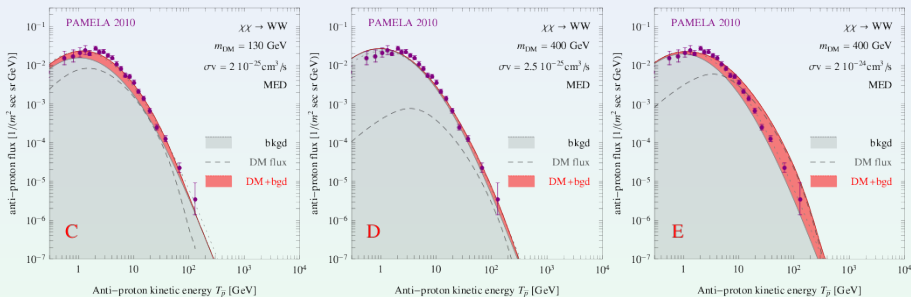
$T_0 = 30$ GeV : pivot energy

normalisation of the background spectrum : $0.6 < A < 1.4$

spectral index : $-0.1 < p < +0.1$

BACKUP

ID constraints from \bar{p}

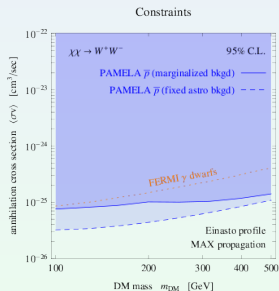
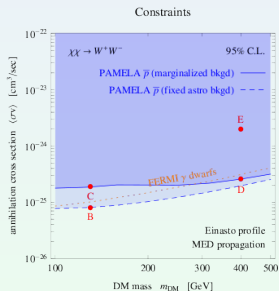
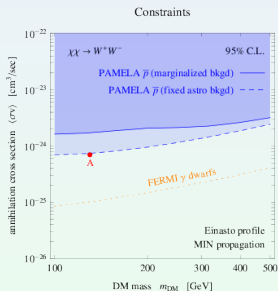


“Conservative” procedure approximately independent of m_{DM} : \bar{p} flux from heavy DM negligible at low energy, where PAMELA set very small error bars

We consider diffuse γ -ray constraints from dwarf spheroidal galaxies and \bar{p} constraints using ‘MED’ propagation parameters + marginalized background

BACKUP

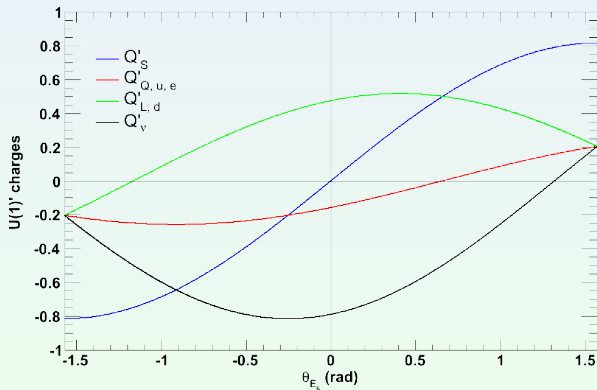
Generic bounds on DM annihilation into W^\pm



BACKUP

Higgs boson contribution to the direct detection cross section for $\tilde{\nu}_R$:

$$\mathbf{g}_{\tilde{\nu}_R \tilde{\nu}_R^* h_i} = -g_1'^2 Q'_\nu \left[v_d Q'_{H_d} Z_{hi1} + v_u Q'_{H_u} Z_{hi2} + v_s Q'_S Z_{hi3} \right]$$



\Rightarrow increase of the cross section for $\theta_{E_6} < 0$ because of Q'_ν