

**Université Joseph Fourier – Grenoble – France**  
Ecole Doctorale Terre – Univers – Environnement  
Observatoire des Sciences de l'Univers de Grenoble  
Laboratoire de Géophysique Interne et de Tectonophysique

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## **Thèse**

présentée à  
**L'Université Joseph Fourier - Grenoble I - France**  
pour obtenir le grade de  
**Docteur de l'Université Joseph Fourier,**  
spécialité « Terre – Univers – Environnement »

par  
**Aloé SCHLAGENHAUF**

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Identification des forts séismes passés sur les failles  
normales actives de la région Lazio-Abruzzo  
(Italie Centrale) par 'datations cosmogéniques' ( $^{36}\text{Cl}$ )  
de leurs escarpements  
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<b>ANNEXES</b>
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Thèse soutenue publiquement le 30 Septembre 2009 devant le jury composé de :

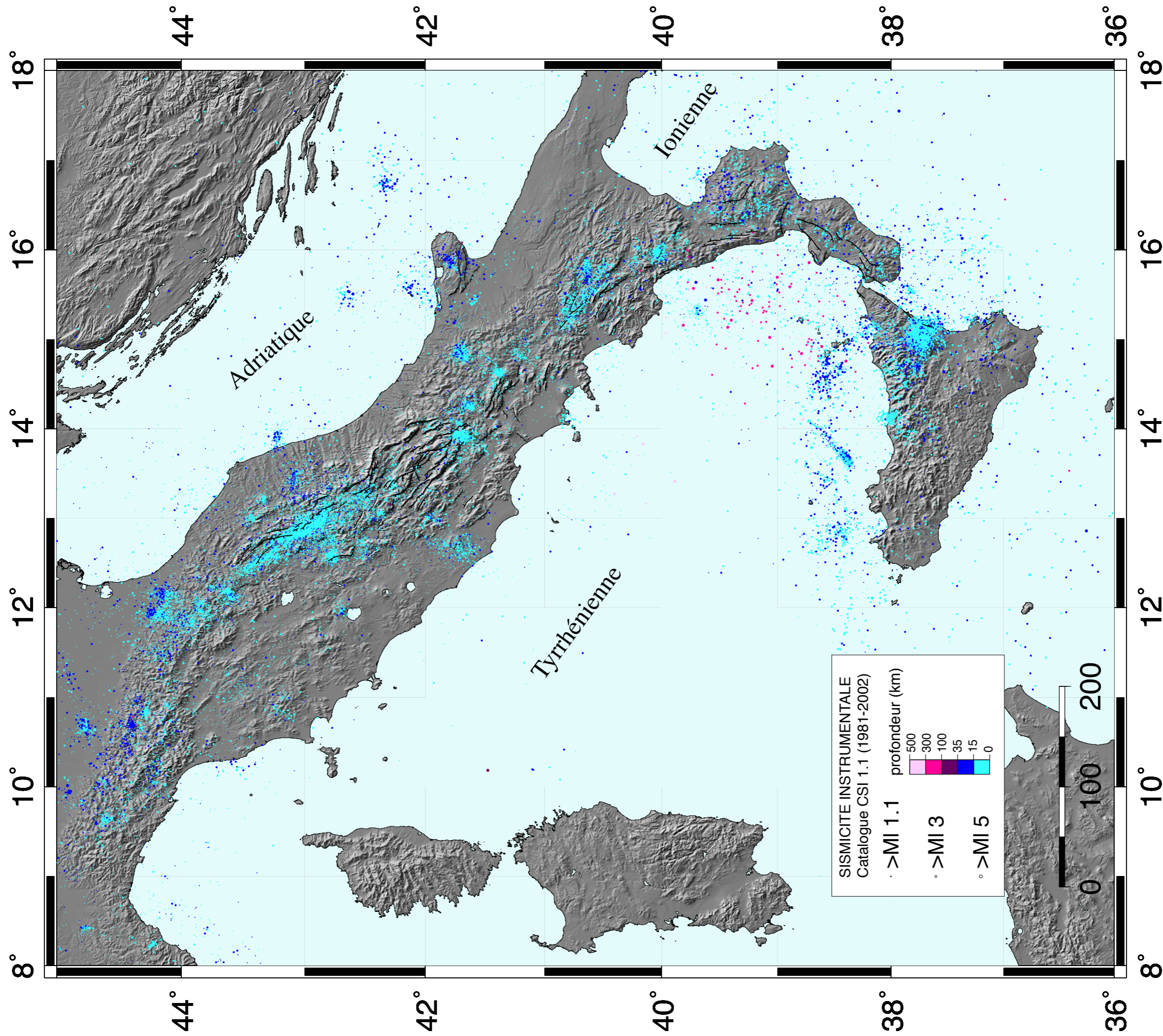
P. Galli	Protezione Civile, Roma, Italie	Rapporteur
G. Hilley	Université de Stanford, Etats-Unis	Rapporteur
M. Campillo	Université J. Fourier, LGIT, Grenoble	Président du Jury
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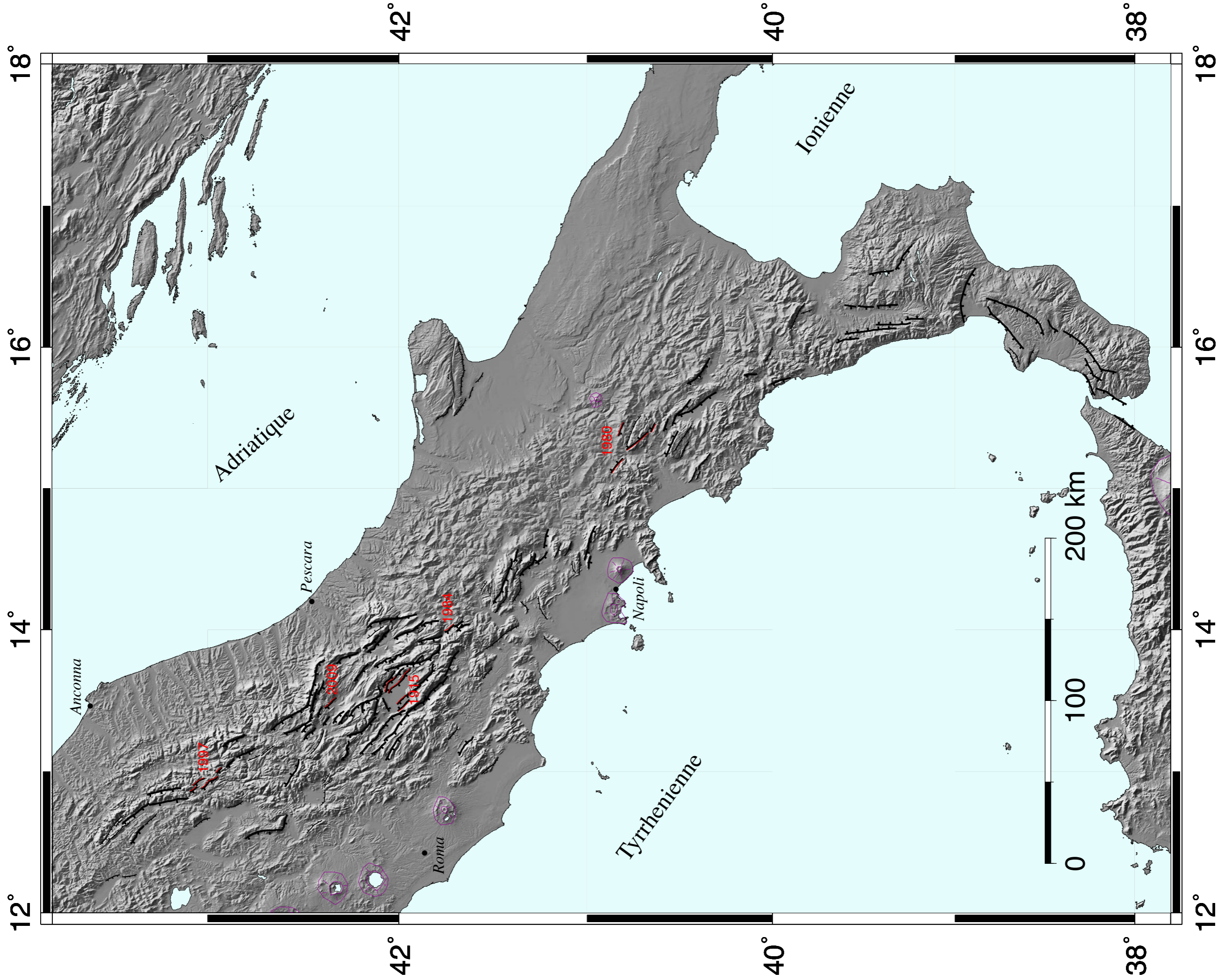


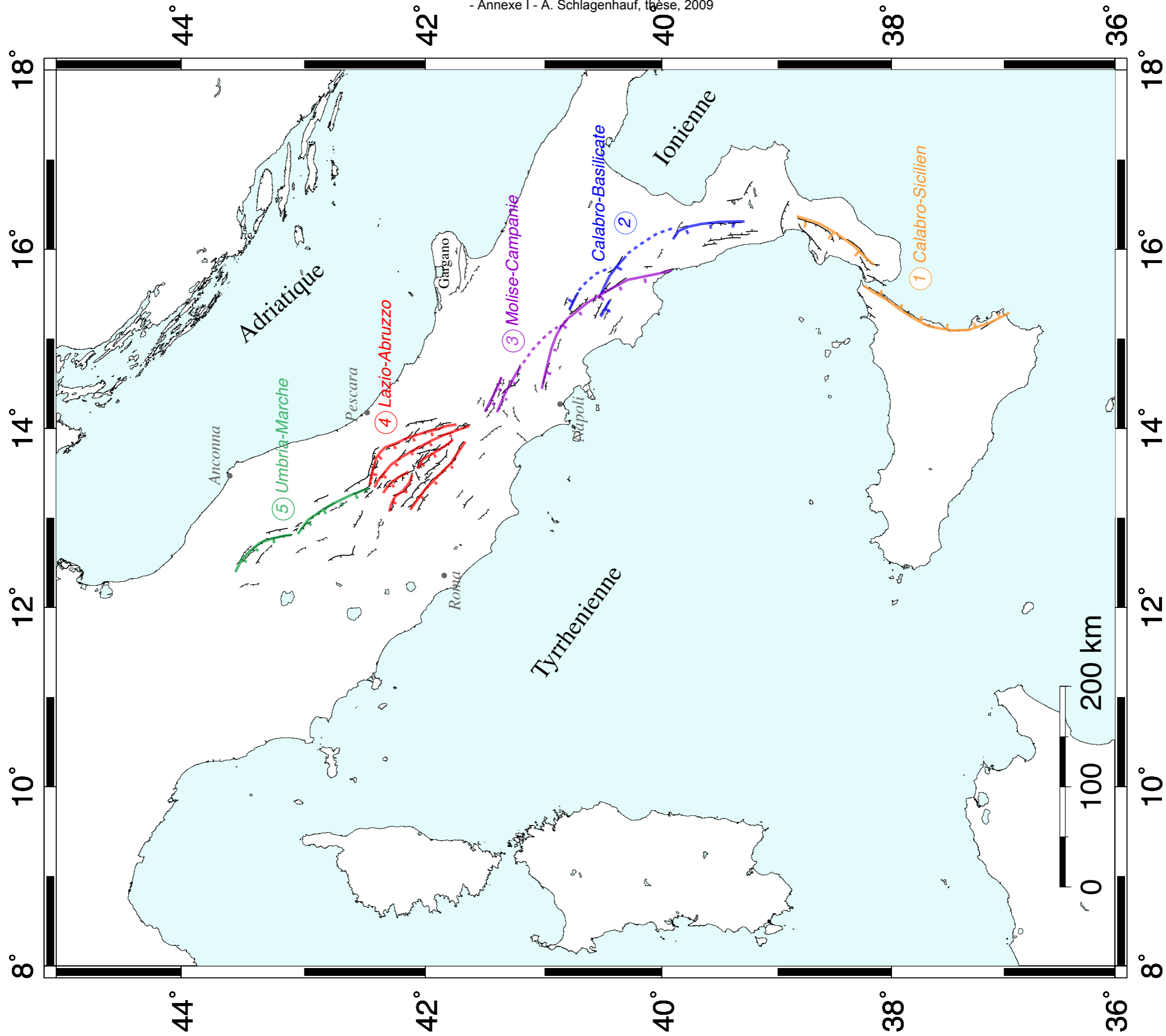
# **Annexe I**

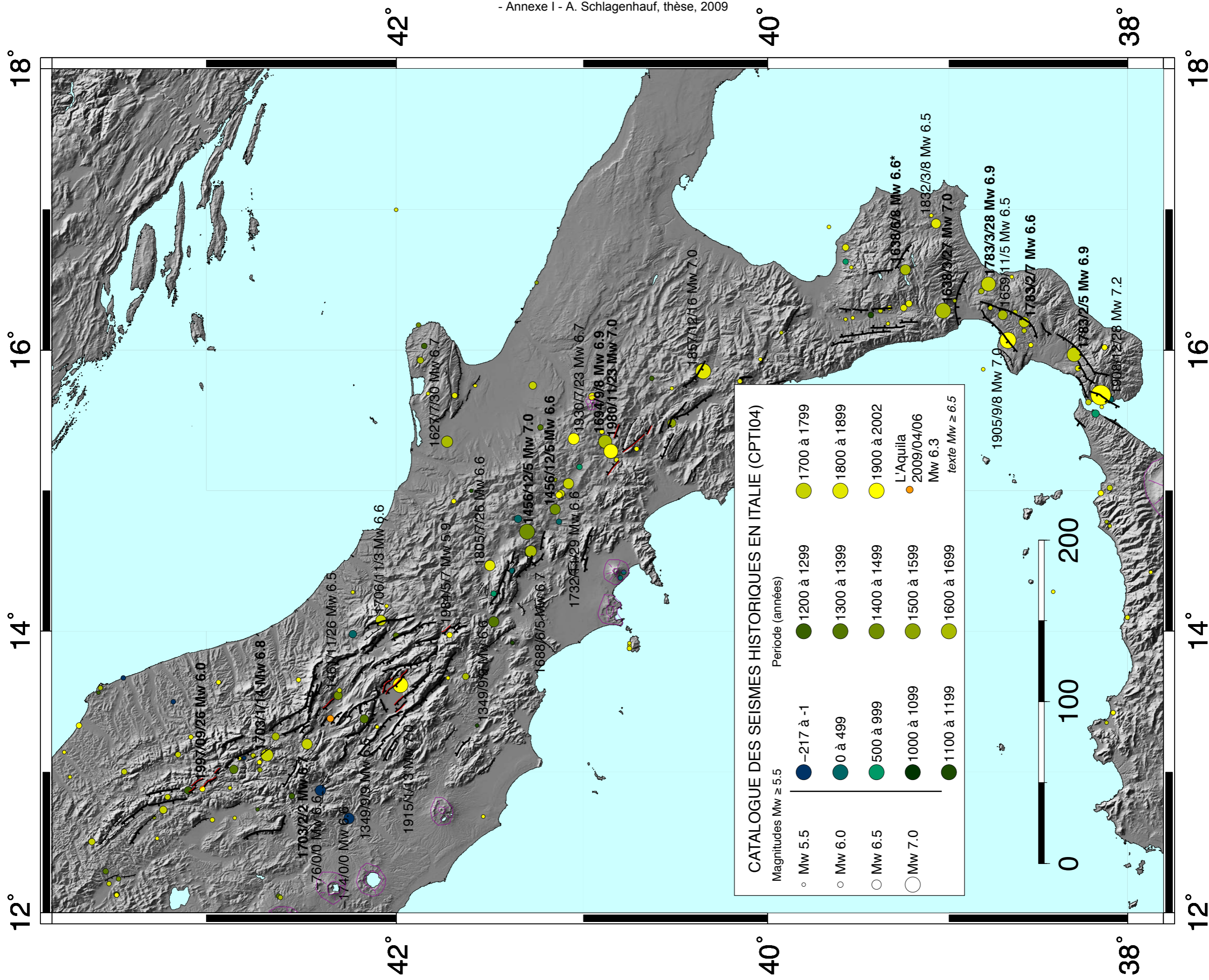
(du chapitre 3)



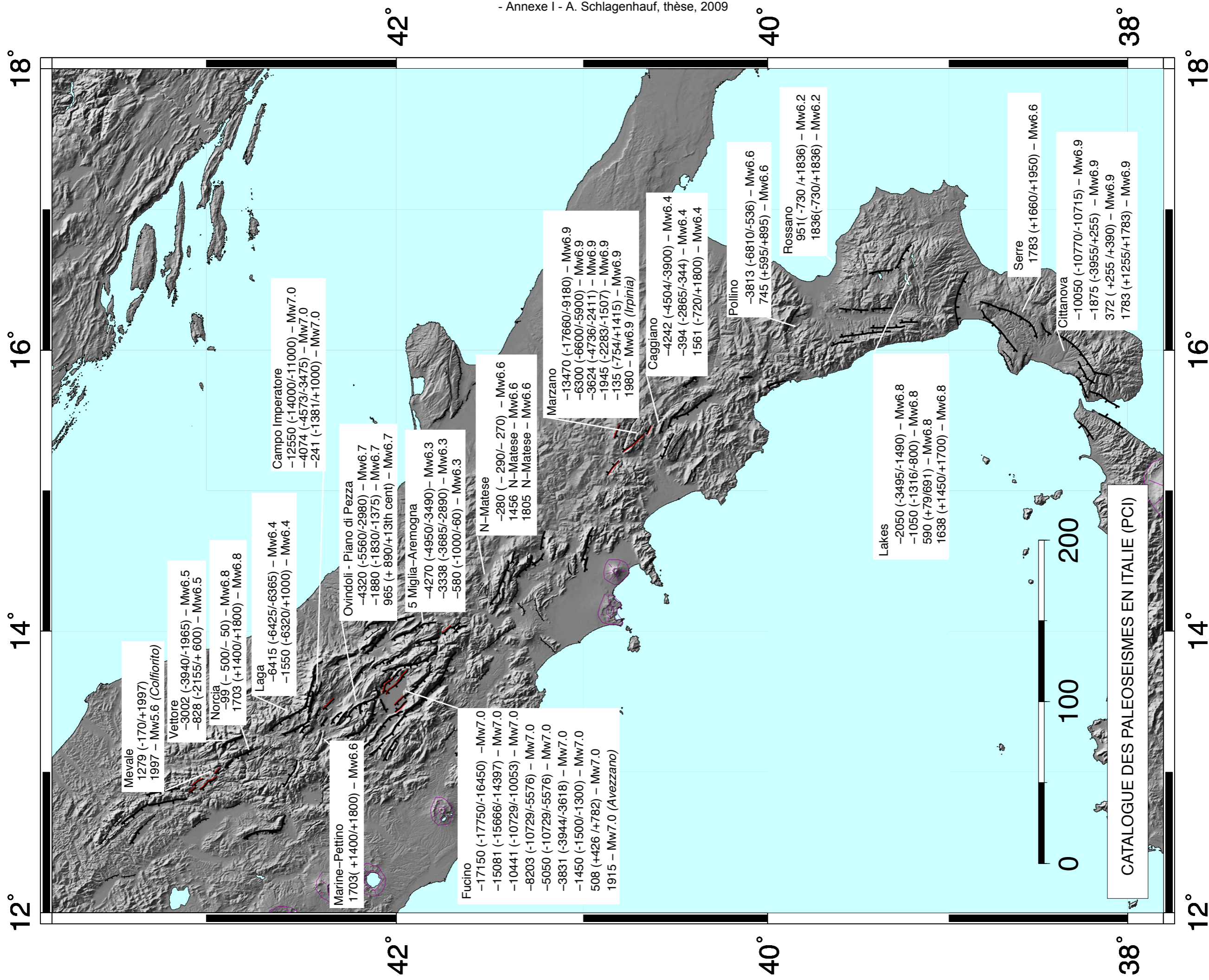


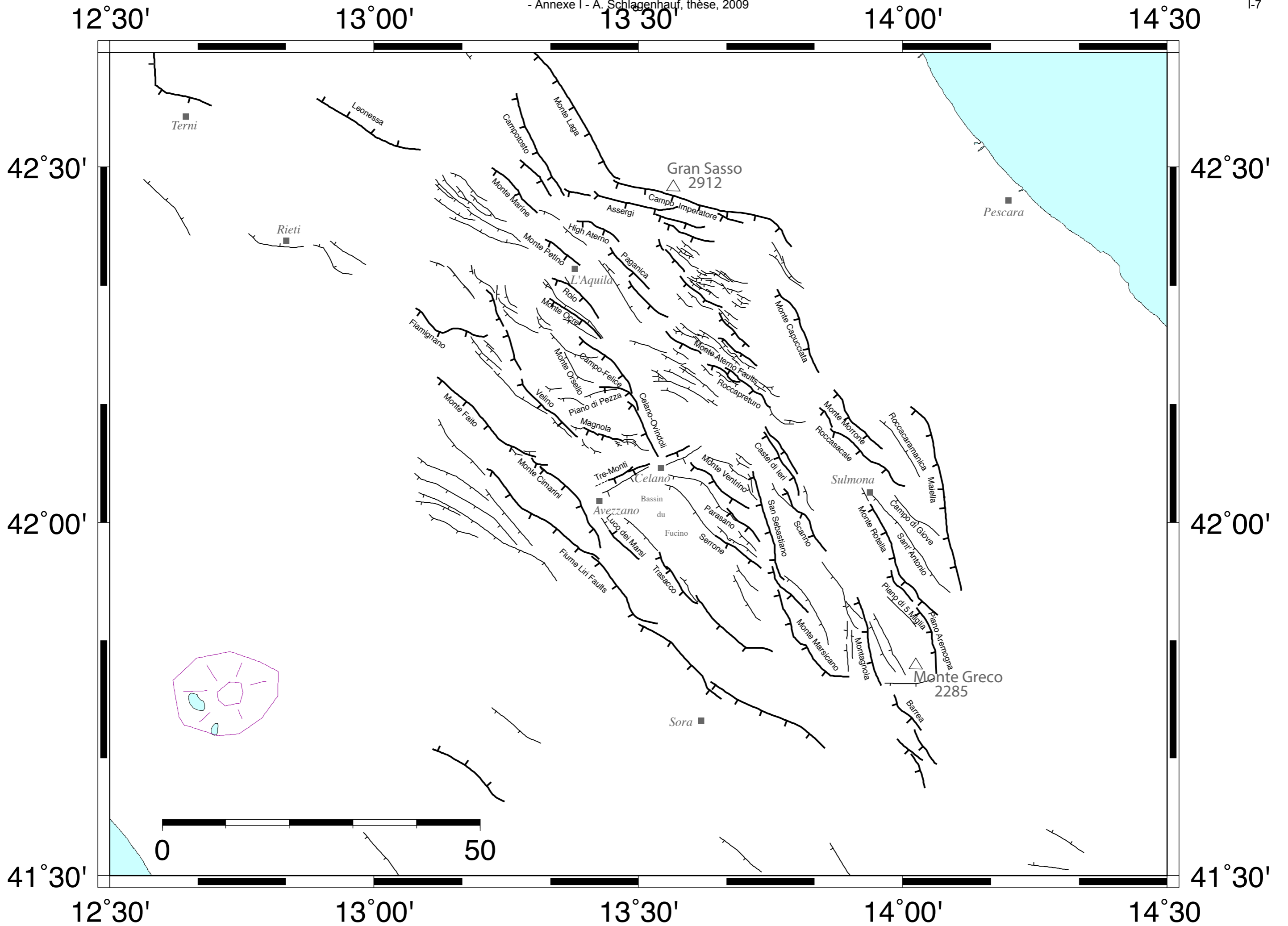


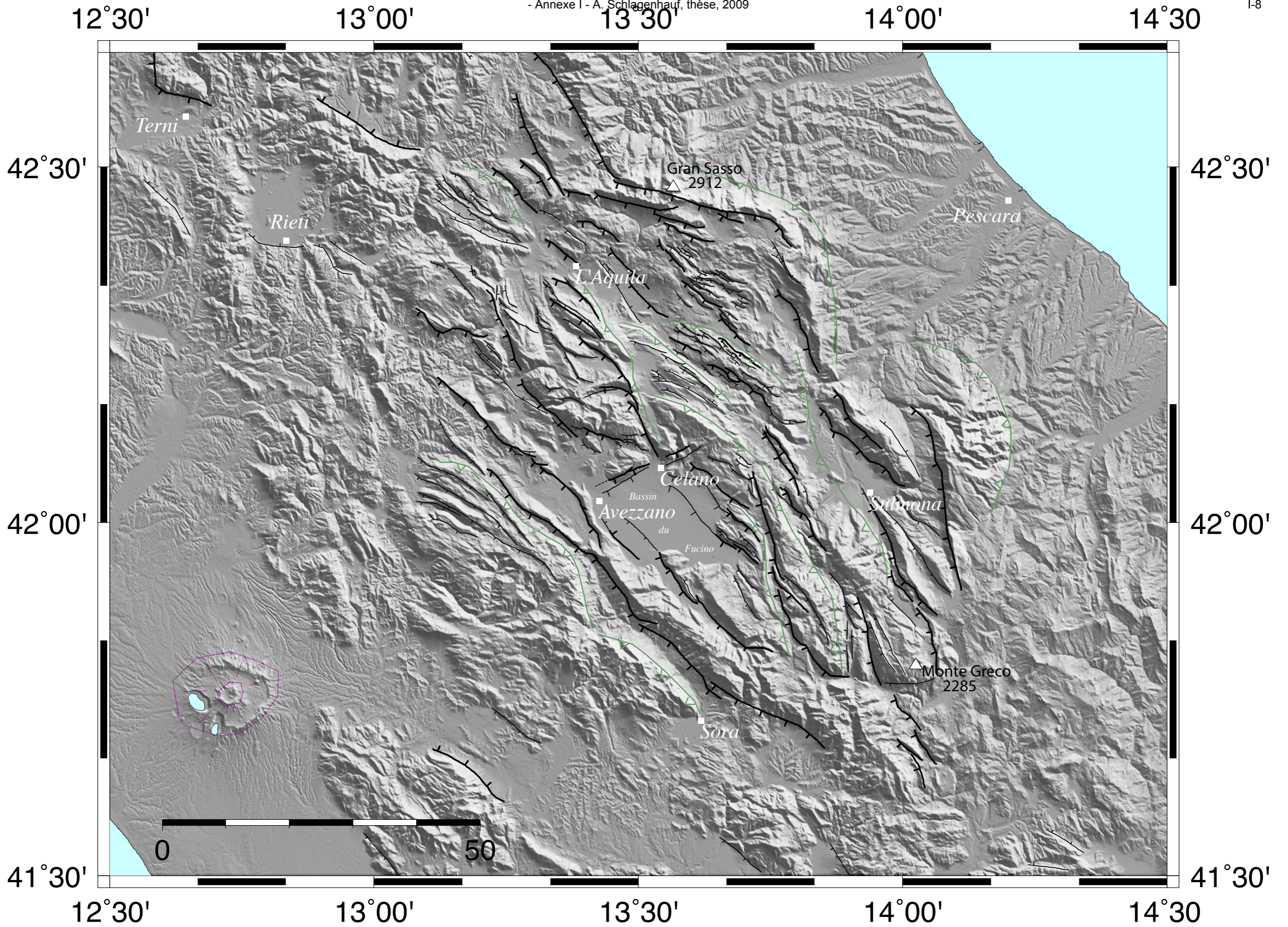


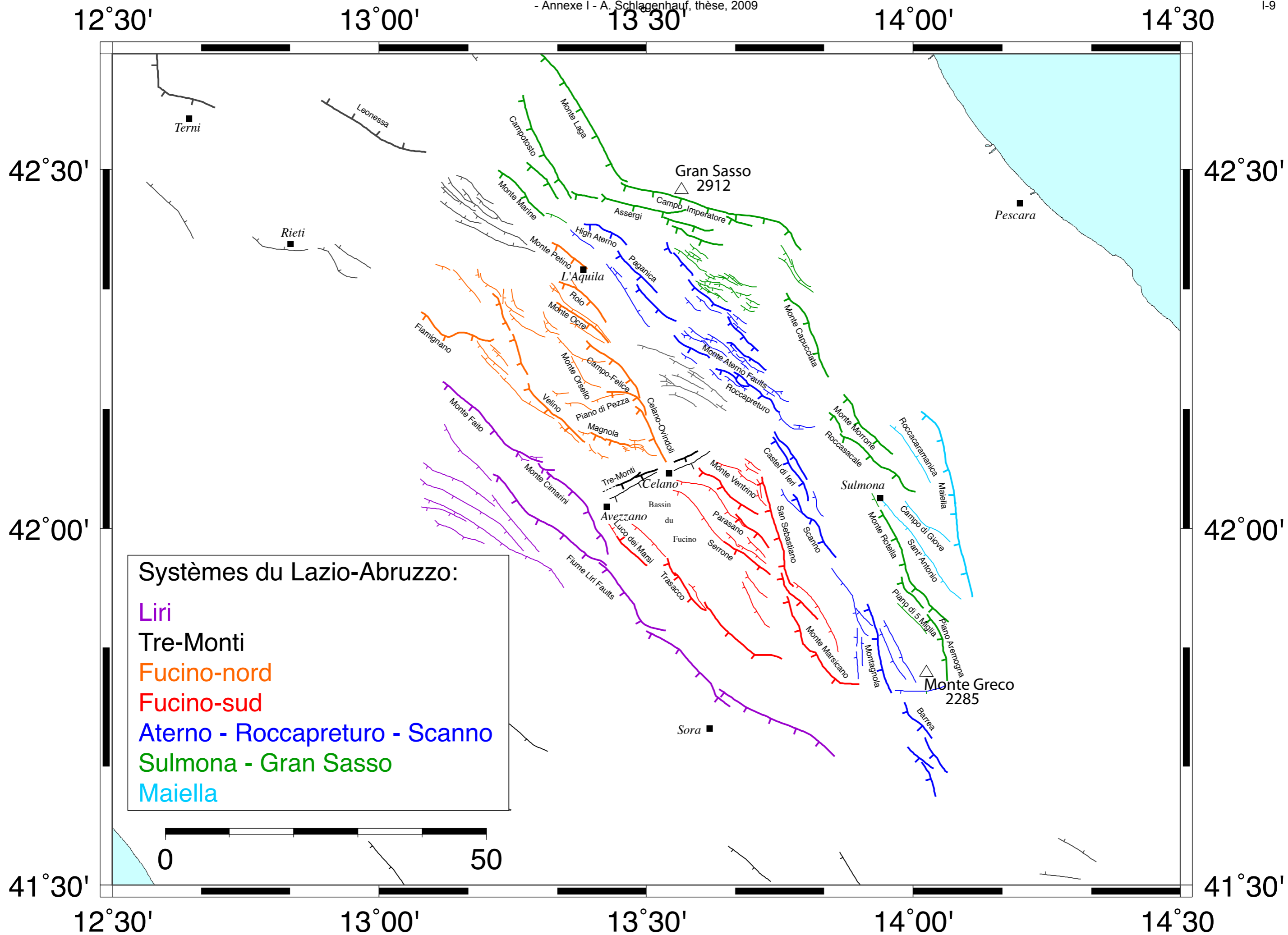


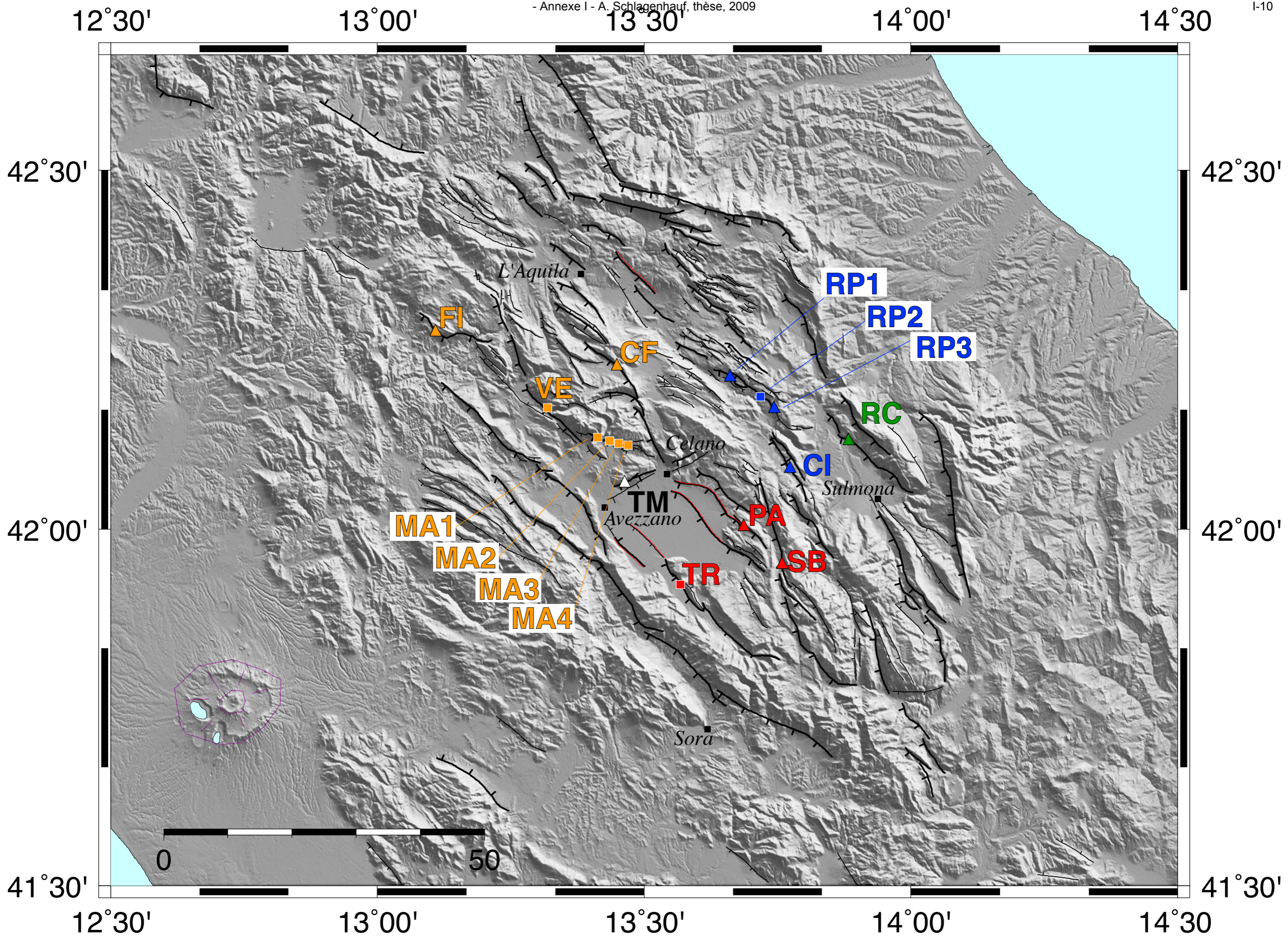














## **Annexe II**

(du chapitre 4)





## APPENDIX

### A Topographic shielding of normal fault sites

*Equations below are only valid for calculation of the fast neutron flux which is not energy dependent.*

The cosmic ray flux intensity arriving from the sky in the direction given by the inclination  $\theta$  (positive upwards) and the azimuth  $\phi$  is generally given by :

$$I(\theta, \phi) = I_o \sin^m \theta \quad (\text{A.1})$$

for  $\theta > 0$  and  $I(\theta, \phi) = 0$  for  $\theta < 0$ .  $I_o$  is the maximum intensity and  $m$  is an empirical constant, generally taken to be equal to 2.3 (e.g. Lal, 1958; Nishiizumi et al., 1989). If the sample exposed is not shielded, it will receive the maximum possible flux given by :

$$\Phi_{\max} = \int_0^{2\pi} \int_0^{\pi/2} I_o \sin^m \theta \cos \theta d\theta d\phi = \frac{2\pi I_o}{m+1} \quad (\text{A.2})$$

In the presence of a shielding topography, the cosmic ray flux  $\Phi$  will be reduced by a shielding factor  $S = \Phi/\Phi_{\max}$ . If the topography is described by a relationship between its vertical angle  $\theta_{\text{topo}}$  (as seen from the sampling site) and the azimuth  $\phi$ , the remaining flux will be :

$$\Phi = \int_0^{2\pi} \int_{\theta_{\text{topo}}(\phi)}^{\pi/2} I_o \sin^m \theta \cos \theta d\theta d\phi \quad (\text{A.3})$$

and the shielding factor  $S$  is given by :

$$S = \frac{m+1}{2\pi} \int_0^{2\pi} \int_{\theta_{\text{topo}}(\phi)}^{\pi/2} \sin^m \theta \cos \theta d\theta d\phi \quad (\text{A.4})$$

Equation A.4 can lead to a closed form in the rare cases where the  $\theta_{\text{topo}}(\phi)$  relationship is simple. Otherwise, it has to be numerically integrated.

#### A simple dipping fault plane

Generally, the fault scarp can be modelled in a first approximation as a semi-infinite plane of dip  $\beta$ . If the azimuth  $\phi$  is taken to be 0 (or  $\pi$ ) in the fault direction and  $\pi/2$  in the direction of maximum scarp slope, the relationship between  $\theta_{\text{topo}}$  and  $\phi$  is given by :

$$\theta_{\text{topo}}(\phi) = \text{atan}(\tan \beta \sin \phi), \quad 0 \leq \phi \leq \pi \quad (\text{A.5a})$$

$$\theta_{\text{topo}}(\phi) = 0, \quad \pi \leq \phi \leq 2\pi \quad (\text{A.5b})$$

Equation A.5a is simply the expression of the apparent slope of the scarp in the azimuth  $\phi$ . The resulting shielding factor  $S(\beta)$  is given by :

$$S(\beta) = \frac{1}{2} + \frac{m+1}{2\pi} \int_0^{\pi} \int_{\text{atan}(\tan \beta \sin \phi)}^{\pi/2} \sin^m \theta \cos \theta d\theta d\phi \quad (\text{A.6})$$

If there is a significant topography across the valley, its shielding effect must be included here, in the  $\pi \leq \phi \leq 2\pi$  domain, by replacing the 1/2 term by an expression similar to equation (A.4) but for  $\pi \leq \phi \leq 2\pi$ .

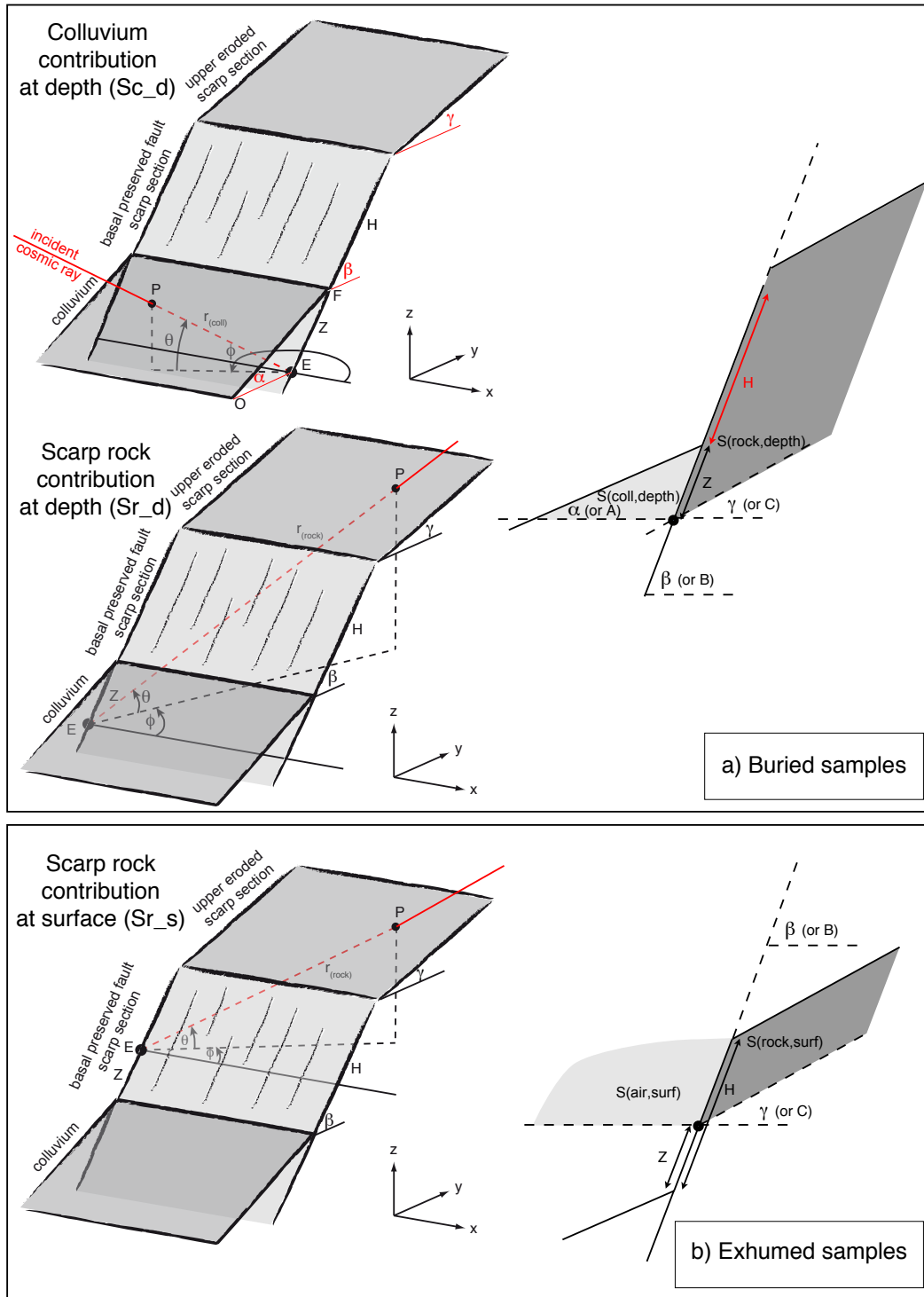


FIGURE A.1 - Schematic representation of the geometry of a normal fault scarp and associated colluvial wedge, and resulting shielding.  $\alpha$  is colluvial wedge dip,  $\beta$  the basal scarp dip, and  $\gamma$  the dip of the upper part of the fault escarpment. H is the actual height of the basal scarp. When  $\gamma$  and  $\beta$  are different, a) buried samples receive cosmic rays passing through the colluvium ('colluvium contribution' of the flux), and cosmic rays passing through the scarp rock ('rock contribution' of the flux), b) the exhumed samples receive cosmic rays traveling through the air, plus cosmic rays passing through the rock.

### Colluvial wedge

As the samples have been buried beneath a colluvial wedge before being exhumed by an earthquake, it is necessary to take into account the partial shielding created by this wedge. As for the fault scarp, we assume that it can be modelled by a semi-infinite plane of dip  $\alpha$ . At any point beneath the colluvial wedge, the cosmic ray flux intensity coming from the direction  $(\theta, \phi)$  will be attenuated from its surface value by a factor  $\exp(-d/\lambda)$  where  $d$  is the distance travelled through the colluvium by the incoming particle and  $\lambda$  the true attenuation length of the corresponding particle flux. In the case of fast neutrons,  $\lambda \approx 208 \text{ g.cm}^{-1}$  (e.g. Gosse and Phillips, 2001).

If the target is at a distance  $Z$  (positive downwards) beneath the colluvium-fault intersection (Figure 4 in the main text), the distance  $d$  travelled through the colluvium by a particle coming from the direction  $(\theta, \phi)$  is given by :

$$d_- = \frac{Z \sin(\beta - \alpha)}{\sin \theta \cos \alpha - \cos \theta \sin \alpha \sin \phi} \text{ if } 0 \leq \phi \leq \pi \quad (\text{A.7a})$$

$$d_+ = \frac{Z \sin(\beta - \alpha)}{\sin \theta \cos \alpha + \cos \theta \sin \alpha \sin \phi} \text{ if } \pi \leq \phi \leq 2\pi \quad (\text{A.7b})$$

Our expression of  $d$  differs from that given by Dunne et al. (1999, eq. 12) as we use a different depth  $Z$ . Combining the shielding provided by the fault scarp of slope  $\beta$  and the attenuation due to the colluvial wedge of slope  $\alpha$ , the resulting flux at  $Z$  is :

$$\begin{aligned} \Phi(Z, \alpha, \beta) = & \int_0^\pi \int_{\text{atan}(\tan \beta \sin \phi)}^{\pi/2} I_0 e^{-d_-/\lambda} \sin^m \theta \cos \theta \, d\theta \, d\phi \\ & + \int_\pi^{2\pi} \int_0^{\pi/2} I_0 e^{-d_+/\lambda} \sin^m \theta \cos \theta \, d\theta \, d\phi \end{aligned} \quad (\text{A.8})$$

In the actual calculation, the distances  $d_-$ ,  $d_+$ , and the attenuation length  $\lambda$  have to be expressed in the same units, e.g. cm or  $\text{g.cm}^{-2}$ . The ratio between  $\Phi$  and  $\Phi_{\text{max}}$  can then be used as a scaling factor  $S_d(Z, \alpha, \beta)$  for the production rate (Dunne et al. 1999) beneath the colluvial wedge. The Matlab<sup>®</sup> program (see Electronic Supplement 1) `scdepth.m` computes  $S_d(Z, \alpha, \beta)$ .

Plotting  $S$  as a function of  $Z$  for given values of  $\alpha$  and  $\beta$  shows (Figures A.2.a,b,c) that it is reasonable to approximate the scaling factor by an exponential decay of the form :

$$S_d(Z, \alpha, \beta) \approx S_Z \exp(-Z/\Lambda_Z) \quad (\text{A.9})$$

where  $\Lambda_Z$  is an apparent attenuation length. In order to facilitate the comparison with the work of Dunne et al. (1999), we calculate the scaling factor as a function of  $z$ , the distance measured perpendicularly to the colluvial upper surface, which is related to  $Z$  by  $z = Z \sin(\beta - \alpha)$ . The scaling factor can also be approximated by an exponential decay in  $z$  of the form  $S_d(z) = S_z \exp(-z/\Lambda_z)$  where  $\Lambda_z$  is the apparent attenuation length relevant to the  $z$  direction, perpendicular to the colluvial outer surface. In the case of a flat colluvium and no scarp,  $\alpha = \beta = 0$ , we find that setting the true attenuation length  $\lambda = 208 \text{ g.cm}^{-2}$  gives  $\Lambda_z \approx 157 \text{ g.cm}^{-2}$ , which is close to the classical value of  $160 \text{ g.cm}^{-2}$  used in vertical profiles (e.g. Gosse and Phillips, 2001). Other calculated values of  $s_{o,f}$  and  $\Lambda_f$  are given in Figures A.2c,d for given  $\alpha$  and  $\beta$  related to Figures 5a,b of the main text.

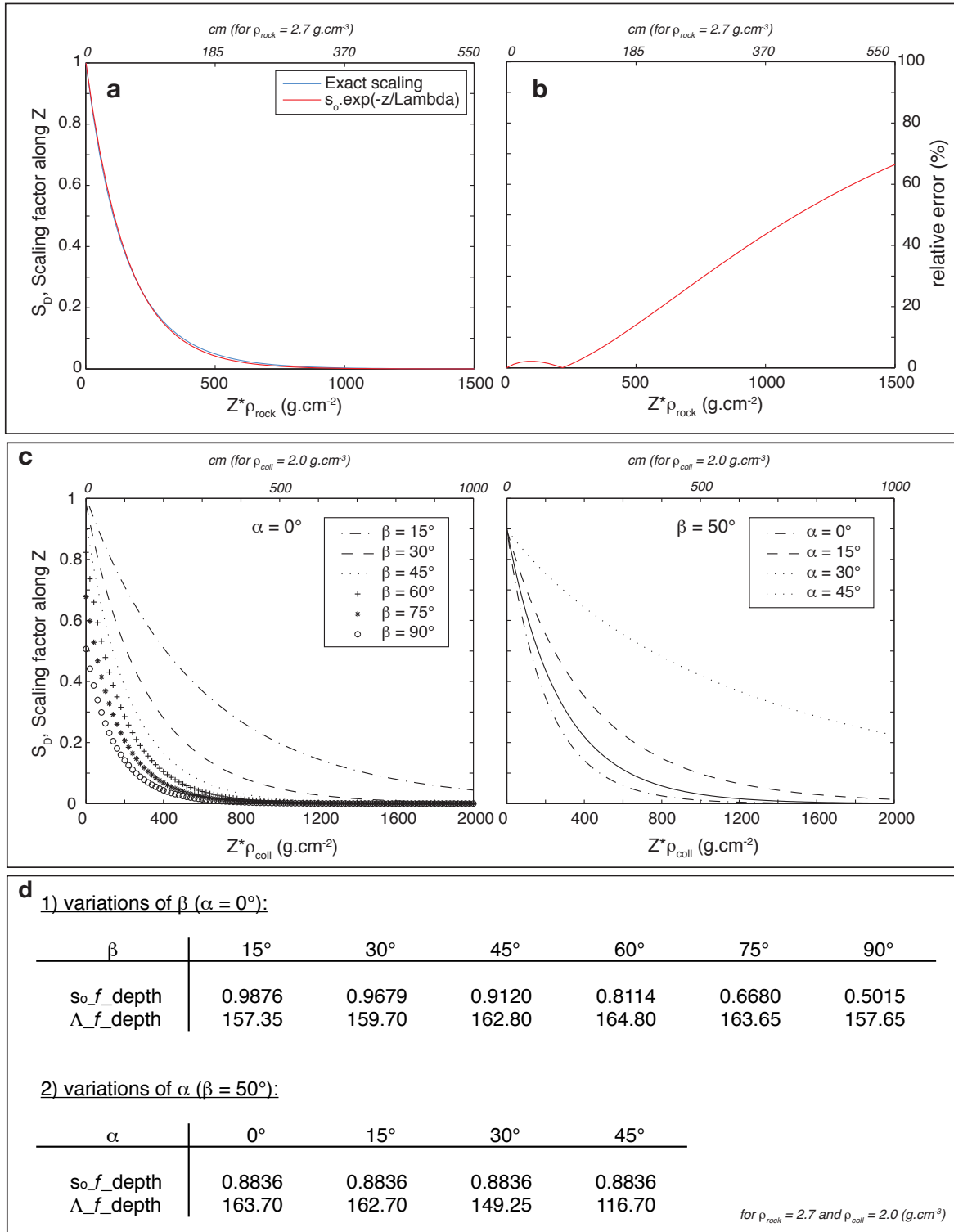


FIGURE A.2 - Fast neutron attenuation at depth. a) Comparison of exact scaling and its fit by an exponential with the fitexp.m Matlab<sup>®</sup> code ( $S = s_o \cdot \exp(-Z/\Lambda_f)$ ). Corresponding values of  $z$  in cm at the top are given for  $\rho_{rock} = 2.7$ . b) Relative error between the two previous curves. c) Scaling with depth for different values of  $\beta$  and  $\alpha$ , and d) corresponding values of  $s_o$  and  $\Lambda_f$ .

### Change in dip of the fault escarpment

It has been observed that, above the steep scarplet, the mountain front has a gentler slope. Samples located near the top of the scarplet will receive not only particles from the atmosphere directly above them but also particles who will have travelled through the upper surface (Figure A.1). Again, if we take the scarp and the mountain front above as planar surfaces of slopes  $\beta$  and  $\gamma$  respectively, the particle flux arriving at the sample from the direction  $(\theta, \phi)$  will be attenuated by a factor  $\exp(-d/\lambda)$  where  $d$  is the distance travelled through the footwall rock from the entry point to the sample and  $\lambda$  is the true attenuation length. If the scarp height is  $H$  and the position of the sample on the scarp surface is measured by  $Z$  (positive upwards), then  $d$  is given by :

$$d = \frac{(H - Z) \sin(\beta - \gamma)}{\sin \theta \cos \gamma - \cos \theta \sin \gamma \sin \phi} \quad (\text{A.10})$$

The net contribution of these particles to the flux reaching the sample will then be :

$$\Phi = \int_0^\pi \int_0^{\text{atan}(\tan \beta \sin \phi)} I_0 e^{-d/\lambda} \sin^m \theta \cos \theta \, d\theta \, d\phi \quad (\text{A.11})$$

In fact, the lower limit of the inner integral may not be 0 but  $\text{atan}(\tan \gamma \sin \phi)$ , the apparent slope of the upper surface in the direction  $\phi$ , as the distance  $d$  becomes infinite for  $\theta \leq \text{atan}(\tan \gamma \sin \phi)$  and attenuation is complete.

The total flux received by the sample can then be calculated and leads to the following scaling factor  $S_s(Z, H, \beta, \gamma)$  :

$$\begin{aligned} S_s(Z, H, \beta, \gamma) &= \frac{1}{2} + \frac{m+1}{2\pi} \int_0^\pi \int_{\text{atan}(\tan \beta \sin \phi)}^{\pi/2} \sin^m \theta \cos \theta \, d\theta \, d\phi \\ &+ \frac{m+1}{2\pi} \int_0^\pi \int_{\text{atan}(\tan \gamma \sin \phi)}^{\text{atan}(\tan \beta \sin \phi)} e^{-d/\lambda} \sin^m \theta \cos \theta \, d\theta \, d\phi \end{aligned} \quad (\text{A.12})$$

The Matlab<sup>®</sup> program (see Electronic Supplement 1) `scurf.m` computes  $S_s(Z, H, \beta, \gamma)$ . This time, the scaling factor  $S_s(Z, H, \beta, \gamma)$  should be modelled, for given values of  $H$ ,  $\beta$ , and  $\gamma$ , as the sum of a constant term and of an exponential, reflecting the contributions of particles reaching directly the sample and particles having travelled through the upper surface of the fault scarp, respectively.

If the scarplet height  $H$  is small enough, it may become necessary to combine the effects of the colluvial wedge and of the upper surface to calculate the scaling factor for samples located at relatively shallow depths beneath the colluvial wedge (in fact, the Matlab<sup>®</sup> program `scdepth.m` mentioned above includes the presence of an upper slope).

### In the presence of erosion

In the presence of erosion, a sample now at the surface will have been partially shielded from the surface flux by a layer of rock of decreasing thickness. It is therefore necessary to calculate the flux reaching a sample beneath the scarp surface. Assuming for the moment that the scarp surface is an infinite plane of slope  $\beta$ , if  $e$  is the rock thickness between the (future) sample and the scarp surface (Figure A.3a), the distance travelled by a particle coming from the direction  $(\theta, \phi)$  is :

$$d_- = \frac{e}{\sin \theta \cos \beta - \cos \theta \sin \beta \sin \phi} \quad \text{if } 0 \leq \phi \leq \pi \quad (\text{A.13a})$$

$$d_+ = \frac{e}{\sin \theta \cos \beta + \cos \theta \sin \beta \sin \phi} \quad \text{if } \pi \leq \phi \leq 2\pi \quad (\text{A.13b})$$

The corresponding scaling factor is :

$$S_r(e, \beta) = \frac{m+1}{2\pi} \int_0^\pi \int_{\text{atan}(\tan \beta \sin \phi)}^{\pi/2} e^{-d_-/\lambda} \sin^m \theta \cos \theta \, d\theta \, d\phi \\ + \frac{m+1}{2\pi} \int_\pi^{2\pi} \int_0^{\pi/2} e^{-d_+/\lambda} \sin^m \theta \cos \theta \, d\theta \, d\phi \quad (\text{A.14})$$

The Matlab<sup>®</sup> program (see Electronic Supplement 1) `scrock.m` computes  $S_r(e, \beta)$ . For a given scarp slope  $\beta$ , the scaling factor  $S_r(e)$  is also well approximated by an exponential decay of the form :

$$S_r(e, \beta) \approx S_e \exp(-e/\Lambda_e) \quad (\text{A.15})$$

### Colluvium + upper surface + erosion

If we wish to model the erosion of a scarp of slope  $\beta$  and height  $H$ , above a colluvial wedge of slope  $\alpha$  and below an upper surface of slope  $\gamma$ , the same methodology as in the above cases should be used : (1) calculating first the distance  $d$  between the entry point at the Earth's surface and the target at depth as a function of the incoming direction  $(\theta, \phi)$ , (2) integrating the attenuated flux intensity over the  $(0 \leq \phi \leq 2\pi, 0 \leq \theta \leq \pi/2)$  domain, skipping regions from which no flux is coming by adjusting the integration limits for  $\theta$  as we did above.

It is easy to see that the expression of  $d$  as a function of the target position given by  $Z$  and  $e$  (Figure A.3a) becomes very complex as several regions have to be considered and that the integration of the flux intensity becomes cumbersome, if not painful. It will be very difficult to easily express the corresponding scaling factor  $S_t(Z, e)$  as a function of  $Z$  and  $e$ . As, however, this scaling factor should vary smoothly with  $Z$  and  $e$ , we believe that it is reasonable to assume that  $S_t$  can be expressed as the product of two scaling factors, one accounting for the shape of the scarp itself and being a function of  $Z$ , and one accounting for the thickness  $e$  of footwall rock separating the target from the scarp surface :

$$S_t(Z, e) = f(Z)g(e) \quad (\text{A.16})$$

At the surface of the scarp ( $e = 0$ ),  $S_t(Z, 0)$  should be equal to the scaling factors  $S_d$  or  $S_s$  we have calculated above in equations (A.9) and (A.12). Above the colluvial wedge and far enough from the top of the scarp surface,  $S_t(Z, e)$  should be equal to  $S_r(e)$ . The solution is to take :

$$S(Z, e) = S_r(e)\bar{S}(Z) \quad (\text{A.17})$$

where  $\bar{S}(Z)$  is either  $S_d(Z)$  or  $S_s(Z)$ , depending on the location of the target, normalized to the (identical) value they take at the colluvium-fault scarp limit (Figure A3.a). This ensures that the effect of the scarp slope  $\beta$  is not taken into account twice. With this normalization,  $\bar{S}(Z)$  decrease exponentially from 1 towards 0 for increasing depths beneath the colluvium, while it increases slightly towards the top of the scarp, reflecting the increasing contribution of particles entering through the upper surface (Figure A.3b).

## B Matlab<sup>®</sup> code description

### The production rate $P$

Production rates are calculated using the full description of the production sources from Schimmelpfennig et al. (2009; after Gosse and Phillips, 2001). Both `clrock.m` and `clcoll.m` have equivalent equations. They differ only by the fact that `clcoll.m` also takes into account the colluvial wedge mean chemical composition : it calculates the production rate in the sample (scarp rock composition) buried under a thickness  $z$  of colluvium having its proper chemical composition.

Thickness factors ( $Q_x$ ; with  $x$  standing for each production source) from Gosse and Phillips (2001) are only valid for samples at the surface vicinity. We re-calculated the thickness factors to express them as a function of depth. The new thickness factors are defined at the sample center, or thickness half, named

$th_2$  (in 'e' direction) and are the following (using Schimmelpfennig et al., 2009 nomenclature) :

For spallation :

$$Q_{sp} = 1 + \frac{1}{6} \left( \frac{th_2}{\Lambda_{f,e}} \right)^2 \quad (\text{B.1})$$

For direct capture of slow negative muons :

$$Q_{\mu} = 1 + \frac{1}{6} \left( \frac{th_2}{\Lambda_{\mu}} \right)^2 \quad (\text{B.2})$$

For epithermal neutrons :

$$\begin{aligned} Q_{eth} = \frac{1}{P_{eth}} & \left[ \phi_{eth}^* \left( \frac{f_{eth}}{\Lambda_{eth}} \right) (1 - p(E_{th})) \exp \left( \frac{-e}{\Lambda_{f,e}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{\Lambda_{f,e}} \right)^2 \right) \right. \\ & + (1 + R_{\mu} R_{eth}) (F \Delta \phi)_{eth}^* \left( \frac{f_{eth}}{\Lambda_{eth}} \right) (1 - p(E_{th})) \exp \left( \frac{-e}{L_{eth}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{L_{eth}} \right)^2 \right) \\ & \left. + R_{\mu} \phi_{eth}^* \left( \frac{f_{eth}}{\Lambda_{eth}} \right) (1 - p(E_{th})) \exp \left( \frac{-e}{\Lambda_{\mu}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{\Lambda_{\mu}} \right)^2 \right) \right] \quad (\text{B.3}) \end{aligned}$$

For thermal neutrons :

$$\begin{aligned} Q_{th} = \frac{1}{P_{th}} & \left[ \phi_{th}^* \left( \frac{f_{th}}{\Lambda_{th}} \right) \exp \left( \frac{-e}{\Lambda_{f,e}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{\Lambda_{f,e}} \right)^2 \right) \right. \\ & + (1 + R'_{\mu}) (\mathfrak{S} \Delta \phi)_{eth}^* \left( \frac{f_{th}}{\Lambda_{th}} \right) \exp \left( \frac{-e}{L_{eth}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{L_{eth}} \right)^2 \right) \\ & + (1 + R'_{\mu} R_{th}) (\mathfrak{S} \Delta \phi)_{th}^* \left( \frac{f_{th}}{\Lambda_{th}} \right) \exp \left( \frac{-e}{L_{th}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{L_{th}} \right)^2 \right) \\ & \left. + R'_{\mu} \phi_{th}^* \left( \frac{f_{th}}{\Lambda_{th}} \right) \exp \left( \frac{-e}{\Lambda_{\mu}} \right) \left( 1 + \frac{1}{6} \left( \frac{th_2}{\Lambda_{\mu}} \right)^2 \right) \right] \quad (\text{B.4}) \end{aligned}$$

Apparent attenuation length for fast neutrons is calculated as explained above in part 1 of the Appendix using either `scdepth.m`, `scurf.m`, or `scrock.m` depending of sample position in  $Z$  direction (Figure 5 of the main text). As slow muons flux is energy dependent (e.g. Gosse and Phillips, 2001) the same calculation of the apparent attenuation length of slow muons cannot be achieved as for fast neutron. We use the  $1500 \text{ g.cm}^{-2}$  defined for a flat surface (e.g. Gosse and Phillips, 2001, p 1504).

### The number of atoms $N$

The equation governing the evolution of the number  $N$  of  $^{36}\text{Cl}$  atoms is :

$$\frac{dN}{dt} = -\lambda_{36}N + P \quad (\text{B.5})$$

where  $\lambda_{36}$  is the decay constant of  $^{36}\text{Cl}$  and  $P$  is the production rate. In certain conditions, equation (B.5) can be solved analytically. Here we want to model the evolution of a fault scarp whose geometry changes repeatedly with time, we wish also to study the influence of a variable magnetic field on the production rate, and to account for possible erosion of the scarp surface. This has lead us to solve equation (B.5) numerically by dividing the time span of the scarp evolution into small increments of time and writing :

$$N(t + \Delta t) = N(t) + [P(t) - \lambda_{36}N(t)] \times \Delta t \quad (\text{B.6})$$

These time increments are small enough, typically 100 years, to ensure that the magnetic field and hence its effects on the production rate through the elevation-latitude correcting factors  $S_{el,f}$  for fast neutrons and  $S_{el,\mu}$  for muons, can be considered constant. In that, we follow the method used for instance by Dunai (2001), Pigati and Lifton (2004), Lifton et al. (2005) and Lifton et al. (2008). The changes in geometry, either by earthquake slip or erosion are accounted for by calculating the scaling factor  $S(Z(t), e(t))$ . For instance, if the erosion rate is  $\epsilon$ , the thickness  $e$  of rock between the sample and the scarp surface varies as :

$$e(t + \Delta t) = e(t) - \epsilon \times \Delta t \tag{B.7}$$

Changes in  $Z$  occur every time there is an earthquake, so that  $Z(t)$  values are updated much less frequently than  $e(t)$  values. Finally, each segment of the fault scarp has its own exhumation history. This has led us to write a Matlab<sup>®</sup> program, named `modelscarp.m` (see Electronic Supplement 1 and Figure B.1), whose main structure is based on the nesting of three loops : one for each of the segments, one for the number of earthquakes prior to the exhumation of the current segment (once it is exhumed, the actual  $Z$  values of a segment have no influence the production rate), and one for the small time increments  $\Delta t$ . Before these loops, the program calculates the number of  $^{36}\text{Cl}$  atoms at the end of the pre-exposition time.

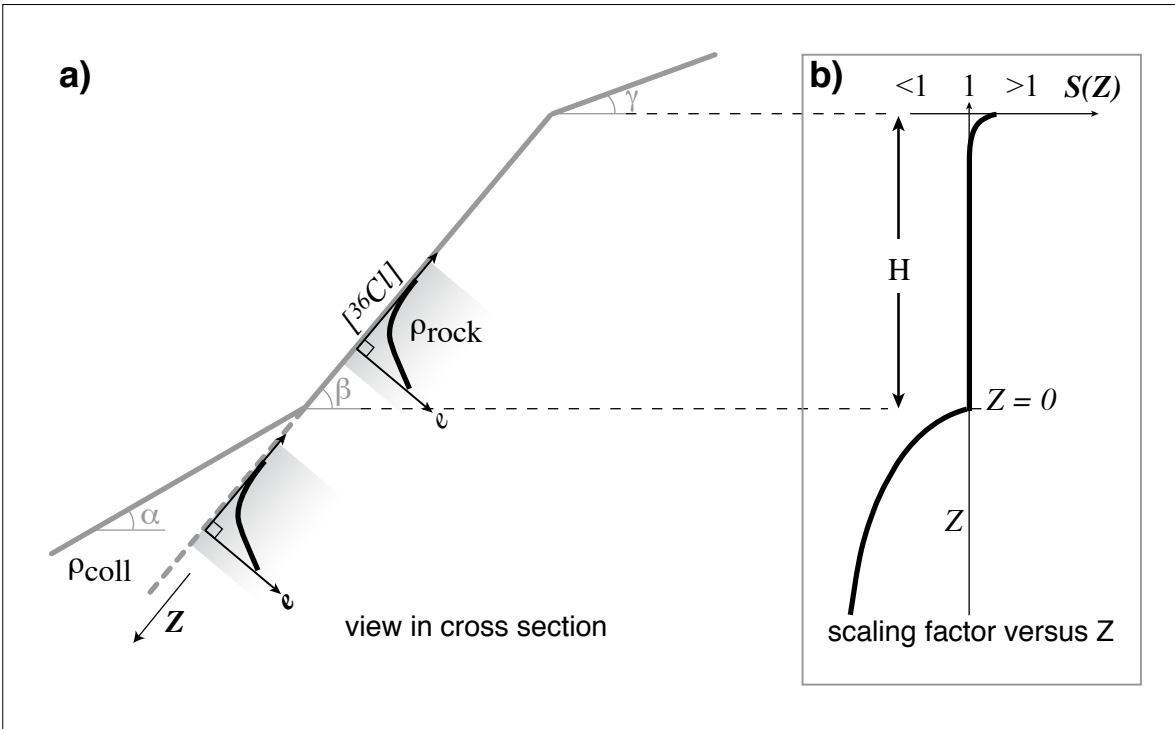
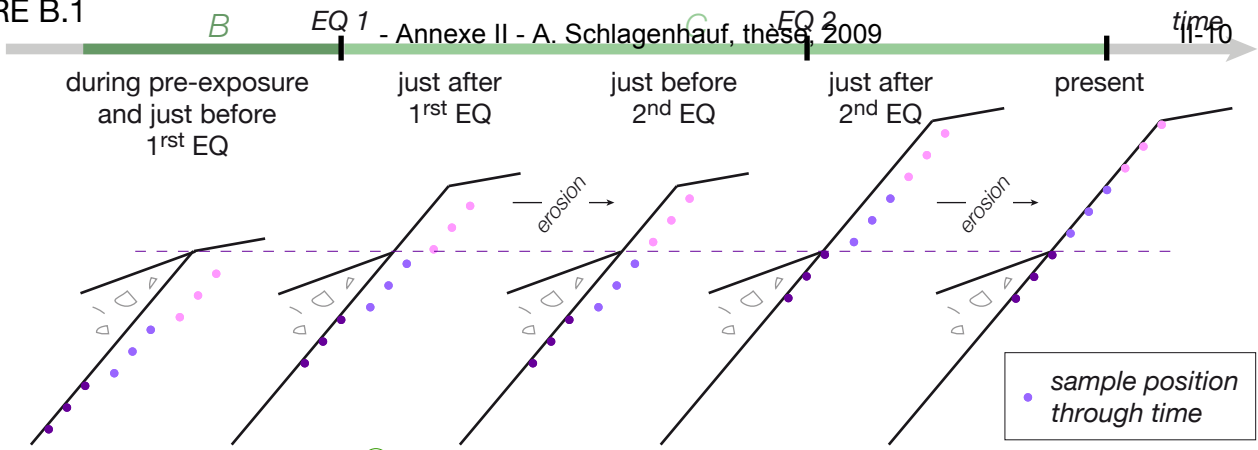


FIGURE A.3 - The attenuations along 'e' and 'Z' directions are separated : a) the  $^{36}\text{Cl}$  profile due to rock attenuation is calculated for a surface of dip  $\beta$  in the direction  $e$  perpendicular, b) a scaling factor  $S(Z)$ , which is only applied to the cosmogenic part of the production, is introduced. It accounts for both attenuation at depth ( $S(Z) < 1$ ) and increased flux at the vicinity of  $\gamma$  ( $S(Z) > 1$ ). It is calculated by scaling the cosmogenic production at surface far from  $\gamma$  (at  $Z = 0$ , for  $\beta$  infinite).

FIGURE B.1 - Schematic representation of the Matlab<sup>®</sup> `modelscarp.m` program given in Electronic Supplement 1. In orange, parameters that need to be modified for each site by the users.  $EQ$  stands for Earthquake;  $f$  for neutrons;  $\mu$  for muons;  $j$  for the sample;  $i$  for the earthquake number;  $t$  for the time step;  $rock$  for the sample chemical composition;  $coll$  for the mean colluvial wedge composition;  $P_{cosmo}$  for the cosmogenic production;  $P_{rad}$  for the radiogenic production;  $\epsilon$  for the erosion rate (in mm/yr).



FIGURE B.1



**modelscarp.m (Matlab<sup>®</sup> code):**

**A - Initialization**

CONSTANTS DECLARATION:  $\alpha, \beta, \gamma, \rho_{rock}, \rho_{coll}, H_{final}$  (per site)  
 $\Psi^{36}Cl\_Ca,0$  (function of Earth magnetic field description used)

CONSTANTS LOADING:  $\lambda_{36}, \Lambda_f, S_{el,f}(t), S_{el,\mu}(t)$

SURFACE SCALING (along  $z \geq 0$ )  $S(z \geq 0)$  equal to 1 if  $\beta$  infinite ;  $S(z) \geq 1$  (for  $z \geq 0$ ) if  $H, \gamma$   
[scsurf.m](#)

DEPTH SCALING (along  $z \leq 0$ ) after each EQ(i)  $S(z, H_i) \leq 1$  (for  $z \leq 0$ ) ; loop on  $H_i$   
 $S(z \leq 0, H_i)$  fitted with [fitexp.m](#)  $\rightarrow$   $sof_{d,i}$  and  $\Lambda_{f,d,i}$  (d: depth)  
[scdepth.m](#)

ROCK SCALING (along  $e \leq 0$ ) calculates  $S_e(e)$  (for  $e \leq 0$ ) independently of position on  $z$   
 $S_e(e)$  fitted with [fitexp.m](#)  $\rightarrow$   $sof_{f,e}$  and  $\Lambda_{f,e}$   
[scrock.m](#)

VARIABLES INITIALIZATION Initial samples (j) positions along  $e \rightarrow e_0(j)$

**B - Calculation of <sup>36</sup>Cl depth profile after pre-exposure (all samples buried under the colluvium)**

sample (j) time (t)  $\left[ \begin{array}{l} \text{clock.m ( rock(j), } e_0(j), \Lambda_{f,e}, sof_{f,e}, S_{el,f}(t), S_{el,\mu}(t) ) \rightarrow P_{cosmo(j)} \text{ and } P_{rad(j)} \text{ at } e_0(j) \\ \text{ccoll.m( coll(j), rock(j), } \Lambda_{f,d,i=0}, sof_{f,d,i=0}, S_{el,f}(t), S_{el,\mu}(t) ) / \text{ccoll.m (... , z=0,...) } \rightarrow S_{coll,cosmo(j)} \\ P_{tot,pre-exp(j)} = P_{rad(j)} + P_{cosmo(j)} * S_{coll,cosmo(j)} - \lambda_{36} * P_{tot,pre-exp(j,t-1)} * t \end{array} \right. \rightarrow P_{tot,pre-exp(j)} \text{ depth}$

**C - Calculation of <sup>36</sup>Cl depth and surface profiles during seismic phase**

earthquake (i) sample (j) time (t)  $\left[ \begin{array}{l} \text{clock.m ( rock(j), } e_0(j), \Lambda_{f,e}, sof_{f,e}, S_{el,f}(t), S_{el,\mu}(t) ) \rightarrow P_{cosmo(j)} \text{ and } P_{rad(j)} \text{ at } e(j) \\ \text{ccoll.m( coll(j), rock(j), } \Lambda_{f,d,i=0}, sof_{f,d,i=0}, S_{el,f}(t), S_{el,\mu}(t) ) / \text{ccoll.m (... , z=0,...) } \rightarrow S_{coll,cosmo(j)} \\ P_{tot,depth(j)} = P_{rad(j)} + P_{cosmo(j)} * S_{coll,cosmo(j)} - \lambda_{36} * P_{tot(j,t-1)} * t \end{array} \right. \rightarrow P_{tot,depth(j)} \text{ depth}$

sample (j) time (t)  $\left[ \begin{array}{l} \text{clock.m ( rock(j), } e(j,t), \Lambda_{f,e}, sof_{f,e}, S_{el,f}(t), S_{el,\mu}(t) ) \rightarrow P_{cosmo(j)} \text{ and } P_{rad(j)} \text{ at } e(j,t) \\ \text{(scsurf.m) } S(j) / S(z=0) \rightarrow S_{surf,cosmo(j)} \text{ (scaling)} \\ P_{tot,depth(j)} = P_{rad(j)} + P_{cosmo(j)} * S_{surf,cosmo(j)} - \lambda_{36} * P_{tot(j,t-1)} * t \\ e(j,t+1) = e(j,t) - \epsilon * t \end{array} \right. \rightarrow P_{tot,EQ(j)} \text{ surface}$

$$P(j) = P_{tot,pre-exp(j)} + P_{tot,EQ(j)}$$

**Matlab<sup>®</sup> command window:**

```
data = load('datarock.txt');           66 columns * x lines (x samples)
coll = load('datacolluvium.txt');      62 columns * 1 line (mean colluvium composition)
EL = load('datamagfield.txt');        4 columns * t lines
dataset loading

modelscarp(data, coll, age, slip, preexp, EL, epsilon)
calculation
```



## Legend of Electronic Supplement 1 – Schlagenhauf et al., 2009

**For a new sampling site:**

1- Create the 3 input files as the ones given for Magnola MA3 site (text format with tabulations – no coma, use points for decimal separator):

- ‘datarock.txt’: 66 columns by x lines (x : number of samples) ; see Tables 1 and 2 of main text
- ‘datacolluvium.txt’: 62 columns by 1 line (mean colluvium composition) ; see Table 1 of main text
- ‘datamagfield.txt’: 4 columns by t lines (variations of scaling factors  $S_{el,s}$  and  $S_{el,u}$ );  
from 1rst to 4th column: time period (yrs) – time steps (yrs) -  $S_{el,s}$  -  $S_{el,u}$

*example (for a variable magnetic field):*

0	0	2.744	1.784
100	100	2.706	1.78
200	100	2.763	1.786
300	100	2.825	1.792
400	100	2.777	1.787
500	100	2.692	1.778
600	100	2.302	1.732
700	100	2.102	1.697
800	100	2.33	1.736
900	100	2.482	1.756
1000	100	2.722	1.792

2- Modify the Matlab® code *modelscarp.m* (actual values are the one corresponding to MA3 Magnola site):

- alpha (degrees): colluvium wedge dip
- beta (degrees): preserved scarp dip
- gamma (degrees): upper eroded scarp dip
- Hfinal: final height (present height) of the fault scarp of dip  $\beta$  (m)
- rho\_coll: colluvial wedge mean density
- rho\_rock: scarp rock mean density
- Psi\_36Cl\_Ca\_0: spallation production rate of  $^{36}\text{Cl}$  at surface from  $^{40}\text{Ca}$  corresponding to the description of magnetic field used (ex: 48.8 at  $^{36}\text{Cl}/\text{g Ca /yr}$ , from Stone et al., 1996 using a constant magnetic field).

3- In the Matlab® command window:

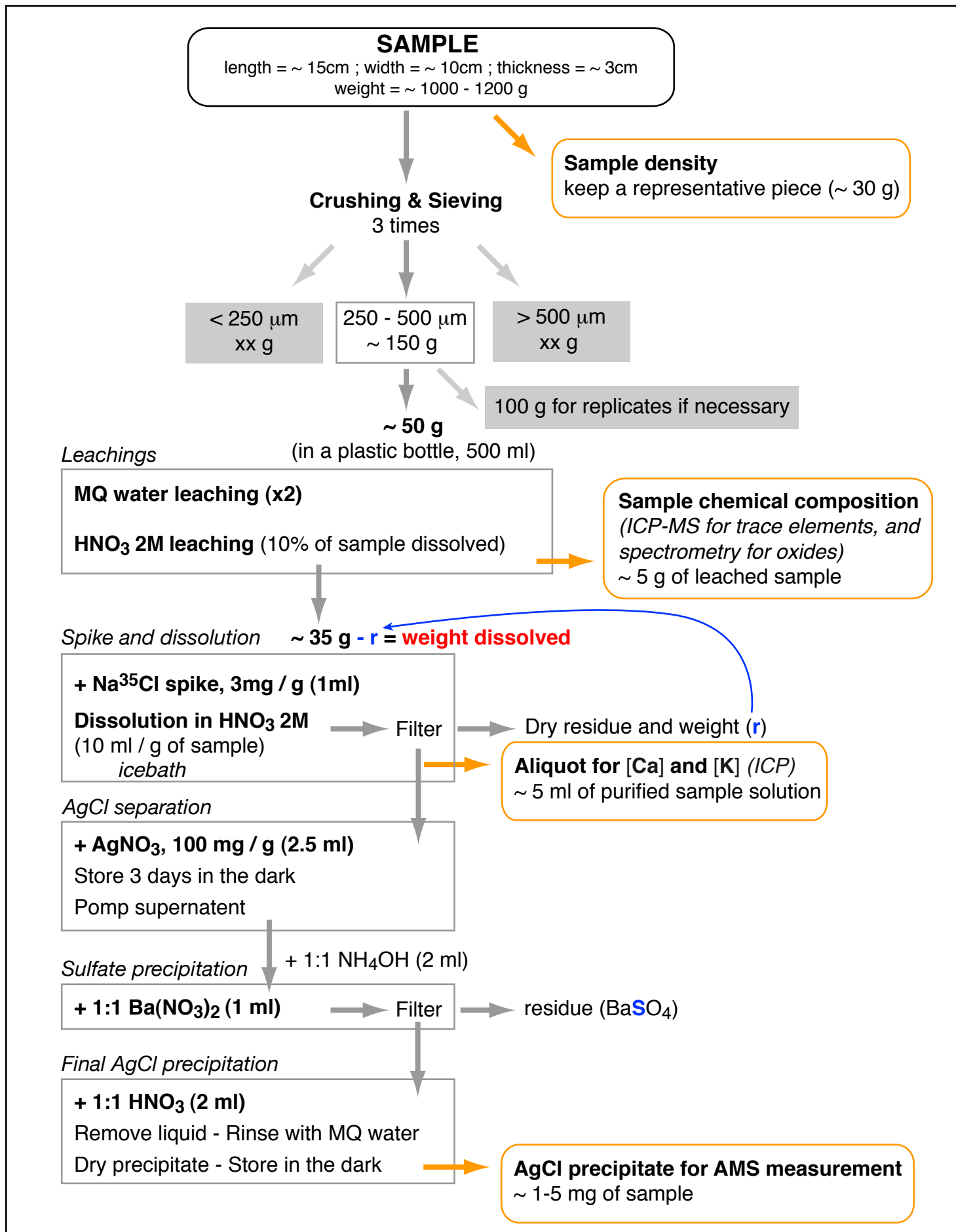
```
data = load('datarockMA3.txt');
coll = load('datacolluviumMA3.txt');
EL = load('datamagfieldMA3Stone2000.txt');
age = [7200 4900 4000 3400 1500 0]; % Earthquakes ages (years), here last event is at t=0 to model the buried samples
slip = [190 205 160 360 200 400]; % Corresponding earthquakes displacements (cm ; 400 cm for the buried samples)
preexp = 1900; % pre-exposure duration (years)
epsilon = 0; % erosion rate (mm/yr)
modelscarp(data,coll,age,slip,preexp,EL,epsilon)
```

**Output:** it can take a few minutes to run. A Figure window of Matlab will appear displaying the results. Black dots are the [ $^{36}\text{Cl}$ ] concentrations with horizontal bars for the corresponding uncertainties. Blue open dots are the modelled [ $^{36}\text{Cl}$ ] concentrations. Blue line represent the position of each discontinuities associated with each modelled event.

“Modelscarp” folder contains: all the Matlab® codes for earthquake recovering using  $^{36}\text{Cl}$  cosmogenic isotope on active normal fault scarps, and all input files for MA3 site on the Magnola Fault.

*modelscarp.m* is the main routine iterating the other six (see Appendix for more details).

- *clcoll.m* and *clock.m* calculate production of  $^{36}\text{Cl}$  in a sample buried under colluvium or at surface, respectively,
- *scdepth.m*, *scsurf.m* calculate the scaling of the samples along Z when buried or at surface, respectively,
- *scrock.m* calculates the scaling along direction e,
- *fitexp.m* calculates the fit by an exponential form of the 3 scaling factors described above to derived the attenuation length  $\Lambda$ .



**Sample chemical preparation for Cl extraction in limestone rocks**  
(modified after Stone et al., 1996)

## Electronic Supplement 3 – Schlagenhauf et al. 2009

*Equations for [Cl] and [<sup>36</sup>Cl] concentrations calculations, from AMS ratios.  
The example is given for a <sup>35</sup>Cl-spiked sample.*

*More details can be found in Desilets et al., 2006*

*[Desilets, D., Zreda, M., Almasi, P.F., Elmore, D., 2006. Determination of cosmogenic <sup>36</sup>Cl in rocks by isotope dilution: innovations, validation and error propagation. Chemical Geology 233, 185-195]*

$$\text{AMS ratios: } \left( \frac{{}^{36}\text{Cl}}{{}^{35}\text{Cl}(\text{or } {}^{37}\text{Cl})} \right)_{\text{sample,AMS}} \quad \text{and} \quad \left( \frac{{}^{36}\text{Cl}}{{}^{35}\text{Cl}(\text{or } {}^{37}\text{Cl})} \right)_{\text{blk.i,AMS}} ; \left( \frac{{}^{35}\text{Cl}}{{}^{37}\text{Cl}} \right)_{\text{sample,AMS}} \quad \text{and} \quad \left( \frac{{}^{35}\text{Cl}}{{}^{37}\text{Cl}} \right)_{\text{blk.i,AMS}}$$

### [Cl<sub>nat</sub>] calculation:

$$m_{\text{Cl,spike}} = [\text{Cl}_{\text{spike}}] \times m_{\text{spike}} \times 0.001 \quad (\text{g})$$

with  $[\text{Cl}_{\text{spike}}]$  the spike concentration in Cl (mg/g)  
 $m_{\text{spike}}$  the spike mass (g)

$$n_{\text{Cl,spike}} = m_{\text{Cl,spike}} \times N_{\text{av}} / M_{\text{spike}} \quad (\text{at.})$$

with  $N_{\text{av}}$  the Avogadro number ( $6.022 \cdot 10^{23} \text{ mol}^{-1}$ )  
 $M_{\text{spike}}$  the molar mass of the spike ( $\%_{35}\text{Cl} \times M_{35\text{Cl}} + \%_{37}\text{Cl} \times M_{37\text{Cl}}$ )  
 $M_{35\text{Cl}} = 34.9689 \text{ g}\cdot\text{mol}^{-1}$  and  $M_{37\text{Cl}} = 36.9659 \text{ g}\cdot\text{mol}^{-1}$

*Isotopic dilution equations to calculate  $n_{\text{Cl}_{\text{nat, sample}}}$  :*

$$\text{Cl}_{\text{nat}} \text{ ratio: } R_{\text{Cl}_{\text{nat}}} = \frac{\%_{35}\text{Cl}_{\text{nat}}}{\%_{37}\text{Cl}_{\text{nat}}} = \frac{75.77}{24.23} = 3.127$$

$$\text{Spike ratio: } R_{\text{spike}} = \frac{\%_{35}\text{Cl}_{\text{spike}}}{\%_{37}\text{Cl}_{\text{spike}}} \quad (\text{known})$$

$$\text{Sample ratio from AMS: } R_{\text{sample,AMS}} = \left( \frac{{}^{35}\text{Cl}}{{}^{37}\text{Cl}} \right)_{\text{sample,AMS}} = \frac{n_{35\text{Cl}_{\text{spike}}} + n_{35\text{Cl}_{\text{nat, sample}}}}{n_{37\text{Cl}_{\text{spike}}} + n_{37\text{Cl}_{\text{nat, sample}}}}$$

with  $n_{35\text{Cl}_{\text{spike}}} = n_{\text{Cl}_{\text{spike}}} \times \%_{35}\text{Cl}_{\text{spike}}$       and       $n_{37\text{Cl}_{\text{spike}}} = n_{\text{Cl}_{\text{spike}}} \times \%_{37}\text{Cl}_{\text{spike}}$   
 $n_{35\text{Cl}_{\text{nat, sample}}} = n_{\text{Cl}_{\text{nat, sample}}} \times \%_{35}\text{Cl}_{\text{nat}}$       and       $n_{37\text{Cl}_{\text{nat, sample}}} = n_{\text{Cl}_{\text{nat, sample}}} \times \%_{37}\text{Cl}_{\text{nat}}$

$$R_{\text{sample}} = \frac{(n_{\text{Cl}_{\text{spike}}} \times \%_{35}\text{Cl}_{\text{spike}}) + (n_{\text{Cl}_{\text{nat, sample}}} \times \%_{35}\text{Cl}_{\text{nat}})}{(n_{\text{Cl}_{\text{spike}}} \times \%_{37}\text{Cl}_{\text{spike}}) + (n_{\text{Cl}_{\text{nat, sample}}} \times \%_{37}\text{Cl}_{\text{nat}})}$$

$$\text{for the sample: } n_{\text{Cl}_{\text{nat, sample}}} = n_{\text{Cl}_{\text{spike}}} \times \left( \frac{\%_{35}\text{Cl}_{\text{spike}} - \left( \frac{{}^{35}\text{Cl}}{{}^{37}\text{Cl}} \right)_{\text{sample,AMS}} \times \%_{37}\text{Cl}_{\text{spike}}}{\left( \frac{{}^{35}\text{Cl}}{{}^{37}\text{Cl}} \right)_{\text{sample,AMS}} \times \%_{37}\text{Cl}_{\text{nat}} - \%_{35}\text{Cl}_{\text{nat}}} \right)$$

$$\text{and for the blank: } n_{Cl_{nat},blk,i} = n_{Cl_{spike}} \times \left( \frac{\left( \frac{35}{37} Cl \right)_{blk,i,AMS} \times \%_{37} Cl_{spike} - \%_{35} Cl_{spike}}{\left( \frac{35}{37} Cl \right)_{blk,i,AMS} \times \%_{37} Cl_{nat} - \%_{35} Cl_{nat}} \right)$$

So, the blank corrected sample number of mole is:  $n_{Cl_{nat},corr} = n_{Cl_{nat},sample} - n_{Cl_{nat},blk,i}$

$$\text{Finally, } [Cl_{nat}]_{sample} = \frac{n_{Cl_{nat},corr} \times M_{Cl_{nat}} \times 10^6}{N_{av} \times m_{sample}} \quad (\text{ppm})$$

### [<sup>36</sup>Cl] calculation:

$$\text{For a } ^{35}\text{Cl spike: } \left( \frac{36}{35} Cl \right)_{sample,AMS} = \frac{n_{36} Cl_{sample}}{n_{35} Cl_{sample}} \Leftrightarrow n_{36} Cl_{sample} = \left( \frac{36}{35} Cl \right)_{sample,AMS} \times n_{35} Cl_{sample}$$

$$\text{with } n_{35} Cl_{sample} = n_{35} Cl_{spike} + n_{35} Cl_{nat,sample} \Leftrightarrow n_{35} Cl_{sample} = \left( n_{Cl_{spike}} \times \%_{35} Cl_{spike} \right) + \left( n_{Cl_{nat,sample}} \times \%_{35} Cl_{nat} \right)$$

$$\text{for the sample: } n_{36} Cl_{sample} = \left( \frac{36}{35} Cl \right)_{sample,AMS} \times \left[ \left( n_{Cl_{spike}} \times \%_{35} Cl_{spike} \right) + \left( n_{Cl_{nat,sample}} \times \%_{35} Cl_{nat} \right) \right]$$

$$\text{and for the blank: } n_{36} Cl_{blk,i} = \left( \frac{36}{35} Cl \right)_{blk,i,AMS} \times \left[ \left( n_{Cl_{spike}} \times \%_{35} Cl_{spike} \right) + \left( n_{Cl_{nat,blk,i}} \times \%_{35} Cl_{nat} \right) \right]$$

So, the blank corrected number of mole in the sample is:  $n_{36} Cl_{sample,corr} = n_{36} Cl_{sample} - n_{36} Cl_{blk,i}$

$$\text{Finally, } [^{36}Cl] = \frac{n_{36} Cl_{sample,corr}}{m_{sample}} \quad (\text{at.g}^{-1})$$

And the uncertainty on the calculated [<sup>36</sup>Cl] concentration is:

$$\Delta [^{36}Cl] = [^{36}Cl] \times \sqrt{\left( \frac{\Delta^{36}Cl / ^{35}Cl}{^{36}Cl / ^{35}Cl} \right)^2 + \left( \frac{\Delta^{35}Cl / ^{37}Cl}{^{35}Cl / ^{37}Cl} \right)^2}$$

*blk,i*: blank associated to chemistry series *i*  
*sample*: sample prepared in the series *i*

*Cl<sub>nat</sub>*: natural chlorine (75.77% <sup>35</sup>Cl; 24.23% <sup>37</sup>Cl)  
*m<sub>sample</sub>*: sample weight

**Electronic Supplement 4 - Schlagenhauf et al. 2009**

*Scaling factors for fast neutrons ( $S_{el,s}$ ) and slow muons ( $S_{el,\mu}$ ) according to different Earth paleomagnetic models*

**a - Stone (2000)**

**Spallation rate at latitude  $\lambda$  and elevation  $z$**  (normalized to its value at 60° latitude and at standard sea level pressure) :

$$S_{el,s}(P) = a + b \exp\left(\frac{-P}{150}\right) + cP + dP^2 + eP^3 \quad (\text{a.1})$$

with  $a, b, c, d, e$  : scaling coefficients function of latitude from 0 to > 60° with a 10° step (see Stone, 2000), and  $P$  : Pressure (hPa) - ( from Lal., 1991) :

$$P(z) = P_s \exp\left[-\frac{g_o M}{R \xi} (\ln T_s - \ln(T_s - \xi z))\right] \quad (\text{a.2})$$

$P_s$  : Sea level pressure = 1013.25 hPa;       $M$  : Molar weight of air  
 $T_s$  : Sea level temperature = 288.15 K;       $g_o$  : acceleration of gravity  
 $\xi$  : adiabatic lapse rate = 0.0065 K.m<sup>-1</sup>;       $R$  : gaz constant

}  $\frac{g_o M}{R} = 0.03417 \text{ K.m}^{-1}$

**Slow muon capture at latitude  $\lambda$  and elevation  $z$  :**

$$S_{el,\mu}(P) = M_{\lambda,1013.25} \exp\left(\frac{1013.25 - P}{242}\right) \quad (\text{a.3})$$

with the attenuation length of muons in the air = 242,  
 and  $M_{\lambda,1013.25}$  : scaling coefficients for latitudes from 0 to > 60° with a 10° step (see Stone, 2000).

**b - Dunai (2001)**

**Spallation rate at latitude  $\lambda$ , longitude  $\theta$  and elevation  $z$  :**

$$S_{el,s} = \frac{\sum N_i}{i}; \text{ with } N_{(z,Rc,t)} = N_{(1030,Rc,f)} \exp\left(\frac{z(h)}{\Lambda_{(Rc)}}\right) \quad (\text{b.1})$$

$N_{(1030,Rc,f)}$  : neutron flux at sea level,  
 $z(h)$  ; difference in atmospheric depth,  
 $\Lambda_{(Rc)}$  : attenuation coefficient.

• with  $N_{1030,Rc} = Y + \frac{A}{\left[1 + \exp\left(\frac{-(Rc-X)}{B}\right)\right]^C}$  ; and A, B, C, X, Y scaling coefficients (see Dunai, 2001)

(b.2)

Depending on the age of the site,  $R_c$  is calculated following one of the two equations below :

★ > 10-20 ka (using Yang et al., 2000 and Guyodo and Valet, 1999 record of dipole moment ) :

$$R_c(t) = \frac{M(t)\mu_o c}{16\pi R^2} \quad (b.3)$$

$M(t)$  : dipole moment over time ( $\text{Am}^2$ ); 0-11ka : Yang et al., 2000; > 11 ka : Guyodo and Valet, 1999;  
 $\mu_o$  : permeability of free space =  $1.2566 \cdot 10^{-6}$   
 $c$  : speed of light =  $2.99 \cdot 10^8$  (m/s);  
 $R$  : Earth radius =  $6.37 \cdot 10^6$  (m)

★ 0 - 11 ka (using Ohno and Hamano record of dipole moment, 1993) :

$$R_c(t) = \frac{R}{4} \frac{H(t)c}{(1 + 0.25 \tan^2 I(t))^{3/2}} \quad (b.4)$$

$H(t)$  : horizontal field strength over time,  
 $I(t)$  : inclination of the site over time.

$$H(t) = \frac{M(t)\mu_o \sin \theta(t)}{4\pi R^3} \quad (b.5)$$

$M(t)$  : dipole moment over time (values from Ohno and Hamano, 1993),  
 $\theta(t)$  : paleocolatitude over time,  $\cos \theta(t) = \sin \lambda_s \sin \lambda_N(t) + \cos \lambda_s \cos \lambda_N(t) \cos(\theta_N(t) - \theta_s)$ ,  
 $\lambda_s$  : present geographic latitude of the site,  
 $\lambda_N(t)$  : paleolatitude of the North geomagnetic pole over time (values from Ohno and Hamano, 1993),  
 $\theta_s$  : present geographic longitude of the site,  
 $\theta_N(t)$  : paleolongitude of the North geomagnetic pole over time (values from Ohno and Hamano, 1993).

$$I(t) = I_{(REF,measured)}(t) - I_{(REF,expected)} + I_{(s,expected)} \quad (b.6)$$

$I_{(REF,measured)}(t)$  : paleoinclination at the reference site over time,  
 $\tan(I_{(REF,expected)}) = 2 \tan \lambda_{(REF)}$  : paleoinclination at reference site if field would be a pure axial dipole,  
 $I_{(s,expected)} = 2 \tan \lambda_{(s)}$  : paleoinclination of the site if field would be a pure axial dipole.

• with  $z_{(h)} = \frac{10}{g_o}(P_o - P_h)$  ; and  $P_o = 1013.25$  mbar ;  $P_h$  : atmospheric pressure (mbar) (b.7)

$$P_h = P_o \left[ 1 - \beta_o \frac{z}{T_o} \right]^{\left( \frac{g_o}{R_d \beta_o} \right)} \quad (b.8)$$

$\beta_o$  : temperature decrease with  $z = 0.0065$  K/m,  
 $z$  : altitude (m);  $T_o$  : Temperature at sea level = 288.15 K,  
 $g_o$  : acceleration due to gravity =  $9.80665$  m.s<sup>-2</sup>;  $R_d$  : gas constant = 287.05 (J/kg/K).

• with  $\Lambda(R_c) = y + \frac{a}{\left[ 1 + \exp \frac{-(R_c - x)}{b} \right]^c}$  ; and a, b, c, x, y scaling coefficients (see Dunai, 2001) (b.9)

**Slow muon capture at latitude  $\lambda$  and elevation  $z$  (time independent) :**

$$S_{el,\mu} = N_{(1030,\mu)} \exp \frac{z(h)}{\Lambda_\mu} \quad (b.10)$$



$z_{(h)}$  ; difference in atmospheric depth (see equation b.7),  
 $\Lambda_{\mu} = 247 \text{ (g.cm}^{-2}\text{)}$  : muon attenuation length in atmosphere,  
 $N_{(1013.25,\mu)}$  : neutron flux at sea level.

$$\text{with } N_{1013.25,\mu} = Y + \frac{A}{\left[1 + \exp \frac{-(I-X)}{B}\right]^C} \quad ; \text{ and A, B, C, X, Y scaling coefficients for muons (see Dunai, 2001)} \quad (\text{b.11})$$

$\tan(I) = 2 \tan(\lambda)$  : inclination (time independent).

## c - Pigati and Lifton (2004), using Desilets and Zreda (2003)

Spallation rate at latitude  $\lambda$ , longitude  $\theta$  and elevation  $z$  :

$$S_{el,s} = \left[ \frac{J}{J_o} \right]_{(f,lat)} \left[ \frac{J_x}{J_{1033}} \right]_{(f,alt)} \quad (\text{c.1})$$

$$\bullet \text{ with } \left[ \frac{J}{J_o} \right]_{(f,lat)} = 1 - \exp(-\alpha_f R_c^{-k_f}) \text{ (equation from Desilets and Zreda, 2003)} \quad (\text{c.2})$$

$J$  : cosmic-ray flux at site,  
 $J_o$  : cosmic ray flux at Sea Level and High Latitude (SLHL),  
 $alt$  stands for altitude and  $lat$  for latitude,  
 $\alpha_f = 10.275$  and  $k_f = 0.9615$  : fitting parameters for neutrons.

$$R_c = \sum_{i=6}^{i=0} \left[ e_i + f_i \left( \frac{M}{M_o} \right) \right] \theta(t)^i \text{ (equation from Desilets and Zreda, 2003)} \quad (\text{c.3})$$

$M$  : geomagnetic field intensity (or dipole moment) at time t ( $\text{A.m}^{-2}$ ) from : 0-11 ka, Yang et al. (2000) and >11 ka, Guyodo and Valet (SINT800, 1999),  
 $M_o = 8.075 \cdot 10^{22} \text{ (Am}^{-2}\text{)}$  : modern geomagnetic field intensity (in 1945),  
 $\theta(t)$  : geomagnetic paleolatitude over time (see above in part (b) for equation) : 0-2 ka, Merrill and McElhinny (1983), Champion (1980) and 2-11 ka, Ohno and Hamano (1993).

$$\bullet \text{ with } \left[ \frac{J_x}{J_{1033}} \right]_{(f,alt)} = \exp \left[ \frac{x_2 - x_1}{\Lambda_{(eff,f,x)}} \right] \text{ (equation from Desilets and Zreda, 2003)} \quad (\text{c.4})$$

$x_2 = 1033 \text{ (g.cm}^{-2}\text{)}$  : atmospheric depth at sea level,  
 $x_1$  : atmospheric depth at site of interest ( $\text{g.cm}^{-2}$ ).  
 $\Lambda_{(eff,f,x)}$  : effective attenuation length for high energy (spallogenic) neutrons ( $\text{g.cm}^{-2}$ ).

$$x_1 = \left[ \frac{10P_o}{g_o} \right] \left[ \frac{T_o}{T_o + \beta_o z} \right]^{\left( \frac{g_o M}{1000 R \beta_o} \right)} \quad (c.5)$$

$P_o = 1013.25$  (hPa) : Sea level pressure,  
 $g_o = 9.8$  (m.s<sup>-2</sup>) : gravitational constant,  
 $T_o = 288.15$  (K) : Sea level temperature,  
 $\beta_o = -0.0065$  (K.m<sup>-1</sup>) : adiabatic lapse rate,  
 $M = 28.96$  (g.mol<sup>-1</sup>) : mean molecular weight of atmosphere,  
 $R = 8.31451$  (m.mol<sup>-1</sup>K<sup>-1</sup>) : gas constant.

$$\Lambda_{(eff,f,x)} = \frac{x_2 - x_1}{\left[ \frac{n(1 + \exp(-aR_c^{-k}))^{-1}(x_2 - x_1) + 1/2(b_0 + b_1R_c + b_2R_c^2)(x_2^2 - x_1^2) + 1/3(b_3 + b_4R_c + b_5R_c^2)(x_2^3 - x_1^3) + 1/4(b_6 + b_7R_c + b_8R_c^2)(x_2^4 - x_1^4)}{1} \right]} \quad (c.6)$$

$n, a, k, b_1$  to  $b_8$  : fitting parameters from Desilets and Zreda (2003).

**Slow muon capture at latitude  $\lambda$  longitude  $\theta$  and elevation  $z$  :**

$$S_{el,\mu} = \left[ \frac{J}{J_o} \right]_{(\mu,alt)} \left[ \frac{J_x}{J_{1033}} \right]_{(\mu,lat)} \quad (c.7)$$

• with  $\left[ \frac{J}{J_o} \right]_{(\mu,alt)} = 1 - \exp(-\alpha_\mu R_c^{-k_\mu})$  (equation from Desilets and Zreda, 2003) (c.8)

$\alpha_\mu$  and  $k_\mu$  : fitting parameters for muons.

• with  $\left[ \frac{J_x}{J_{1033}} \right]_{(\mu,lat)} = \exp \left[ \frac{x_2 - x_1}{\Lambda_{(eff,\mu,x)}} \right]$  (equation from Desilets and Zreda, 2003) (c.9)

and  $\Lambda_{(eff,\mu,x)} = 233 + 3.68R_c$  (c.10)

**d - Lifton et al. (2005)**

**Spallation rate at latitude  $\lambda$ , longitude  $\theta$  and elevation  $z$  :**

$$S_{el,s} = \exp [\ln(Flux, sp)] \quad (d.1)$$

with  $\ln(Flux, sp) = c_1 \ln(x_1 S) - S \exp \left[ \frac{c_2 S}{(R_c + 5S)^{2S}} \right] + c_3 (x_1)^{c_4} + c_5 [x_1 (R_c + 4S)]^{c_6} + c_7 (R_c + 4S)^{c_8}$  (d.2)

$x_1$  : atmospheric depth at site of interest (see equation c.5 in part (c)),  
 $S$  : sunspot modulation (over time) from Solanki et al. (2004),  
 $c_1$  to  $c_6$  : spallogenic coefficients (Lifton et al., 2005).

$$\text{and } R_c = d_1 \frac{M}{M_o} \cos^{d_2} \theta(t) \text{ (Cutoff rigidity)} \quad (\text{d.3})$$

$M$  : geomagnetic field intensity (or dipole moment) at time t ( $\text{A.m}^{-2}$ ) from : 0-11 ka, Yang et al. (2000) and >11 ka, Guyodo and Valet (SINT800, 1999),  
 $M_o = 8.075.10^{22}$  ( $\text{Am}^{-2}$ ) : modern geomagnetic field intensity (in 1945),  
 $\theta(t)$  : geomagnetic paleolatitude over time (see above in part (b) for equation) : 0-2 ka, Merrill and McElhinny (1983) and 2-11 ka, Ohno and Hamano (1993),  
 $d_1 = 15.765$  and  $d_2 = 3.800$  : scaling coefficients.

**Slow muon capture at latitude  $\lambda$  longitude  $\theta$  and elevation  $z$  :**

$$S_{el,\mu} = \exp [\ln(\text{Flux}, sm)] \quad (\text{d.4})$$

$$\text{with } \ln(\text{Flux}, sm) = a_1 + a_2 x_1 + a_3 x_1^2 + a_4 x_1 R_c + a_5 R_c + a_6 R_c^2 \quad (\text{d.5})$$

$x_1$  : atmospheric depth at site of interest (see equation (c.5)),  
 $R_c$  : Cutoff rigidity over time (equation (d.3)),  
 $a_1$  to  $a_6$  : coefficients for slow muons (Lifton et al., 2005).

**e - Pigati and Lifton (2004) modified using Desilets et al. (2006)**

Same equations as part (c), but  $\Lambda_{(eff,f,x)}$  calculated using new fitting parameters ( $n, a, k, b_1$  to  $b_8$ ) from Desilets et al. (2006).

**f - Lifton et al. (2008)**

Same equations as Lifton et al. (2005, part (d)) but  $R_c$  :

★ 0 - 7 ka :  $R_c$  values extracted from  $R_c$  grids from Korte and Constable (2005a,b ; 5° latitude and 15° longitude grid spacing),

★ > 7 ka :  $R_c$  calculation based on mean field over the entire 0 - 7 ka period

$$R_c = \frac{M}{M_o} \sum_{i=1}^6 d_{2i-1} \cos^i \left( \theta(t) + \frac{d_{2i}}{M/M_o} \right) \quad (\text{f.1})$$

$M$  : geomagnetic field intensity (or dipole moment) at time  $t$  ( $\text{A.m}^{-2}$ ) from : 0-11 ka, Yang et al. (2000) and >11 ka, Guyodo and Valet (SINT800, 1999),  
 $M_o = 8.075 \cdot 10^{22}$  ( $\text{A.m}^{-2}$ ) : modern geomagnetic field intensity (in 1945),  
 $\theta(t)$  : geomagnetic paleolatitude over time (see above in part (b) for equation) : 0-2 ka, Merrill and McElhinny (1983) and 2-11 ka, Ohno and Hamano (1993),  
 $d_1$  to  $d_{12}$  : coefficients for  $R_c$  calculation.

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	[Cl]nat	[36Cl]	uncertainty on [36Cl]
MA3c-18	11,1	1,40E+05	3,34E+03
MA3c-18R	11,4	1,45E+05	3,59E+03
<b>MA3c-18 mean</b>	<b>11,2</b>	<b>1,43E+05</b>	<b>2,79E+03</b>
MA3c-28	12,9	1,25E+05	3,14E+03
MA3c-28R	13,1	1,27E+05	3,76E+03
<b>MA3c-28 mean</b>	<b>13,0</b>	<b>1,26E+05</b>	<b>1,02E+03</b>
MA3c-42	10,8	1,05E+05	2,61E+03
MA3c-42R	11,1	1,05E+05	3,20E+03
<b>MA3c-42 mean</b>	<b>10,9</b>	<b>1,05E+05</b>	<b>3,64E+02</b>

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Replicates values from Magnola MA3c data-set (samples below the colluvium).

	1	2	3	4	5	6	7	8	9
Sample name	As ppm	Ba ppm	Be ppm	Bi ppm	Cd ppm	Ce ppm	Co ppm	Cr ppm	Cs ppm
MA3c - S1	4.162	120.6	1.724	0.191	1.087	37.14	5.989	26.82	3.697
MA3c - S2	8.421	221.0	2.748	0.335	1.465	63.26	9.904	49.34	6.626
MA3c - S4	5.608	153.7	1.949	0.217	1.085	43.16	7.314	33.04	4.373
MA3c - S8	3.924	130.5	1.577	0.199	0.660	36.71	5.470	20.67	4.998
MA3c - S10	0	41.27	0	0	0.422	9.897	2.653	8.241	1.303
MA3c - Soil mean	4.423	133.4	1.600	0.188	0.944	38.03	6.266	27.62	4.199
MA3c-P1	0	2.534	0	0	0.391	0.426	1.203	0	0
MA3c-P2	0	7.110	0	0	0.339	0.451	1.130	0	0
MA3c-P4	0	0	0	0	0.360	0.359	1.214	0	0
MA3c-P8	0	14.82	0	0	0.480	4.135	1.664	0	0.598
MA3c-P10	0	12.86	0	0	0.366	2.623	1.617	0	0.350
MA3c -Pebbles mean	0	7.465	0	0	0.387	1.599	1.366	0	0.190

	10	11	12	13	14	15	16	17	18	19
Cu ppm	Dy ppm	Er ppm	Eu ppm	Ga ppm	Gd ppm	Ge ppm	Hf ppm	Ho ppm	In ppm	
15.74	1.963	1.094	0.586	5.574	2.413	0.473	1.674	0.388	0	
21.86	3.264	1.746	0.982	10.37	3.933	0.836	2.862	0.615	0	
15.62	2.259	1.221	0.673	7.070	2.715	0.576	1.89	0.435	0	
8.956	1.713	0.898	0.556	4.394	2.223	0.402	1.625	0.324	0	
0	0.602	0.338	0.18	1.481	0.692	0.145	0.498	0.117	0	
12.43	1.960	1.059	0.595	5.778	2.395	0.486	1.71	0.376	0	
0	0.123	0.069	0.036	0	0.128	0	0.025	0.025	0	
0	0.134	0.086	0.027	0	0.121	0	0	0.029	0	
0	0.130	0.086	0.029	0	0.131	0	0.020	0.029	0	
0	0.414	0.218	0.139	0.512	0.487	0	0.194	0.078	0	
0	0.240	0.147	0.071	0.366	0.303	0	0.117	0.050	0	
0	0.208	0.121	0.06	0.176	0.234	0	0.071	0.042	0	

	20	21	22	23	24	25	26	27	28	29
La ppm	Lu ppm	Mo ppm	Nb ppm	Nd ppm	Ni ppm	Pb ppm	Pr ppm	Rb ppm	Sb ppm	
20.43	0.167	0	7.465	15.62	20.67	11.88	4.333	31.86	0.325	
32.61	0.265	0.649	12.67	25.70	29.60	21.54	6.997	61.88	0.509	
22.81	0.179	0	8.225	17.90	23.97	14.13	4.881	42.08	0.373	
20.06	0.133	0	6.382	15.21	19.28	11.77	4.219	30.69	0.352	
5.831	0.051	0	1.780	4.608	14.55	3.382	1.247	10.05	0	
20.348	0.159	0.13	7.304	15.808	21.614	12.54	4.335	35.312	0.312	
1.346	0.010	0	0.059	0.951	10.05	0	0.260	0	0	
0.886	0.012	0	0.053	0.627	9.672	0	0.162	0.368	0	
1.176	0.012	0	0	0.856	10.21	0	0.231	0.303	0	
4.864	0.032	0	0.683	3.722	10.72	1.801	1.011	5.409	0	
2.290	0.023	0	0.410	1.879	11.41	1.250	0.496	2.966	0	
2.112	0.018	0	0.241	1.607	10.412	0.61	0.432	1.809	0	

	30	31	32	33	34	35	36	37	38	39
Sm ppm	Sn ppm	Sr ppm	Ta ppm	Tb ppm	Th ppm	Tm ppm	U ppm	V ppm	W ppm	
2.873	1.143	145.7	0.528	0.359	6.849	0.159	1.174	32.12	0.653	
4.760	1.932	130.3	0.891	0.567	12.06	0.254	1.775	60.01	1.206	
3.351	1.401	138.9	0.588	0.392	7.813	0.171	1.256	41.91	0.849	
2.728	1.067	122.0	0.418	0.310	8.572	0.130	1.021	30.18	0.633	
0.847	0.457	96.89	0.129	0.102	1.978	0.048	0.486	10.82	0	
2.912	1.2	126.758	0.511	0.346	7.454	0.152	1.142	35.008	0.668	
0.160	0	148.9	0	0.018	0.062	0.010	0.231	0	0	
0.118	0	155.9	0	0.019	0.047	0.012	0.383	0	0	
0.143	0	154.1	0	0.020	0.036	0.012	0.407	0	0	
0.660	0	122.5	0.048	0.070	1.010	0.032	0.383	4.952	0	
0.373	0	106.3	0.029	0.042	0.515	0.020	0.223	2.975	0	
0.291	0	137.54	0.015	0.034	0.334	0.017	0.325	1.585	0	

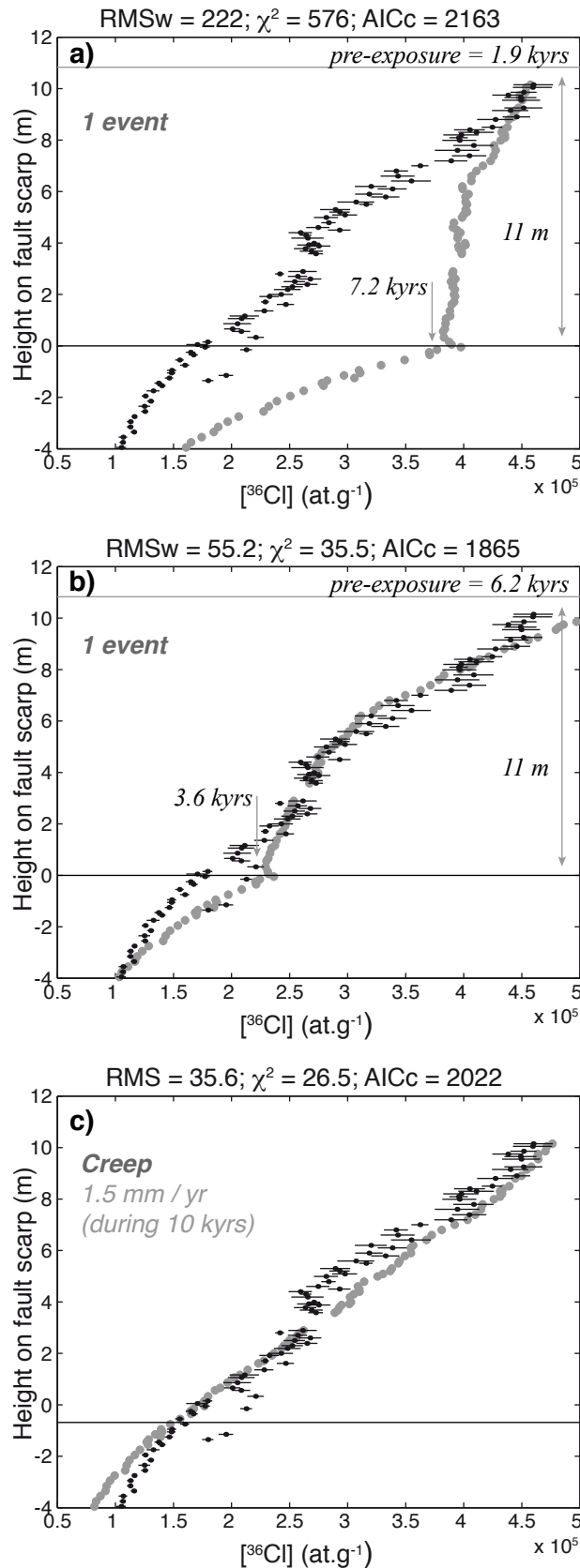
	40	41	42	43	44	45	46	47	48	49
Y ppm	Yb ppm	Zn ppm	Zr ppm	SiO2 (Si) %	Al2O3 (Al) %	Fe2O3 (Fe) %	MnO (Mn) %	MgO (Mg) %	CaO (Ca) %	
13.41	1.046	42.40	73.07	9.915	4.575	1.722	0.036	0.785	43.54	
20.94	1.705	73.40	124.3	19.07	8.355	3.228	0.065	1.115	33.46	
14.64	1.198	50.42	81.76	12.56	5.525	2.167	0.045	0.901	40.76	
10.56	0.869	33.92	74.15	9.294	4.024	1.551	0.031	0.722	45.15	
4.260	0.317	15.41	20.73	3.098	1.246	0.515	0.011	0.43	52.23	
12.762	1.027	43.11	74.802	10.787	4.745	1.837	0.038	0.791	43.03	
1.157	0.07	0	1.358	0	0	0	0.004	0.55a	55.21	
1.385	0.078	9.205	0	0	0	0	0.006	0.543	55.41	
1.443	0.077	8.114	0	0	0	0	0.004	0.556	55.18	
3.089	0.202	8.739	9.120	0.943	0.458	0.202	0.007	0.509	54.15	
1.962	0.135	0	5.759	0.604	0.297	0.130	0.004	0.389	54.31	
1.807	0.112	5.212	3.247	0.309	0.151	0.066	0.005	0.499	54.85	

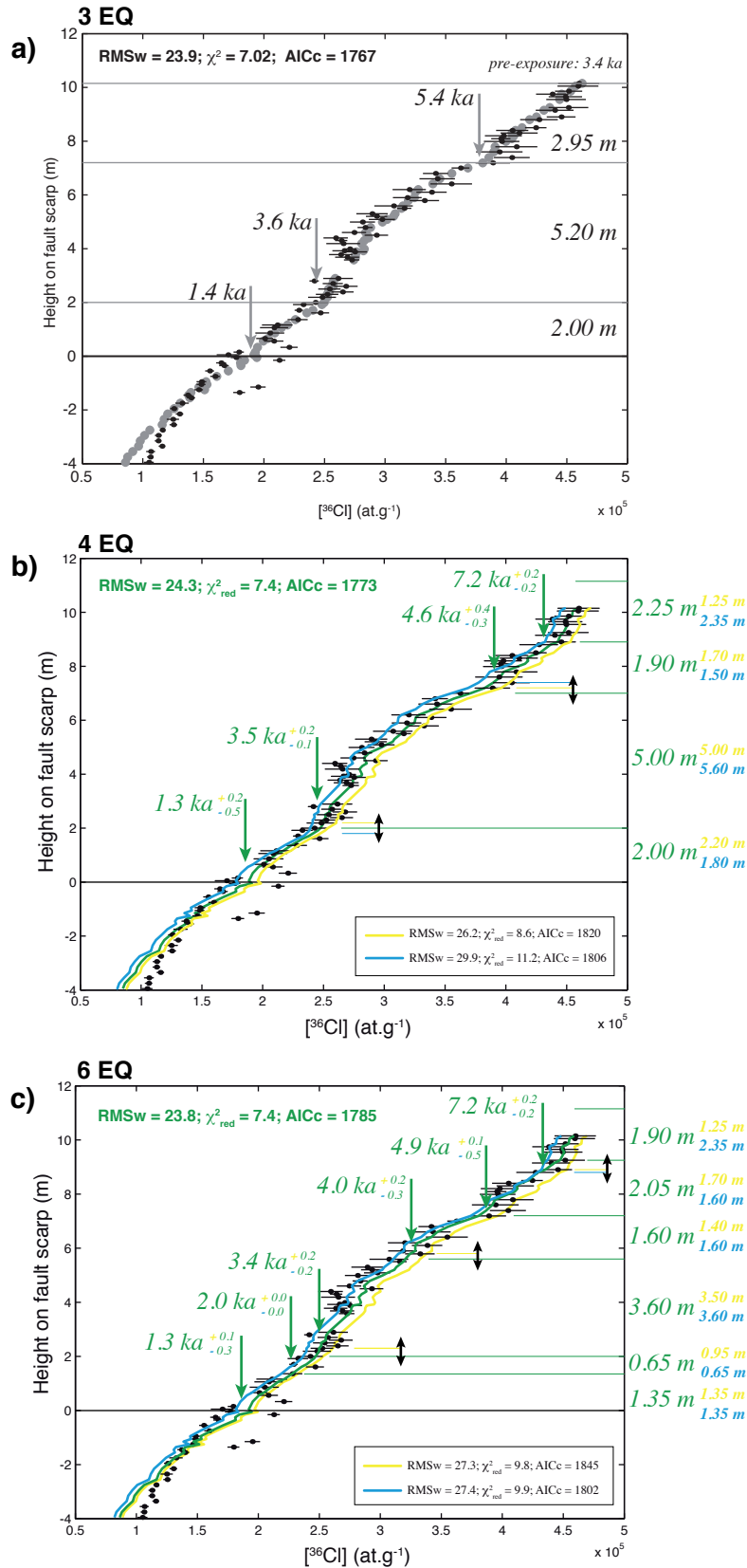
	50	51	52	53	54	55	56	57	58
Na2O (Na) %	K2O (K) %	TiO2 (Ti) %	P2O5 (P) %	B ppm	Li ppm	H2Otot (H) %	Stot (S) %	CO2 tot (C) %	
0	0.370	0.213	0.286	11.00	15.00	3.790	0.040	37.92	
0.151	0.762	0.386	0.342	23.00	27.60	6.020	0.060	31.91	
0.094	0.497	0.259	0.256	17.00	17.70	4.470	0.040	36.10	
0.078	0.413	0.192	0.216	11.00	11.30	3.050	0.020	36.56	
0	0.146	0.064	0.106	6.000	3.900	1.280	0.020	41.58	
0.065	0.438	0.223	0.241	13.6	15.1	3.722	0.036	36.814	
0	0	0	0.061	2.000	1.000	0.230	0.020	44.18	
0	0	0	0.143	2.000	0.900	0.240	0.010	44.16	
0	0	0	0.086	2.000	1.000	0.240	0.020	44.17	
0	0.084	0.024	0.079	2.000	2.000	0.430	0.010	43.16	
0	0	0.017	0.044	2.000	1.600	0.490	0.010	43.42	
0	0.017	0.008	0.083	2	1.3	0.326	0.014	43.818	

Chemical composition of five samples from the colluvial wedge at MA3 site (Magnola Fault) and corresponding mean compositions. See Figure 15 for position of the samples. Pebbles and soil were analyzed separately. Note that when using the 'most concentrated' sample S2 as mean composition of the colluvial wedge, production rate is similar to the one modeled with the mean composition over the five samples.

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Magnola fault data set (MA3 site). Samples below  $z = 0$  are those from the buried scarp under the colluvium. a,b) Bests fits for a unique event (one earthquake or a landslide). a) Fit of the topmost segment, one event of 10 meters 7.2 kyrs ago (with pre-exposure fixed at 1.9 kyrs as Fig. 18a). b) Best fit of the complete data set with an event of same displacement 3.6 kyrs ago (pre-exposure parameter free: best fit for 6.2 kyrs). c) Modeling creep on the fault. Best fit is obtained for slip-rate of 1.5 mm/yr over 10 kyrs. In all cases, fit is not satisfactory.



Best fit of the Magnola data set from MA3 site (black dots and respective uncertainties as horizontal lines) by 3, 4 and 6 large slip exhumations. Models have been calculated with no erosion and for a constant Earth's magnetic field.

a) 3 events model. The data best-fit function is in grey. Earthquake ages and displacements are indicated.  
 b-c) 4 and 6 events models. The data best-fit function is in green, while fitting the data uncertainties produce the yellow and blue curves. Double arrows indicate the uncertainties on the displacements. Earthquake ages are indicated together with their uncertainties.



## **Annexe III**

(du chapitre 5)



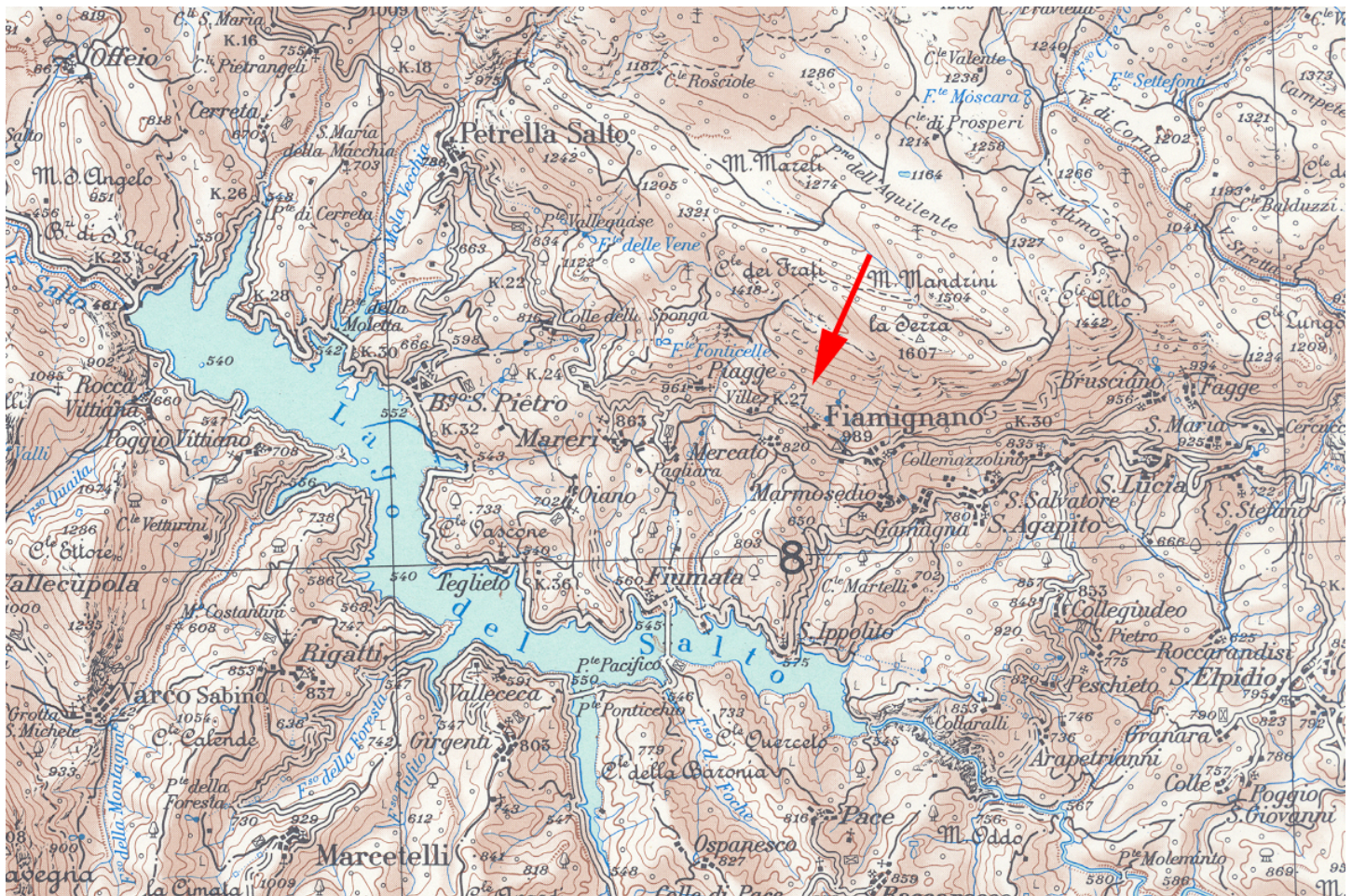
## RECAPITULATIF SITES ABRUZZES (ITALIE)

FAILLE	Nom du site	NOUVEAU NOM	date éch.	rho_rock	rho_coll	alpha	beta	gamma	H	altitude	latitude	longitude	EL_f Stone 2000	EL_mu Stone 2000	H echantillons
<b>Magnola</b>	M1	<b>MA3</b>	2003	2.7	1.5 (mesuré)	30°	45°	30°	20m	1255m +/- 5m	N42.11900	E013.44800	2.750	1.703	10.15 + 4.0 m
	M1c		Mai 2007												
<b>Magnola</b>	A1c (W)	<b>MA1</b>	Sept. 2005	2.67	1.5 ?	25°	40°	35°	15m	1265m +/- 5m	N42.12800	E013.41371	2.771	1.711	8.10 m
<b>Magnola</b>	A4c (E)	<b>MA4</b>	Sept. 2005	2.64	1.5 ?	30°	42°	30°	7m	1300m +/- 10m	N42.11840	E013.46135	2.846	1.737	7.00 m
<b>Magnola</b>	A6c	<b>MA2</b>	Avril 2006	2.67	1.5 ?	35°	50°	35°	8.6m	1200m +/- 5m	N42.12400	E013.42737	2.636	1.662	8.55 m
<b>Velino</b>	<b>VE</b>		Mai 2007	2.71	1.5 ?	30°	40°	35°	9.50	1014m +/-5m	N42.16858	E013.31302	2.281	1.531	7.01 m
<b>Campo Felice</b>	<b>CF</b>		Mai 2007	2.66	1.5 ?	30°	55°	45°	17m	1595m +/-5m	N42.22832	E013.4437	3.558	1.974	8.85 m
<b>Fiamignano</b>	<b>FI</b>		Mai 2007	2.68	1.5 ?	12°	63°/40°	38°	10m	1178m +/-5m	N42.27208	E013.11618	2.597	1.648	7.10 m
<b>Tre Monti</b>	<b>TM</b>		Mai 2007	2.67	1.5 ?	20°	72°	35°	3.5m	970m +/-5m	N42.06505	E013.4627	2.200	1.499	2.70 m
<b>Trasacco</b>	<b>TR</b>		Mai 2007	2.64	1.5 ? - 2.0 ?	25°	65°	25°	4.7m	800m +/-5m	N41.92805	E013.56700	1.916	1.386	4.45 m
<b>Parasano</b>	<b>PA</b>		Mai 2007	2.64	1.5 ?	25°	55°	35°	4.5m	1270m +/-5m	N41.99613	E013.70275	2.777	1.713	3.30 m
<b>San Sebastiano</b>	<b>SB</b>		Mai 2007	2.62	1.5 ?	25°	62°	35°	8.5m	1214m +/-5m	N41.94667	E013.76185	2.658	1.671	4.45 m
<b>Roccapreturo</b>	A2R (W)	<b>RP1</b>	Avril 2006	2.69	1.5 ?	20°	60°	35°	10m	795m +/- 5m	N42.21846	E013.65015	1.915	1.386	7.35 m
<b>Roccapreturo</b>	A1R	<b>RP2</b>	Avril 2006	2.66	1.5 ?	25°	55°	35°	10.25m	805m +/-10m	N42.19143	E013.71244	1.930	1.392	10.20 m
<b>Roccapreturo</b>	A3R (E)	<b>RP3</b>	Avril 2006	2.65	1.5 ?	25°	60°	40°	10m	783m +/-7m	N42.18127	E013.72890	1.896	1.378	5.35 m
<b>Castel di Ieri</b>	<b>CI</b>		Mai 2007	2.64	1.5 ?	15°	45°	40°	4m	570 +/- 5 m	N42.1057	E013.90263	1.601	1.253	3.0 m
<b>Rocccasale</b>	<b>RC</b>		Mai 2007	2.68	1.5 ?	30°	55°	25°	5.7m	392m +/-5m	N42.12515	E013.88250	1.366	1.147	4.70 m

# Faïlle de Fiamignano



Panorama de la faille de Fiamignano (vue vers le nord)



Extrait de la feuille n°145 (AVEZZANO) au 1:100 000ème, IGM

## Faille de Fiamignano - Site FI-A

*Mission Mai 2007*

**GPS FI-A:** alt. 1178 m  $\pm$  5m ; N42.27208 - E013.11618

**Ecrantage :** seulement le plan de faille, à la base **115 SW 63** (1m de hauteur), au dessus **115 SW 40**

**Nbre échantillons :** 95 (FI-1 à FI-95)

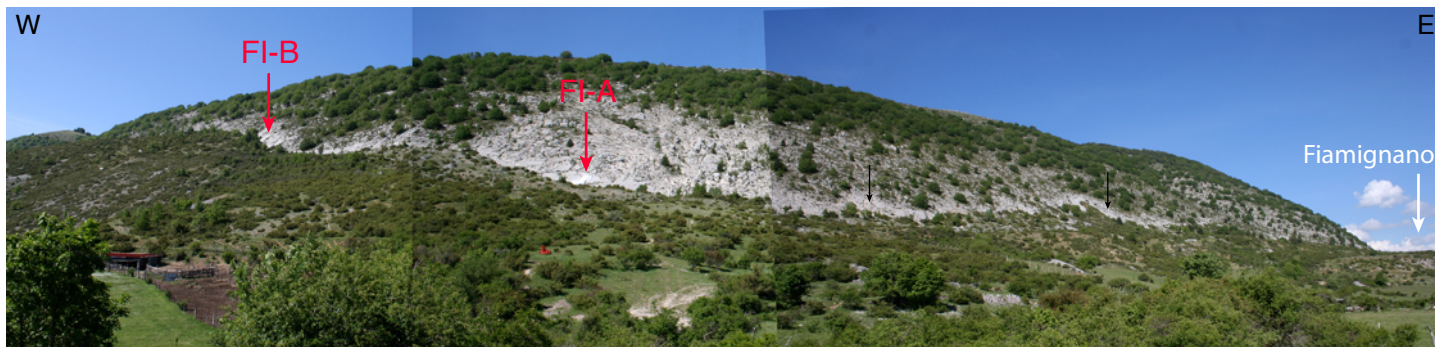
+ 2 échantillons au dessus : **FI-TOP1** (alt. 1188 m  $\pm$  5m ;  $\gamma = 38^\circ$ )

**FI-TOP2** (alt. 1191 m  $\pm$  5m ;  $\gamma = 38^\circ$ )

+ **FI-COL** (alt. 1176 m  $\pm$  5m ;  $\alpha = 12^\circ$ )

$\alpha = 12^\circ$  ;  $\beta = 63^\circ$  puis  $40^\circ$  ;  $\gamma = 38^\circ$

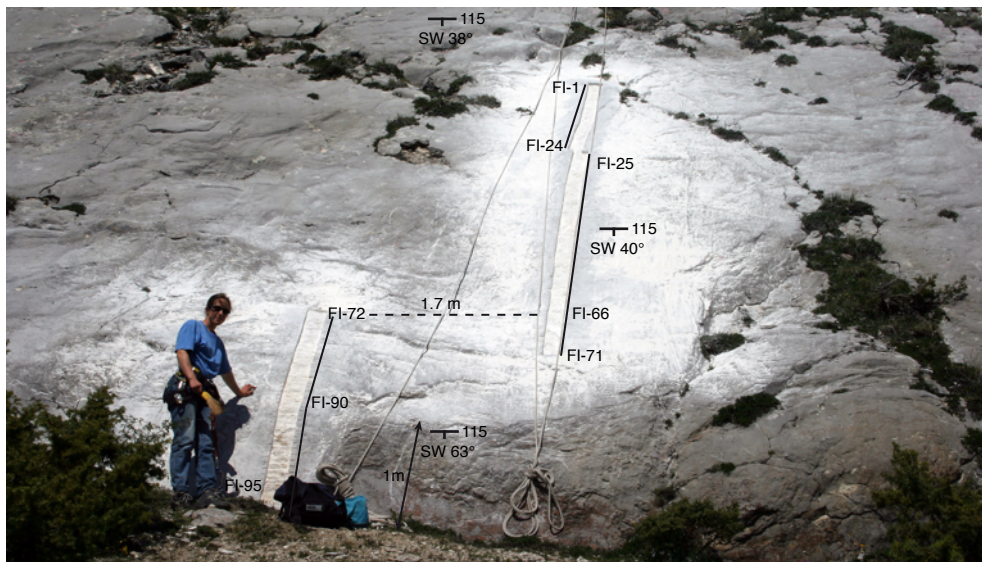
H = 10 m



Vue vers le nord des sites FI (A et B) au nord-ouest du village de Fiamignano



Site principal FI-A  
(échantillons FI-1 à FI-95)



## Faille de Fiamignano - Site FI-B

*Mission Mai 2007*

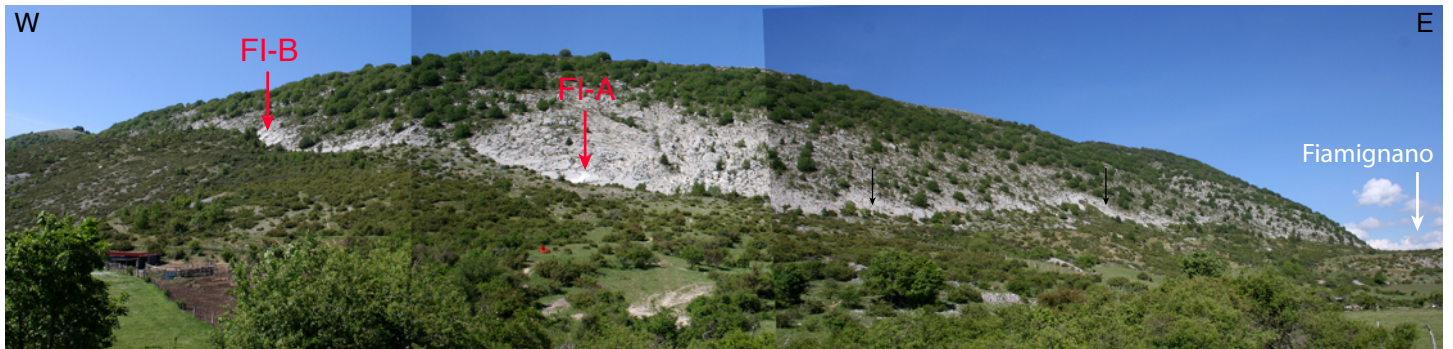
**GPS FI-B:** alt. 1200 m  $\pm$  5m ; N42.27378 - E013.11400

**Ecrantage :** seulement le plan de faille **135 SW 40**

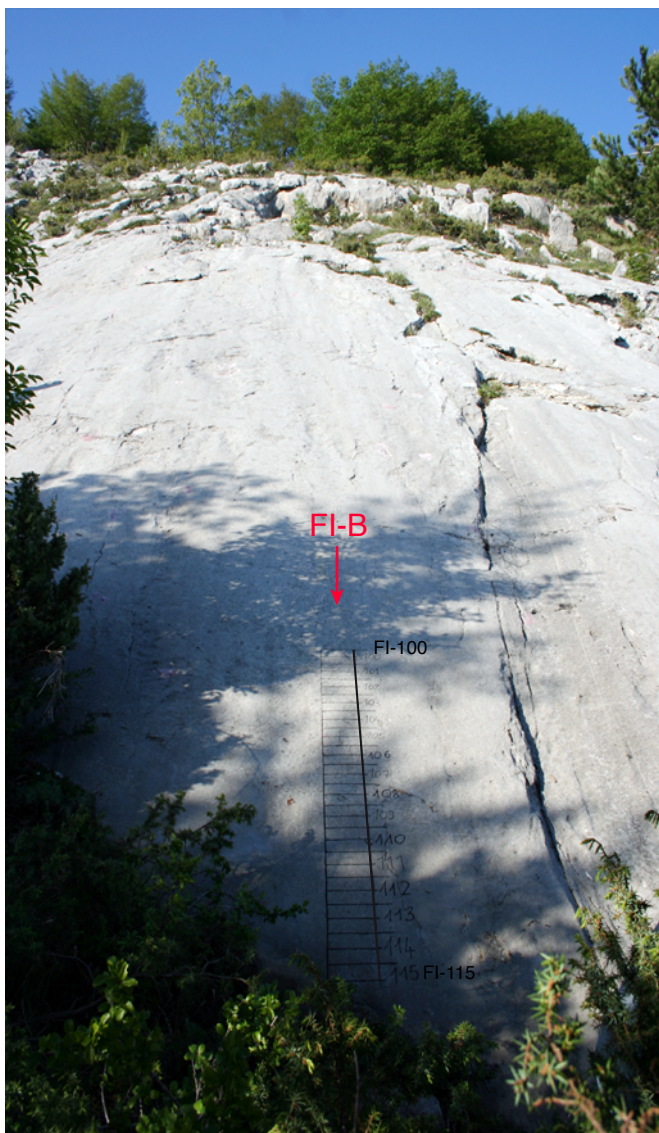
**Nbre échantillons :** 15 (FI-100 à FI-115)

$\alpha = 20^\circ$  ;  $\beta = 40^\circ$  ;  $\gamma = 38^\circ$

H = 7 m



*Vue vers le nord des sites FI (A et B) au nord-ouest du village de Fiamignano*



*Site secondaire FI-B  
(échantillons FI-100 à FI-105)*

## Faille de Campo Felice

Mission Mai 2007

GPS : alt. 1595 m  $\pm$  5m ; N42.22832 - E013.44370

Ecrantage : seulement le plan de faille 105 S 55

Nbre échantillons : 77 (CF-1 à CF-77)

+ 1 échantillon au dessus : **CF-TOP** (alt. 1612 m  $\pm$  5m ;  $\gamma = 45^\circ$ )

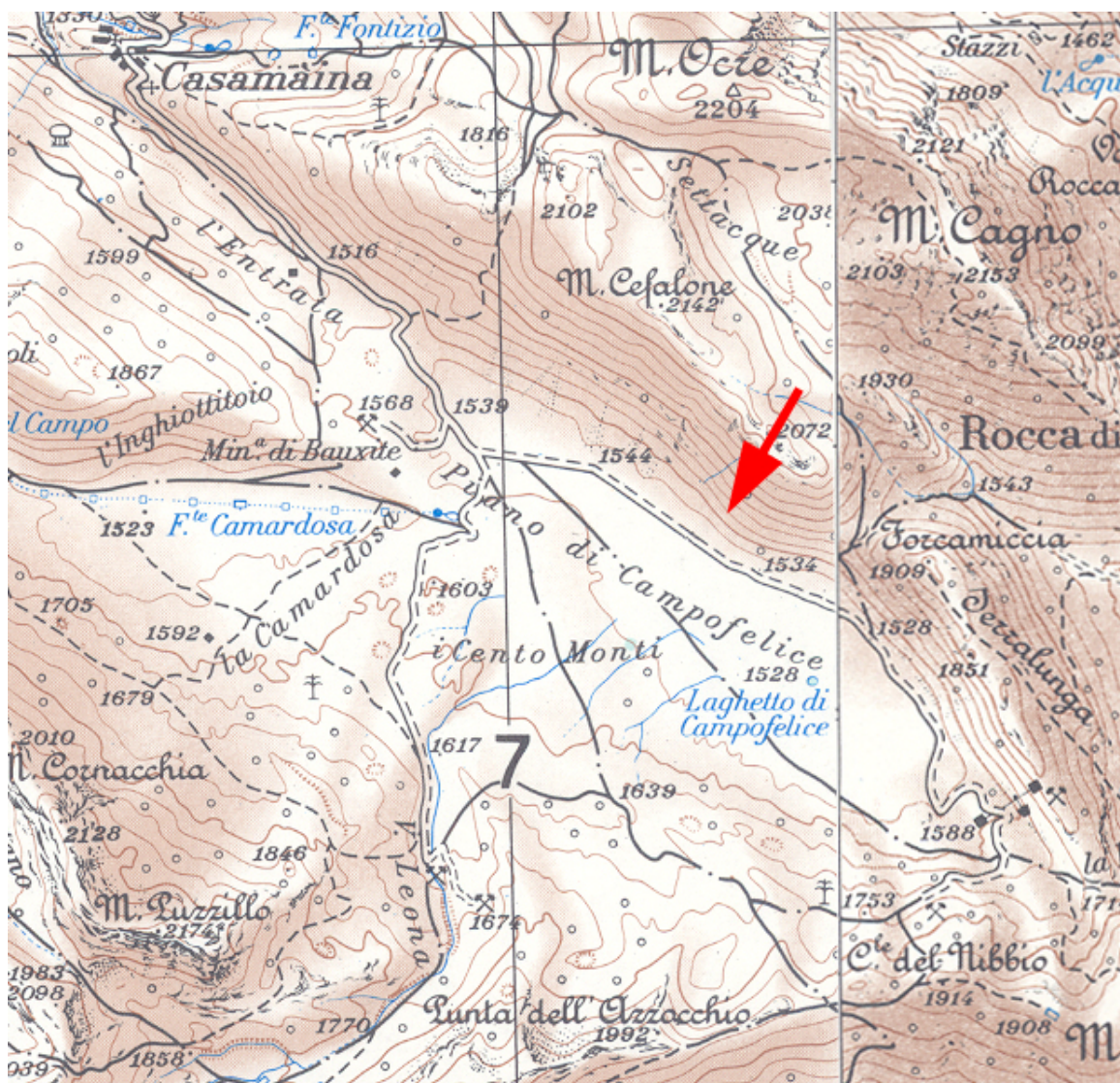
+ **CF-COL** (alt. 1590 m  $\pm$  5m ;  $\alpha = 30^\circ$ )

$\alpha = 30^\circ$  ;  $\beta = 55^\circ$  ;  $\gamma = 45^\circ$

H = 17 m



Panorama de la faille de Campo Felice (vue vers le nord)

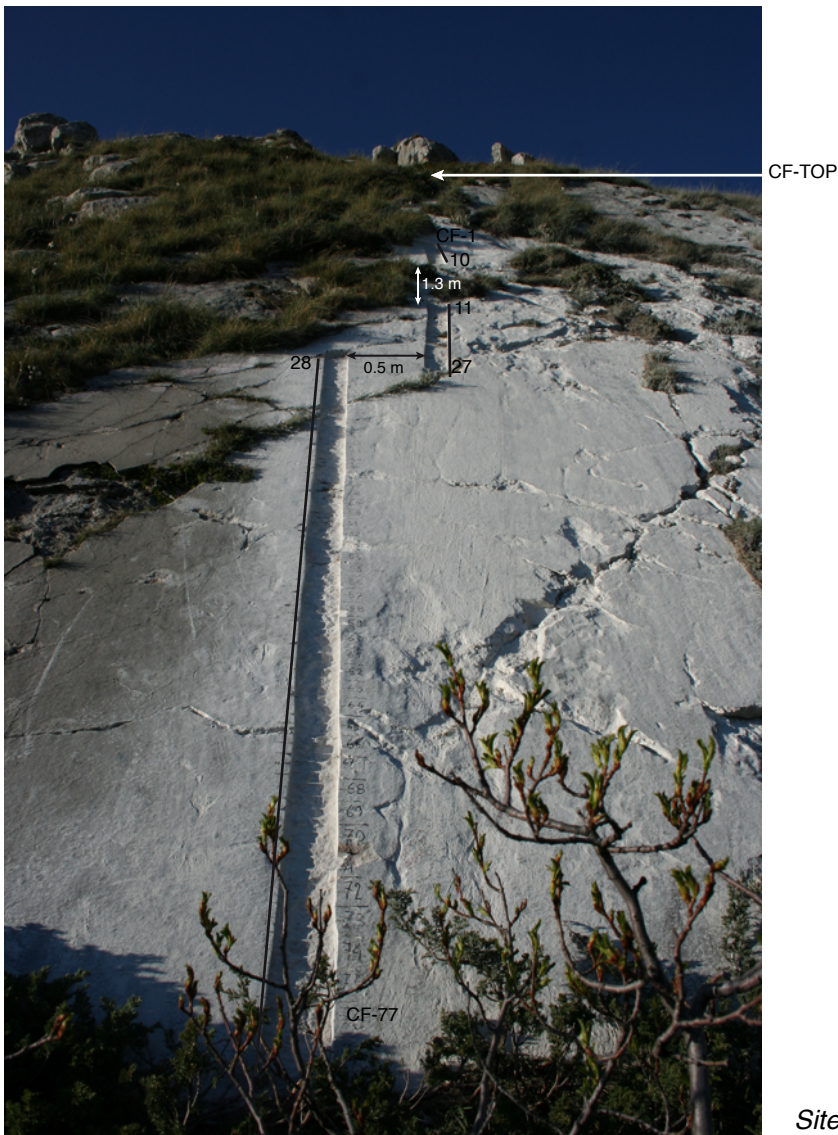


Extrait des feuilles n°145 (AVEZZANO) et n°146 (SULMONA) au 1:100 000ème, IGM

# Site CF



*Faille de Campo Felice, localisation du site d'échantillonnage CF*





## Faille du Mont Velino - Site VE

Site au dessus du village de Cartore

Mission Mai 2007

GPS : alt. 1014 m  $\pm$  5m ; N42.16858 - E013.31302

Ecrantage : seulement le plan de faille 130 SW 40

Nbre échantillons : 76 (VE-1 à VE-77 ; pas de VE-40)

+ 2 échantillons au dessus : VE-TOP1 (alt. 1023 m  $\pm$  5m ;  $\gamma = 35^\circ$ )

VE-TOP2 (alt. 1022.5 m  $\pm$  5m ;  $\gamma = 35^\circ$ )

+ VE-COL (alt. 1011 m  $\pm$  5m ;  $\alpha = 30^\circ$ )

$\alpha = 30^\circ \pm 3^\circ$  ;  $\beta = 40^\circ \pm 2^\circ$  ;  $\gamma = 35^\circ \pm 5^\circ$

H = 9.5 m



Faille du Mont Velino vue depuis le village de Cartore (vue vers le nord)



Extrait de la feuille n°145 (AVEZZANO) au 1:100 000ème, IGM



Site VE

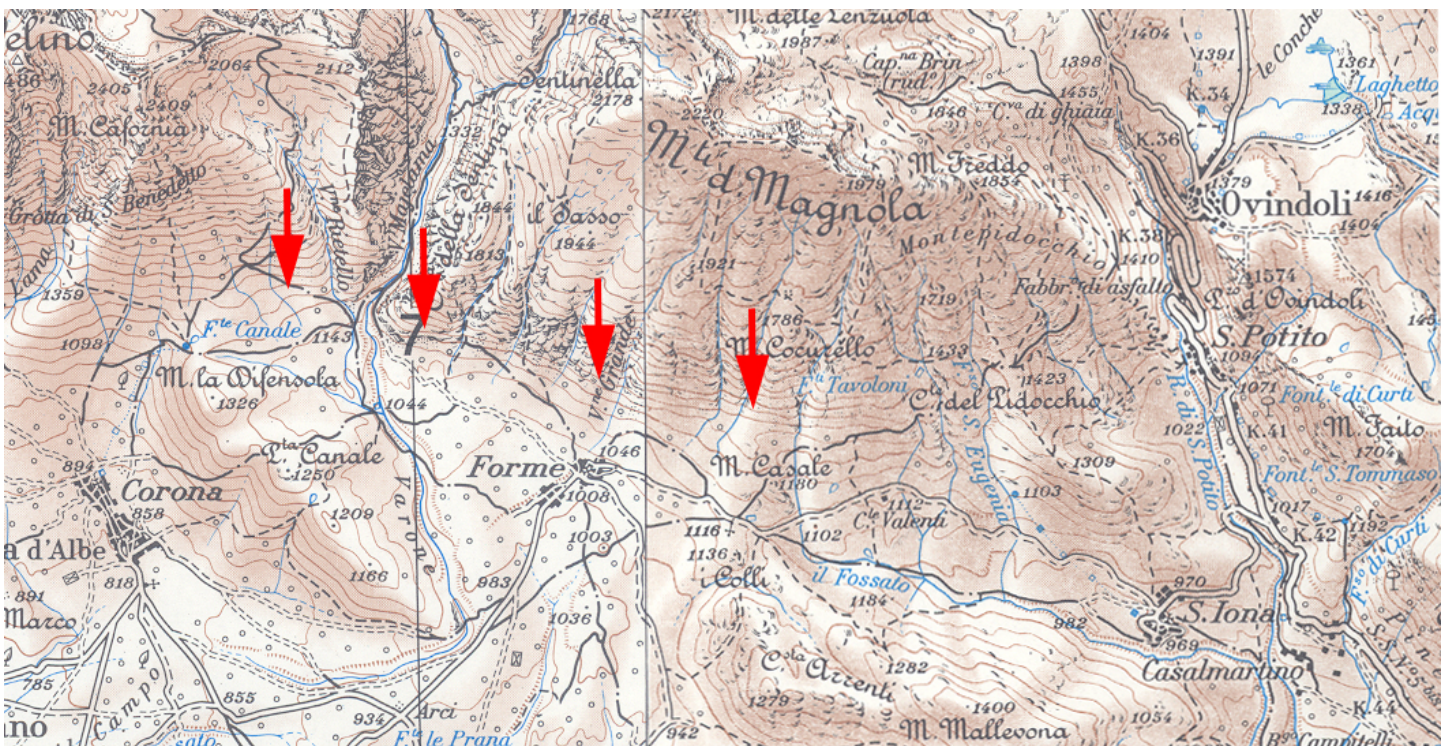
## Faille de la Magnola



Panorama de la faille de la Magnola vue depuis la route au sud du village de Forme (vue vers le nord)



Panorama de la faille de la Magnola vue depuis le village de Forme (vue vers le nord)



Extrait des feuilles n°145 (AVEZZANO) et n°146 (SULMONA) au 1:100 000ème, IGM

## Faille de la Magnola - Site MA1

Terminaison Ouest de la faille de la Magnola (à l'ouest de la Vallée Majelama)

*Mission Sept. 2005*

**GPS** : alt. 1265 m  $\pm$  5m ; N42.12800 - E013.41371

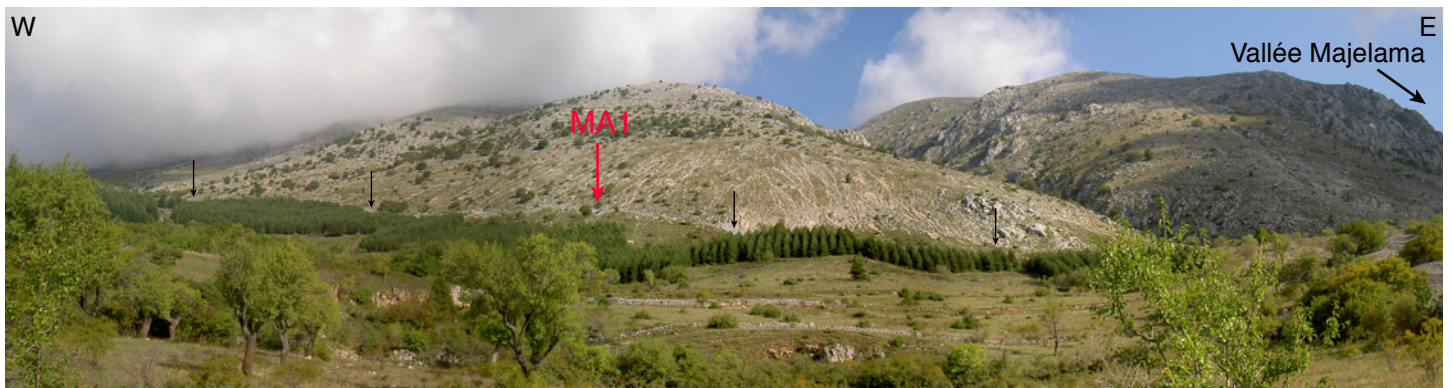
**Ecrantage** : seulement le plan de faille **110 S 42**

**Nbre échantillons** : **70** (A1c-1 à A1c-70)

+ 1 échantillon au dessus : **A1c-TOP** (alt. 1305 m  $\pm$  5m ;  $\gamma = 35^\circ$ )

$\alpha = 25^\circ \pm 3^\circ$  ;  $\beta = 40^\circ \pm 2^\circ$  ;  $\gamma = 35^\circ \pm 5^\circ$

H = 15 m



*Vue vers le nord de l'extrémité ouest de la faille de la Magnola et localisation du site MA1*



*Site MA1 (anciennement nommé A1c)*

## Faille de la Magnola - Site MA2

(à l'est de la Vallée Majelama)

Mission Avril 2006

GPS : alt. 1200 m  $\pm$  5m ; N42.12400 - E013.42737

Ecrantage : seulement le plan de faille **100 S 50**

Nbre échantillons : **74** (A6c-1 à A6c-74)

$\alpha = 35^\circ \pm 3^\circ$  ;  $\beta = 50^\circ \pm 2^\circ$  ;  $\gamma = 35^\circ \pm 5^\circ$

H = 8.6 m



Vue vers le nord de la faille de la Magnola (à l'est de la vallée Majelama) et localisation du site MA2



Site MA2 (anciennement nommé A6c)

## Faille de la Magnola - Site MA3

Site central de la faille de la Magnola (au dessus du village de Forme)

Echantillonnage "MAG" lors de la thèse de L. Palumbo ; "M1c" mission Mai 2007

GPS : alt. 1255 m  $\pm$  5m ; N42.11900 - E013.44800

Ecrantage : seulement le plan de faille **110 S 45**

Nbre échantillons : **65** (MAG-1 à MAG-65)

+ échantillons du scarp sous le colluvion: **42** (M1c-1 à M1c-42)

+ cailloux dans le colluvion : **14** (M1S-1 à M1S-14)

+ <sup>14</sup>C : **2** échantillons (M1carb-1 = 31.6  $\pm$  0.5 ka ; M1carb-2 = 38.8  $\pm$  1.2 ka)

$\alpha = 30^\circ \pm 3^\circ$  ;  $\beta = 45^\circ \pm 2^\circ$  ;  $\gamma = 30^\circ \pm 5^\circ$

H = 20 m



Vue vers le nord de la partie centrale de la faille de la Magnola et localisation du site MA3



Site MA3 (anciennement nommé M1)

## Faille de la Magnola - SITE A4c

Terminaison Est de la faille de la Magnola

Mission Sept. 2005

GPS : alt. 1300 m  $\pm$  10m ; N42.11840 - E013.46135

Ecrantage : seulement le plan de faille 97 S 42

Nbre échantillons : 64 (A4c-1 à A4c-64)

$\alpha = 30^\circ \pm 3^\circ$  ;  $\beta = 42^\circ \pm 2^\circ$  ;  $\gamma = 30^\circ \pm 5^\circ$

H = 7 m



Vue vers le nord de la terminaison Est de la faille de la Magnola et localisation du site MA4



Site MA4 (anciennement nommé A4c)

## Faille de Tre Monti - Site TM

Site au dessus du hameau de San Pelino Vecchio  
Mission Mai 2007

GPS : alt. 970 m  $\pm$  5m ; N42.06505 - E013.46270

Ecrantage : seulement le plan de faille **65 SSE 72**

Nbre échantillons : **27** (TM-1 à TM-27)

+ 1 échantillon au dessus : **TM-TOP** (alt. 977 m  $\pm$  5m ;  $\gamma = 35^\circ$ )

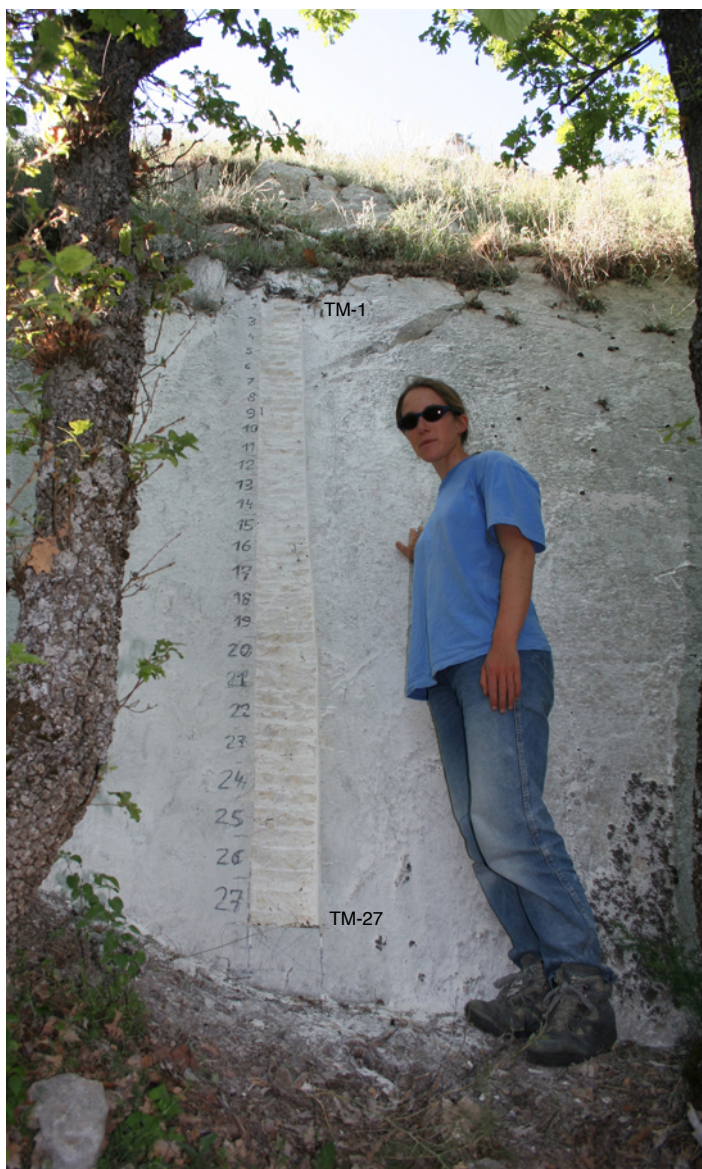
+ **TM-COL** (alt. 968 m  $\pm$  5m ;  $\alpha = 20^\circ$ )

$\alpha = 20^\circ$  ;  $\beta = 72^\circ$  ;  $\gamma = 35^\circ$

H = 3.5 m



Faille de Tre Monti vue depuis le le bassin du Fucino (vue vers le nord)



Extrait des feuilles n°145 (AVEZZANO) et n°146 (SULMONA)  
au 1:100 000ème, IGM

### Site TM

La base sur 60 cm de hauteur a un aspect plus lisse tout  
au long de la faille de Tre Monti. Stries normales-dextres.

## Faille de Trasacco (site TR)

Au sud-est du village de Trasacco

Mission Mai 2007

GPS : alt. 800 m  $\pm$  5m ; N41.92805 - E013.56700

Ecrantage : seulement le plan de faille **110 S 65**

Nbre échantillons : **45** (TR-1 à TR-45)

+ 1 échantillon au dessus : **TR-TOP** (alt. 810 m  $\pm$  5m ;  $\gamma = 25^\circ$ )

+ 9 échantillons du colluvion exhumé (à 20 m à l'est du site principal, 90cm exhumés, 50cm à l'avant du scarp principal : TR-C1 à TR-C9)

$\alpha = 25^\circ \pm 3^\circ$  ;  $\beta = 65^\circ \pm 2^\circ$  ;  $\gamma = 25^\circ \pm 5^\circ$

H = 4.7 m



Faille de Trasacco (vue vers le nord)



Extrait de la feuille n°152 (SORA) au 1:100 000ème, IGM



## Site TR



Site TR. A cet endroit, le colluvion induré est décalé verticalement de 30cm, 2m devant le plan principal



TR-C1 à TR-C9 : Colluvion induré exhumé (90cm de hauteur, 50cm devant le plan principal)

## Faille de Parasano (site PA)

*Mission Mai 2007*

**GPS** : alt. 1270 m  $\pm$  5m ; N41.99613- E013.70275

**Ecrantage** : seulement le plan de faille **125 SW 55**

**Nbre échantillons** : **33** (PA-1 à PA-33)

$\alpha = 25^\circ$  ;  $\beta = 55^\circ$   $\gamma = 35^\circ$

H = 4.5 m



## Faille de San Sebastiano (site SB)

*Mission Mai 2007*

**GPS** : alt. 1214 m  $\pm$  5m ; N41.94667- E013.76185

**Ecrantage** : seulement le plan de faille **0 W 62**

**Nbre échantillons** : 45 (SB-1 à SB-45)

$\alpha = 25^\circ$  ;  $\beta = 62^\circ$   $\gamma = 35^\circ$

H = 8.5 m



## Faille de Castel di Ieri (site CI)

*Mission Mai 2007*

**GPS** : alt. 570 m  $\pm$  5m ; N42.1057 - E013.90263

**Ecrantage** : seulement le plan de faille **135 SW 45**

**Nbre échantillons** : **29** (CI-1 à CI-29)

+ CI-TOP : à 10. 5 m de la base

+ CI-col : à 4 m de la base

$\alpha = 15^\circ$  ;  $\beta = 45^\circ$   $\gamma = 40^\circ$

H = 4 m



## Faille de Roccapreturo - site RP1

*Mission Avril 2006 (ancien nom du site : A2R)*

**GPS** : alt. 795 m  $\pm$  5 m ; N42.21846 - E013.65015

**Ecrantage** : seulement le plan de faille **125 SW 60**

**Nbre échantillons** : 71 (A2R-1 à A2R-71)

(pas de A2R-49!)

$\alpha = 20^\circ \pm 3^\circ$  ;  $\beta = 60^\circ \pm 2^\circ$  ;  $\gamma = 35^\circ \pm 5^\circ$

H = 10 m



## Faille de Roccapreturo - site RP2

Site à ~2km à l'est du village de Roccapreturo

Mission Avril 2006 (ancien nom du site : A1R)

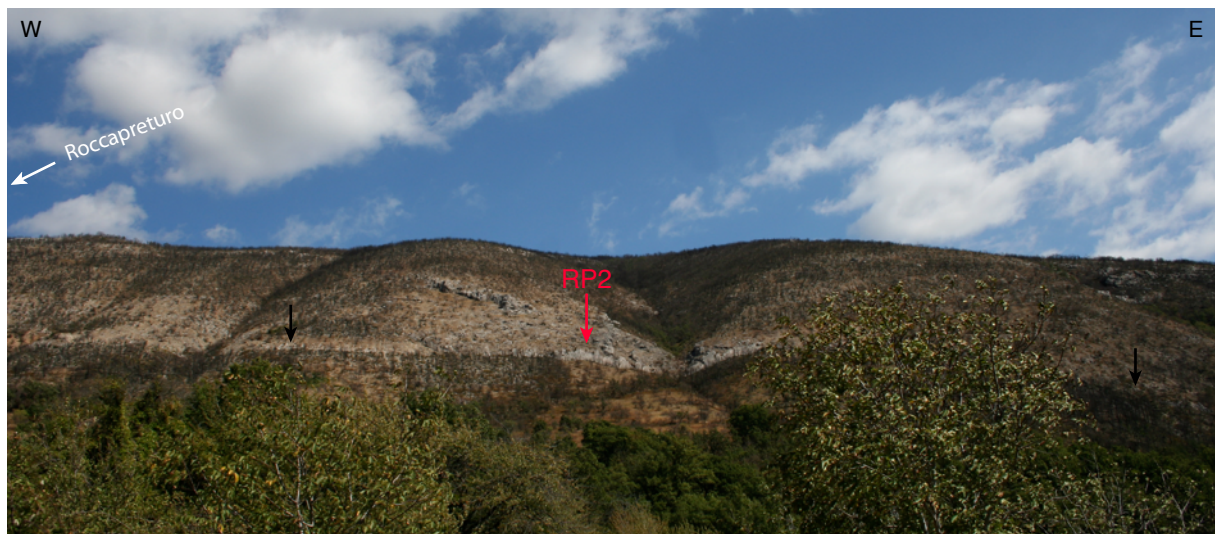
GPS : alt. 805 m  $\pm$  10m ; N42.19143 - E013.71244

Ecrantage : seulement le plan de faille **125 SW 55**

Nbre échantillons (deux verticales) : **75 à droite** (A1R-1 à A1R-75)  
et **25 à gauche** (A1R-76 à A1R-100)

$\alpha = 25^\circ \pm 3^\circ$  ;  $\beta = 55^\circ \pm 2^\circ$  ;  $\gamma = 35^\circ \pm 5^\circ$

H = 10.2 m



Faille du Roccapreturo vue vers le nord



Site RP2 sur la faille de Roccapreturo (anciennement nommé A1R)

## Faille de Roccapreturo - site RP3

*Mission Avril 2006 (ancien nom du site : A3R)*

**GPS** : alt. 783 m  $\pm$  5 m ; N42.18127 - E013.72890

**Ecrantage** : seulement le plan de faille **115 S 60**

**Nbre échantillons** : **53** (A3R-1 à A3R-53)

$\alpha = 25^\circ \pm 3^\circ$  ;  $\beta = 60^\circ \pm 2^\circ$  ;  $\gamma = 40^\circ \pm 5^\circ$

H = 10 m



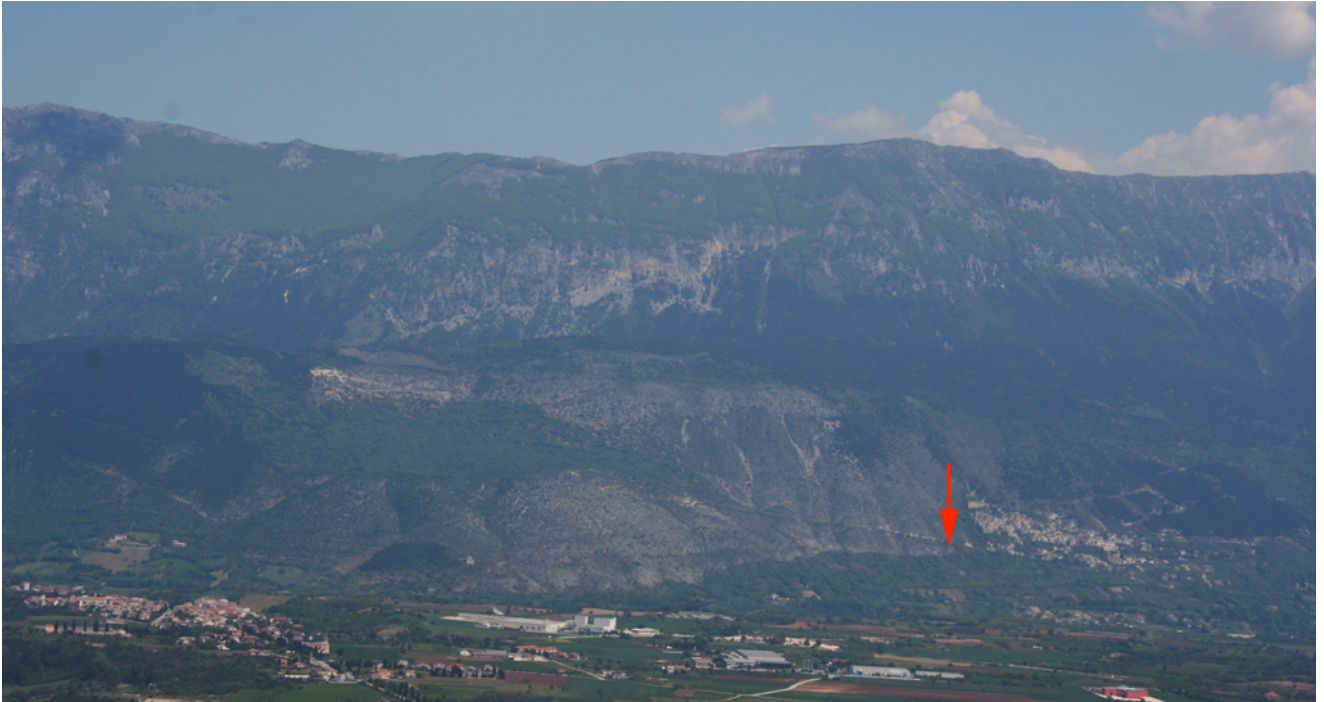
## Faille de Roccacasale - site RC

*Mission Mai 2007*

**GPS** : alt. 392 m  $\pm$  5 m ; N42.12515 - E013.88250

**Ecrantage** : seulement le plan de faille **120 S 55**

**Nbre échantillons** : 47 (RC-1 à RC-47)





SITE	nom blk.	atom Clnat	atom 36Cl	spike name
<b>MA3</b> Palumbo et al 2004 (M1)	<b>B3MAG</b> 01	5,65E+17	1,40E+05	13/02/2003
	<b>B4MAG</b> 01	8,44E+17	2,12E+05	13/02/2003
	<b>B6MAG</b> 01	5,86E+17	1,36E+05	13/02/2003
	<b>B7MAG</b> 01	8,58E+17	4,05E+05	13/02/2003
	<i>moy MA3 =</i>	<b>7,13E+17</b>	<b>2,23E+05</b>	
<b>MA3c</b> (M1c)	BLK M1c 1-20	2,71E+17	1,92E+05	25/09/2007
	BLK M1c 21-42	2,67E+17	1,66E+05	25/09/2007
	<i>moy MA3c =</i>	<b>2,69E+17</b>	<b>1,79E+05</b>	
<b>MA1</b>      (A1c)	A1c-BL3	6,55E+17	1,72E+05	13/01/2006
	AG-BL1	2,42E+17	3,09E+05	14/02/2006
	BLK 22/02/07	4,76E+17	3,67E+05	30/11/2006
	BLK 28/03/07	2,43E+17	2,54E+05	30/11/2006
	BLK 07/03/07	2,79E+17	2,33E+05	30/11/2006
	BLK A1C-2K/54TK	2,28E+17	1,79E+05	25/09/2007
	BLK A1C-42/70	2,76E+17	2,98E+05	30/11/2006
	<i>moy MA1 =</i>	<b>3,43E+17</b>	<b>2,59E+05</b>	
<b>MA4</b>    (A4c)	A4c-BL 1	6,05E+17	3,89E+05	13/01/2006
	BLK 22/02/07	4,76E+17	3,67E+05	30/11/2006
	BLK 07/03/07	2,79E+17	2,33E+05	30/11/2006
	BLK mai 08	1,38E+17	2,13E+05	28/04/2008
	<i>moy MA4 =</i>	<b>3,74E+17</b>	<b>3,01E+05</b>	
<b>MA2</b>      (A6c)	AG-BL1	2,42E+17	3,09E+05	14/02/2006
	BLK A6c-3/12	2,44E+17	4,48E+05	27/04/2007
	BLK A6c-31/47	2,25E+17	4,00E+05	27/04/2007
	BLK A6c-28/03/07	2,43E+17	2,54E+05	30/11/2006
	BLK A6C-64/73	2,68E+17	2,38E+05	30/11/2006
	BLK A6C-13/30	9,23E+16	4,27E+05	27/04/2007
	<i>moy MA2 =</i>	<b>2,19E+17</b>	<b>3,46E+05</b>	
<b>VE</b>	BLK VE42-61	2,42E+17	1,73E+05	25/09/2007
	BLK VE1-20	2,74E+17	1,97E+05	25/09/2007
	BLK VE-21/41	4,03E+17	1,71E+05	25/09/2007
	<i>moy VE =</i>	<b>3,06E+17</b>	<b>1,81E+05</b>	
<b>TR</b>	BLK TR-21/40	2,15E+17	1,87E+05	10/12/2007
	BLK TR-01/20	1,63E+17	2,01E+05	10/12/2007
	<i>moy TR =</i>	<b>1,89E+17</b>	<b>1,94E+05</b>	

Récapitulatif des valeurs des blancs de chimie sur les sites. Entre parenthèse est indiqué l'ancien nom du site lorsqu'il a été simplifié dans le manuscrit.

nom ech.	[Cl]nat	[36Cl]	incertitude sur [36Cl]	uncert. %
pas de réplikat sur MA3				
<b>MA3c</b> (M1c)				
M1C-18	11,120	1,40E+05	3,34E+03	
M1C-18R	11,364	1,45E+05	3,59E+03	
<b>M1C-18 mean</b>	<b>11,242</b>	<b>1,43E+05</b>	<b>2,79E+03</b>	1,96
M1C-28	12,908	1,25E+05	3,14E+03	
M1C-28R	13,074	1,27E+05	3,76E+03	
<b>M1C-28 mean</b>	<b>12,991</b>	<b>1,26E+05</b>	<b>1,02E+03</b>	0,81
M1C-42	10,833	1,05E+05	2,61E+03	
M1C-42R	11,077	1,05E+05	3,20E+03	
<b>M1C-42 mean</b>	<b>10,878</b>	<b>1,05E+05</b>	<b>3,64E+02</b>	0,35
<b>MA1</b> (A1c)				
A1c-54TK	0,887	2,80E+05	8,90E+03	
A1c-54T	0,788	2,77E+05	7,32E+03	
<b>A1c-54T mean</b>	<b>0,837</b>	<b>2,79E+05</b>	<b>1,77E+03</b>	0,63
A1c-54sinK	1,237	3,00E+05	8,35E+03	
A1c-54sin	1,129	2,85E+05	1,45E+04	
<b>A1c-54sin mean</b>	<b>1,183</b>	<b>2,92E+05</b>	<b>7,47E+03</b>	2,56
A1c-1	0,994	5,68E+05	8,05E+02	
A1c-1R	1,071	5,52E+05	1,00E+04	
<b>A1c-1 mean</b>	<b>1,032</b>	<b>5,60E+05</b>	<b>7,98E+03</b>	1,42
A1c-8	1,357	5,09E+05	1,00E+04	
A1c-8K	0,844	5,01E+05	9,25E+03	
<b>A1c-8 mean</b>	<b>1,100</b>	<b>5,05E+05</b>	<b>3,58E+03</b>	0,71
A1c-9	0,870	4,65E+05	8,28E+03	
A1c-9R	0,413	4,56E+05	1,59E+04	
<b>A1c-9 mean</b>	<b>0,642</b>	<b>4,60E+05</b>	<b>4,47E+03</b>	0,97
A1c-10	0,792	4,74E+05	1,12E+04	
A1c-10R	0,364	4,51E+05	1,16E+04	
<b>A1c-10 mean</b>	<b>0,578</b>	<b>4,62E+05</b>	<b>1,15E+04</b>	2,49
A1c-12	0,812	4,33E+05	2,04E+04	
A1c-12K	0,710	4,40E+05	1,59E+04	
<b>A1c-12 mean</b>	<b>0,761</b>	<b>4,37E+05</b>	<b>3,57E+03</b>	0,82
A1c-20	0,787	4,48E+05	1,69E+04	
A1c-20K	0,590	3,94E+05	1,04E+04	
<b>A1c-20 mean</b>	<b>0,688</b>	<b>4,21E+05</b>	<b>2,67E+04</b>	6,34
A1c-28	0,848	3,52E+05	8,06E+03	
A1c-28K	0,844	3,58E+05	9,63E+03	
<b>A1c-28 mean</b>	<b>0,846</b>	<b>3,55E+05</b>	<b>3,04E+03</b>	0,86
A1c-24	2,105	4,04E+05	1,15E+04	
A1c-24K	1,028	3,93E+05	8,06E+03	
<b>A1c-24 mean</b>	<b>1,566</b>	<b>3,98E+05</b>	<b>5,29E+03</b>	1,33
A1c-32	1,442	3,80E+05	9,79E+03	
A1c-32K	1,381	3,57E+05	7,42E+03	
<b>A1c-32 mean</b>	<b>1,412</b>	<b>3,69E+05</b>	<b>1,14E+04</b>	3,10
A1c-38	0,380	3,43E+05	6,58E+03	
A1c-38R	0,646	3,55E+05	1,15E+04	
<b>A1c-38 mean</b>	<b>0,513</b>	<b>3,49E+05</b>	<b>6,13E+03</b>	1,76
A1c-44	0,208	3,16E+05	9,35E+03	
A1c-44K	0,407	3,20E+05	7,50E+03	
<b>A1c-44 mean</b>	<b>0,307</b>	<b>3,18E+05</b>	<b>1,58E+03</b>	0,50
A1c-69	3,042	2,59E+05	4,70E+03	
A1c-69R	3,072	2,78E+05	5,57E+03	
<b>A1c-69 mean</b>	<b>3,057</b>	<b>2,69E+05</b>	<b>9,50E+03</b>	3,54
<b>MA4</b> (A4c)				
A4c-3	8,925	1,01E+06	1,82E+04	
A4c-3R	8,196	9,79E+05	1,74E+04	
<b>A4c-3 mean</b>	<b>8,561</b>	<b>9,94E+05</b>	<b>1,45E+04</b>	1,46
A4c-63	7,508	3,70E+05	6,53E+03	
A4c-63R	7,588	3,71E+05	6,55E+03	
<b>A4c-63 mean</b>	<b>7,548</b>	<b>3,70E+05</b>	<b>5,12E+02</b>	0,14
A4c-10	9,080	9,26E+05	2,19E+04	
A4c-10R	8,783	9,39E+05	2,86E+04	
<b>A4c-10 mean</b>	<b>8,932</b>	<b>9,32E+05</b>	<b>6,31E+03</b>	0,68
A4c-56	2,675	3,83E+05	9,12E+03	
A4c-56R	2,877	3,97E+05	7,67E+03	
<b>A4c-56 mean</b>	<b>2,776</b>	<b>3,90E+05</b>	<b>6,92E+03</b>	1,77
A4c-43	12,758	5,18E+05	9,48E+03	
A4c-43R	13,002	5,41E+05	9,70E+03	
<b>A4c-43 mean</b>	<b>12,880</b>	<b>5,30E+05</b>	<b>1,20E+04</b>	2,26
A4c-35	5,329	6,02E+05	1,10E+04	
A4c-35R	5,381	6,15E+05	1,77E+04	
<b>A4c-35 mean</b>	<b>5,355</b>	<b>6,08E+05</b>	<b>6,34E+03</b>	1,04
A4c-11	9,198	9,23E+05	1,64E+04	
A4c-11R	8,200	9,33E+05	1,71E+04	
<b>A4c-11 mean</b>	<b>8,699</b>	<b>9,28E+05</b>	<b>5,10E+03</b>	0,55
<b>MA2</b> (A6c)				
A6c-72	4,598	3,45E+05	9,91E+03	
A6c-72R	4,809	3,75E+05	1,30E+04	
<b>A6c-72 mean</b>	<b>4,704</b>	<b>3,60E+05</b>	<b>1,50E+04</b>	4,16
A6c-42	9,780	6,25E+05	1,51E+04	
A6c-42R	10,440	6,72E+05	1,66E+04	
<b>A6c-42 mean</b>	<b>10,110</b>	<b>6,49E+05</b>	<b>2,32E+04</b>	3,58
A6c-74	5,687	3,73E+05	8,78E+03	
A6c-74R	5,299	3,46E+05	8,20E+03	
<b>A6c-74 mean</b>	<b>5,493</b>	<b>3,60E+05</b>	<b>1,33E+04</b>	3,71
A6c-30	8,196	7,30E+05	1,77E+04	
A6c-30R	8,536	7,53E+05	1,04E+04	
<b>A6c-30 mean</b>	<b>8,366</b>	<b>7,41E+05</b>	<b>1,16E+04</b>	1,56
A6c-01	7,380	1,15E+06	2,73E+04	
A6c-01R	7,126	1,07E+06	1,92E+04	
<b>A6c-01 mean</b>	<b>7,253</b>	<b>1,11E+06</b>	<b>4,21E+04</b>	3,79
<b>VE</b>				
VE-77	5,121	2,49E+05	6,59E+03	
VE-77R	5,212	2,38E+05	6,55E+03	
<b>VE-77 mean</b>	<b>5,167</b>	<b>2,43E+05</b>	<b>5,65E+03</b>	2,32
VE-29	3,658	3,69E+05	8,49E+03	
VE-29R	3,818	3,88E+05	8,48E+03	
<b>VE-29 mean</b>	<b>3,738</b>	<b>3,79E+05</b>	<b>9,39E+03</b>	2,48

Récapitulatif des réplikats de chimie. Entre parenthèse est indiqué l'ancien nom du site lorsqu'il a été simplifié dans le manuscrit.





46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65	66
Fe2O3 (Fe)	MnO (Mn)	MpO (Mg)	CaO (Ca)	Na2O(Na)	K2O(K)	TO2(Ti)	P2O5(P)	B	Li	H2Otot (H)	Stot (S)	CO2 tot (C)	O.rock	O.water	Cl nat AMS	Ca ICP	Z (position sur le scarp)	Thickness sample	[36C] AMS	uncertainty
%	%	%	%	%	%	%	%	ppm	ppm	%	%	%			ppm	ppm	cm	cm	µg/rock	
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.25	380626	3	3.7	243273	6010
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.25	383741	10	3.5	250243	4421
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.65	394666	20	3.5	267334	6410
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.67	409582	30	3.2	258838	6595
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.84	380814	40	3	262230	6829
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.93	392767	60	3.3	271365	6290
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.9	381115	80	3.3	276166	6354
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.55	383735	100	3.3	281508	6487
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.85	386680	120	3.5	290359	6674
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.68	384921	130	3.5	294968	5212
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.22	371727	141	3.5	298124	5603
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	2.79	384310	151	3.5	300051	6420
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	2.78	387039	161	3	290094	6883
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.57	370577	181	3.5	285949	6597
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.03	353867	191	3	297066	7935
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.08	378063	201	3.3	312197	7541
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.05	379182	221	3	305699	7552
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.97	372668	231	3	306526	7696
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.28	376988	251	3.5	305256	7690
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.78	389280	261	3	331975	8904
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.2	388380	261	2.5	323813	9941
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.9	382178	281	3	336063	9442
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.92	384874	291	3	330886	6788
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.89	381503	302	3	349707	11755
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.22	387857	312	3	323695	6884
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.15	371541	322	2.8	345849	8026
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.09	375201	332	3	342156	6988
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.24	395587	342	3	355592	9715
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.6	380426	352	2.8	347228	8518
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.01	386504	362	2	344035	7972
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.5	388286	372	3	345148	8575
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.7	390236	381	3.3	373889	6640
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	5.5	383682	391	2.8	373201	9901
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.7	384505	401	3.3	371358	9477
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.16	378960	411	3	367715	8465
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.65	378960	431	3	374392	8614
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	4.61	416000	441	2.8	374608	12145
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.53	393748	451	3.3	357616	8245
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.15	395353	461	3	367941	6624
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	3.89	392975	471	3	388659	22231
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.75	396263	482	2.5	390576	6897
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.52	373204	492	3	414749	9545
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.71	380701	501	2.7	393725	8366
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	8.14	382472	511	3	408083	9435
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.48	386319	521	2.7	409963	7249
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.27	388477	531	2.5	427419	10200
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.79	404865	541	2.5	441048	13578
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	8.44	369564	551	3	410586	5283
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	8.07	40728	571	3	407610	1659
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	8.07	40728	581	3	469192	1847
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.08	402948	581	3	438294	8347
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	6.55	403989	591	3	449818	3062
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.41	398719	601	3	453718	8701
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.89	398356	611	3	486398	11301
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	7.99	403535	630	2	499714	12261
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	8.96	394338	640	2.5	484225	9261
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	9.87	401307	650	2.5	512268	12969
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	11.72	393725	671	3	542468	9575
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	14.6	400025	681	2.8	498289	9529
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	12.69	392300	691	2.7	577073	10079
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0	0.0	11.81	400376	701	2.5	527627	15662
0.07	0.0	0.65	55.75	0.0	0.0	0.01	0.0	2.0	0.99	0.29	0.02	43.77	0.0							





46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65	66	
Fe2O3 (Fb) %	MnO (Mn) %	MgO (Mg) %	C-nb (Cn) %	Nu2ONa) %	K2O(K) %	TO2(T) %	P2O5(P) %	B ppm	Li ppm	H2O(H) %	Stc(S) %	CO2 tot(C) %	O.rock	O.water	Cl mix AMS ppm	Ca ICP ppm	Z (position sur le stemp) cm	Thickness sample cm	13C(AMS) aff.rock	uncertainty	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.18	36.5564					
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	2.08	36.792	17	3	2570.95	589	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	3.06	36.8554	37	3	2363.95	595	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	2.45	37.8742	37	3	2603.54	6033	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	2.91	36.4440	57	2	2663.92	8127	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	3.14	36.5564	67	1	2745.33	4884	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	3.62	36.3306	78	2	2765.25	5404	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	3.74	36.5564	88	2.5	2889.94	5983	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	4.46	36.5564	100	2	2863.62	5985	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.52	36.5564	110	2	2863.62	5985	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.34	36.792	120	2	2909.98	10848	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.26	36.5564	140	2.5	2897.89	10182	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.96	36.5564	150	1	3010.92	5938	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.84	36.9673	160	3.5	3009.73	7962	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.79	36.9673	170	0.3	2982.36	7763	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.64	36.5564	180	2	2928.69	7778	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	3.54	36.5564	190	2.5	2928.69	5188	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	2.19	36.2231	200	3.5	2975.52	70547	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.66	36.5564	209	4	2975.52	70547	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.17	40.3358	222	3	311247	6137	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.23	36.5564	232	3	305190	7048	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.16	36.5564	232	2	320141	7889	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.26	36.5564	242	3	320141	7889	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.26	36.9721	262	3	314392	7247	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.26	36.9721	262	3	314392	7247	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.31	37.5844	283	1	317951	1583	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.23	36.6548	303	3	331280	7654	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.17	36.5564	313	1	330562	6517	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.29	36.5564	323	2	338860	15098	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.11	36.5564	323	2.5	338860	15098	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.21	36.5564	353	1	345434	7124	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.38	36.5564	363	2	353964	9539	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.38	36.5564	384	2	356428	14767	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.59	36.5564	384	2	356428	14767	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.4	36.5564	394	1.5	371107	17313	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.41	36.3357	404	2	365559	14338	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.78	36.5564	404	2	365559	14338	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.64	36.5564	416	2	365559	14338	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.64	36.5564	416	2	365559	14338	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.85	37.4550	443	2	353416	7750	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.22	36.5564	443	2	353416	7750	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.09	36.5564	463	1.5	363379	14292	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.57	36.5564	474	1.5	375265	7398	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.81	37.3821	484	2	398214	5291	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.78	36.9743	484	2	398214	5291	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.81	36.5564	494	2.5	398180	9877	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.38	36.5564	515	2	406517	14161	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.69	36.6455	545	2	406517	14161	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.14	36.6455	545	2	421208	28720	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.62	40.0444	555	1.5	403828	10468	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.92	36.5564	565	2.5	421259	8880	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	1.41	36.5564	578	1	433112	7700	
0.0	0.001	0.549	54.818	0.0	0.0	0.001	0.007	1.0	0.811	0.56	0.016	44.249	0.0	0.0	0.81	36.5564	607	2	447781	19437	
0.0	0.001																				







46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65	66	
Fe2O3 (P) %	MnO (Mn) %	MnO (Mn) %	CaO (Ca) %	Na2O (Na) %	K2O (K) %	TiO2 (Ti) %	P2O5 (P) %	B ppm	Li ppm	H2Otot (H) %	Stot (S) %	CO2 tot (C) %	O.rock	O.water	Cl nat. AMS ppm	Ca (CP ppm)	Z (position sur le scarp) cm	Thickness sample cm	[86C] AMS at/g.rock	uncertainty	
0	0	0,001	0,32718182	55,838	0	0	0	0,05590909	2	0,89090909	0,19454545	0,01	43,8681818								
0	0	0	0,33	55,84	0	0	0	0,06	2	0,89	0,19	0,01	43,87	0,0							
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	381823	5,49	5 3,5	359704	13345		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	380816	6,41	14 4,25	356658	6284		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	384881	4,70	24 3,5	359892	14972		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	389379	8,21	34 3,3	371298	10740		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	379360	8,78	44 3,4	375234	6624		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	384046	9,03	54 3,4	398267	6624		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	383228	9,26	64 3,4	399308	9504		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	375275	8,26	74 3,8	398393	9494		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	379661	7,75	84 3,65	413779	7311		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	379322	8,08	92	433255	10390		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	378731	11,15	125 3,1	433255	10390		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	382925	10,44	136 3,5	453813	10934		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	386641	9,21	146 3,5	468026	11260		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	386541	8,81	155 3,5	468026	11260		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	386952	9,67	165 3,5	512128	9123		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	385680	9,53	175 2,75	512128	9123		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	387087	10,99	185 3,6	479717	11514		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	395687	11,22	196 3,4	512206	9046		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	382764	11,42	207 3,4	532539	14456		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	385473	11,63	217 3,5	528552	11608		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	396361	10,41	226 3,4	528552	12747		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	396361	10,41	236 3,2	540116	12038		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	380470	11,66	245 2,85	557177	10068		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	386157	12	256 2,65	557177	10068		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	395563	10,99	266 2,9	569553	10144		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	395057	10,91	277 2,9	561389	10130		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	385716	10,42	287 2,9	562081	13643		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	389388	8,95	297	630573	9912		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	393149	10,44	307 3,35	591916	39163		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	404045	10,41	317 2,65	594991	10748		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	404045	10,41	327 3,25	619496	17117		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	392391	10,11	337 3,1	648573	23214		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	396072	9,46	357 3,1	648573	23214		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	389325	10,04	367	681796	12330		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	397043	11,77	378	663874	14606		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	397483	8,36	386	679544	12108		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	387210	7,3	455 2,75	708165	23117		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	387383	7,03	475	689464	12120		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	382645	8,71	486 3,1	688784	22969		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	378415	8,44	496 3,2	693082	21481		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	398794	8,87	506 2,75	727787	12904		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	398794	8,87	516 2,9	705561	17098		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	388753	8,37	536	741471	11599		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	395904	9,55	546 2,75	748226	12700		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	371766	8,36	556	732543	17678		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	397560	9,11	566 3,7	738176	12510		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	375331	11,25	576 2,7	773378	15170		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	405643	9,86	586 3,5	790232	18970		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	382459	9,94	596	798026	13255		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	386424	9,15	606	788667	13255		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	384482	8,31	616	770119	27158		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	383968	9,76	636 3,2	804358	18465		
0,0	0,0	0,0	55,84	0,0	0,0	0,0	0,06	2	0,89	0,19	0,01	43,87	0,0	0,0	391639	10,02	645 3,3	868283	18662		
0,0	0,0	0,0	55,																		





46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65	66
Fe2O3 (Fs)	MnO (Mn)	MgO (Mg)	CaO (Ca)	Nu2O(Na)	K2O(K)	TiO2(Ti)	P2O5(P)	B	LI	H2Otot(H)	Stc (S)	CO2 tot(C)	O.rock	O.water	Clntr AMS	Ca ICP	Z (position sur le scarp)	Thickness sample	36(C) AMS	uncertainty
%	%	%	%	%	%	%	%	ppm	ppm	%	%	%			ppm	cm	cm	atg/rock		
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	7.79	391559	5	3	377556	8981
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	7.55	395144	15	2.5	370170	512
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	7.14	391267	25	2.5	359016	6589
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.49	395144	35	3	363426	6419
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.59	401105	45	2	380303	6286
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.12	395144	55	2	384483	8773
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.88	382551	65	2	397151	7485
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	3.76	395144	75	1.5	385888	11305
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	2.78	393423	85	2	390038	6917
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	3.61	395144	95	3	401879	7101
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	3.31	397547	105	1.5	417004	9694
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	3.96	402284	115	1.5	402284	13206
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	3.75	408697	125	1.5	402284	13206
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	3.75	386201	145	2	422901	7759
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.74	395144	155	2.5	427983	9074
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.07	394891	164	2	424316	7783
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.53	395144	174	1	438786	8041
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.04	397985	184	1.5	472253	8895
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	11.86	395144	192	3	512394	18027
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	12.18	392703	199	3	536860	9643
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	12.88	398477	209	3.5	529542	11952
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.71	391866	220	3	541076	12872
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.38	393435	241	3	535670	12707
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.68	395144	250	2	532039	9750
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.48	395901	260	2	536074	14736
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.7	391267	269	2	538244	14867
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.7	391267	269	2	538244	14867
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.35	393182	305	2	608387	6339
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.39	395144	315	2.5	608387	6339
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.74	403542	334	3	640950	19535
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.39	395144	345	2	655642	12008
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	7.76	395298	355	2.5	663495	18542
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.18	402284	364	3	677348	13888
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.18	393873	375	2	656380	16174
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.46	395144	385	2	706197	18673
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	4.99	397585	395	2	715888	17635
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.91	397243	405	2	722763	13060
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.56	400169	415	2.5	704679	30847
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.75	395144	425	2.5	750103	15060
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	12.52	397293	437	2.5	718194	15177
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	12.52	397293	437	2.5	718194	15177
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	13.35	397636	496	1.5	818081	21137
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	10.87	395144	506	1.5	835546	20526
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.92	386504	516	1.5	847910	15669
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.51	395144	526	2	799197	19674
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.66	393976	536	2	869821	15509
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.75	395144	546	3	831017	13686
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	11.64	399993	556	2.5	871769	15527
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	10.69	395144	566	1.5	891727	14788
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.7	393916	576	2.5	904663	16158
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.7	393916	586	3	895470	26214
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.93	392005	596	3	927772	5101
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.93	392005	596	3	932388	6306
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.93	392005	607	3	931947	16617
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	10.07	395144	617	3	931947	16617
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	10.07	395144	627	3	931947	16617
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	9.12	396990	637	2	970786	29468
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.16	395144	647	1.5	996097	12524
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.16	391629	670	3	979695	17741
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	8.56	395144	677	2.5	1006405	14351
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	6.79	393071	687	2	993532	14522
0.0	0.0	0.45	54.94	0.0	0.0	0.0	0.03	1.0	0.88	0.4	0.01	44.08	0.0	0.0	5.3	395144	697	2	1025517	14587
0.0	0.0	0.45	54.94	0.0																







46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65	66
Fe <sub>2</sub> O <sub>3</sub> (Fe)	MnO (Mn)	MgO (Mg)	CaO (Ca)	Na <sub>2</sub> O (Na)	K <sub>2</sub> O (K)	TiO <sub>2</sub> (Ti)	P <sub>2</sub> O <sub>5</sub> (P)	B	Li	H <sub>2</sub> O tot (H)	Stat (S)	CO <sub>2</sub> tot (C)	O <sub>2</sub> rock	O <sub>2</sub> water	Cl nat AMS	Ca IP	Z (position sur le scarp)	Thickness sample	(36C) AMS	uncertainty
%	%	%	%	%	%	%	%	ppm	ppm	%	%	%			ppm	ppm	cm	cm	μg/rock	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	7.34	366827	5	3.7	358904	9421	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	6.69	368709	15	4	359660	6486	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	6.38	368804	25	4	384624	6826	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	7.3	369164	35	4	395363	6983	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.35	355697	45	3.7	383815	6830	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.97	378671	55	3.8	404423	6918	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	12.34	370947	65	3.5	425856	10564	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.62	374437	75	3.8	415273	7916	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	13.1	376177	85	3.7	425446	7627	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	18.34	372229	95	3.5	445904	7876	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	12.61	372731	105	3.5	454223	8059	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.85	365761	115	3	454195	8680	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	9.33	366696	125	3.7	458355	8142	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.64	363044	135	3.6	478934	8520	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.85	367403	145	3	498988	8857	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.28	375245	155	3.8	487742	8700	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	6.93	369284	165	3.5	520171	8232	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	5.21	394897	175	3.7	520211	8520	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	4.62	393961	185	3.5	595972	10958	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.92	393960	195	3.5	595972	10958	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.93	393961	205	3.5	595972	10958	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.81	378671	215	3.7	598291	11269	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.81	378671	225	3.7	609064	9975	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.65	393222	235	3.8	648315	12130	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	13.04	390338	245	3.5	648250	11616	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	12.26	397511	255	3.5	651526	12246	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.66	373097	265	3.7	666056	12292	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	14.32	371884	275	3.5	739234	13861	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.84	375630	285	3.5	748831	14432	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.54	400978	295	3.8	752217	21241	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	7.46	400978	305	3.8	757892	16895	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	8.7	379651	315	3.8	801066	15204	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.76	391492	325	3.8	762438	16859	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.83	392962	335	3.8	810591	14877	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	9.99	395543	345	3.7	807493	13207	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.25	383233	355	4	861818	15285	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.09	391954	365	3.5	890517	15736	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.72	388740	375	3.9	957382	17005	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	10.78	388740	385	3.9	984762	40877	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.14	381939	395	4	1027721	15872	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.54	389659	405	3.7	1051175	18563	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.35	383150	415	3.8	1089735	16699	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	12.69	385880	425	3.7	1154112	20354	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	11.32	378671	435	3.7	1187118	18263	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	12.5	378671	445	3.5	130444	23047	
0.0	0.0	0.32	54.9	0.0	0.0	0.0	0.02	2.035	0.28	0.01	44.04	0.0	0.0	12.69	381358	455	3.5	134185	20463	
0.0	0.0	0.46	55.87				0.0	2	1.022	0.01	44.62	0.0	0.0	9.74	347161	1300	2.5	1890688	33373	

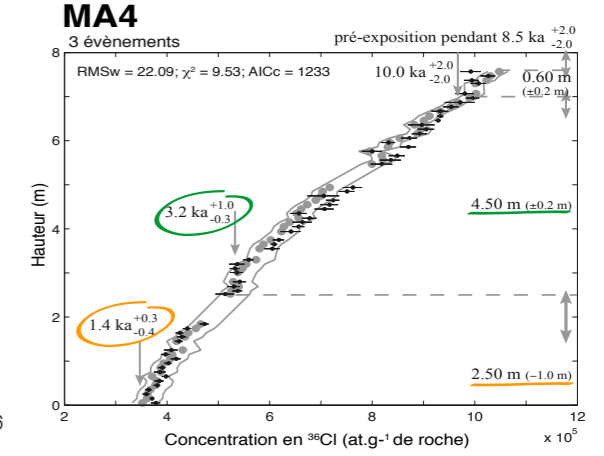
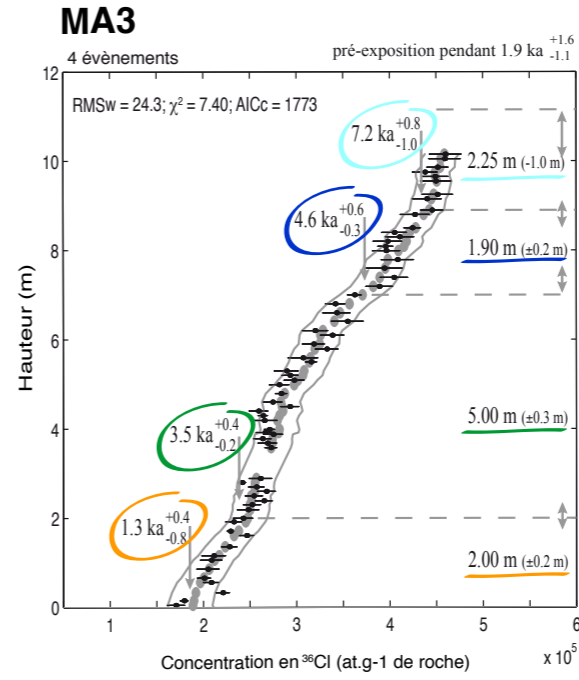
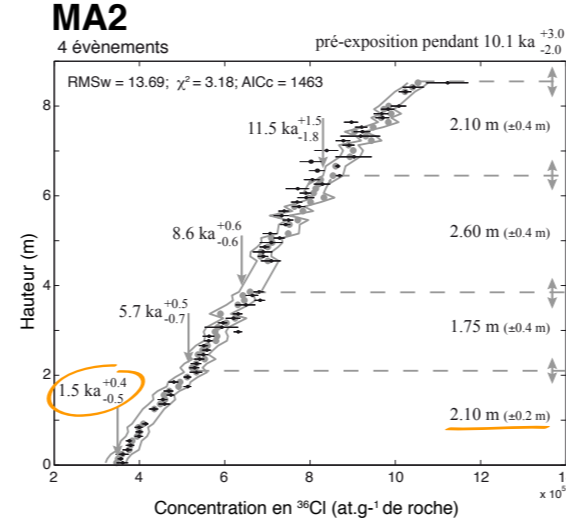
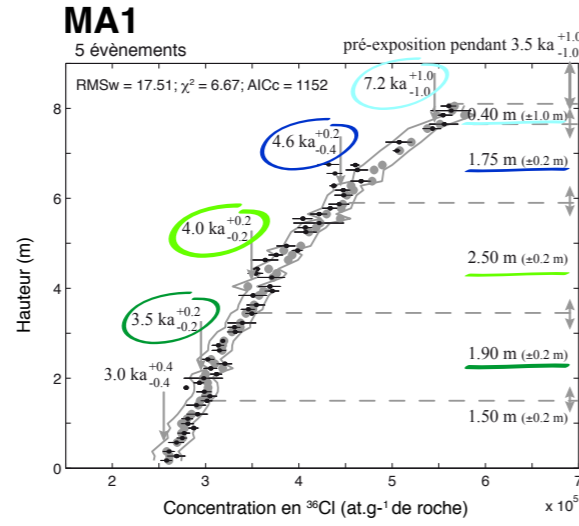
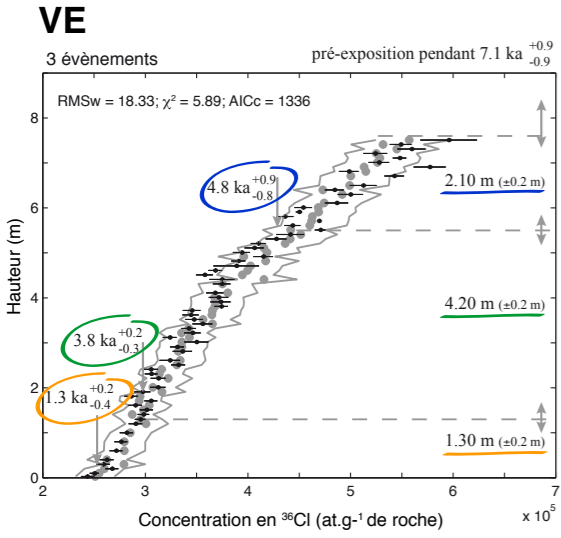
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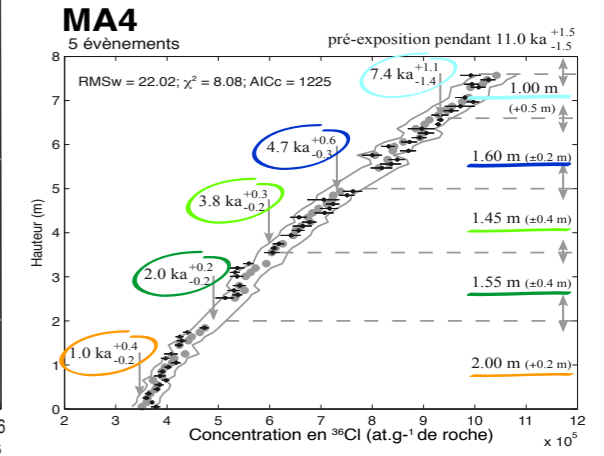
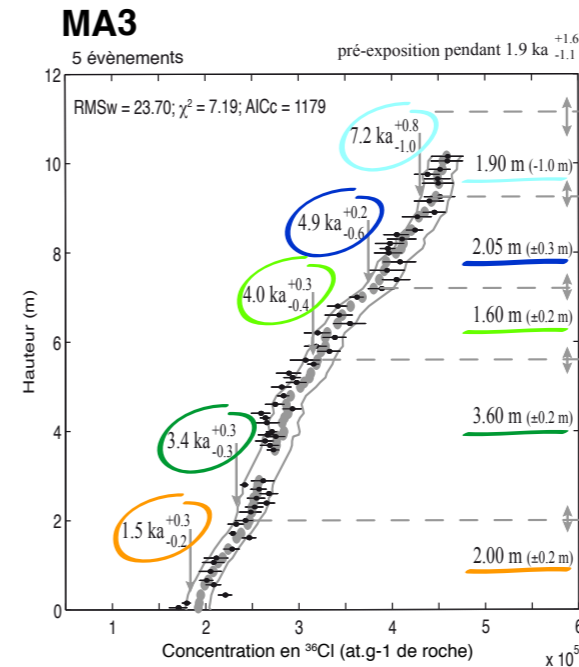
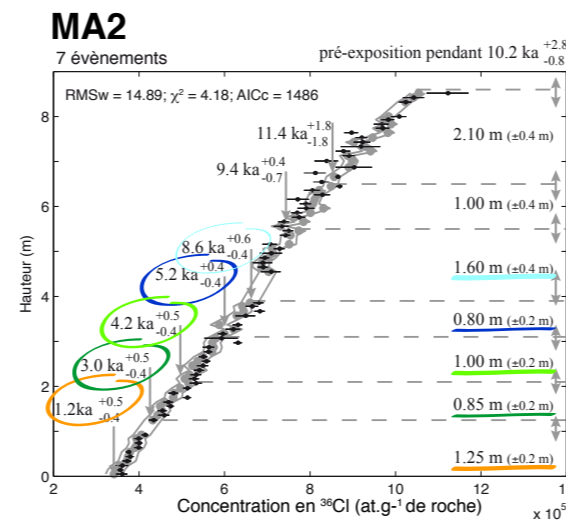
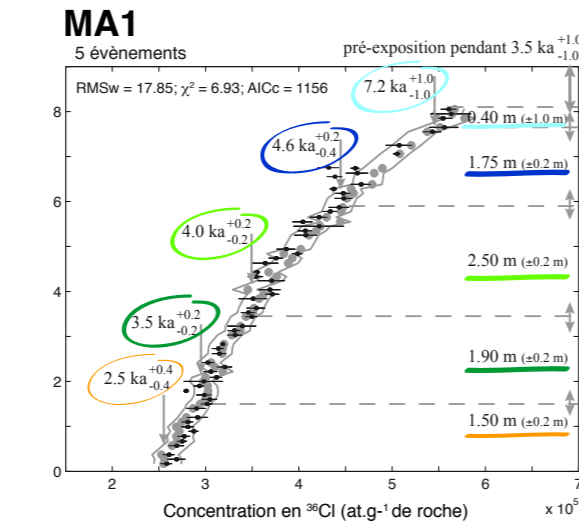
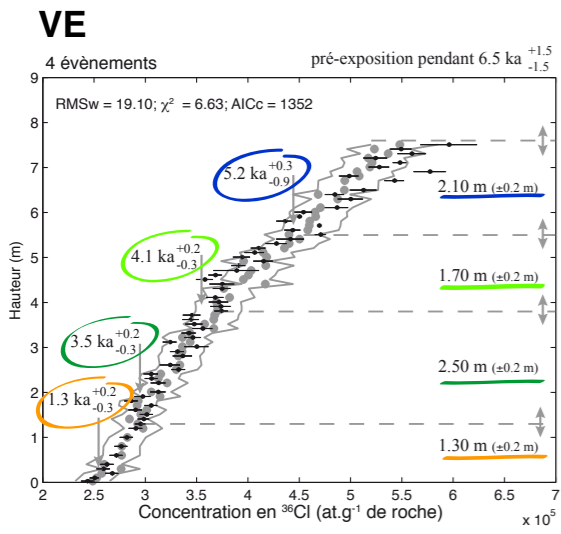


"modèles<sub>pref</sub>"



- n°5 [6.2 - 8.2] ka
- n°4 [4.0 - 5.0] ka
- n°3 ~4.0 ka
- n°2 [3.0 - 4.2] ka
- n°1 [0.5 - 2.0] ka

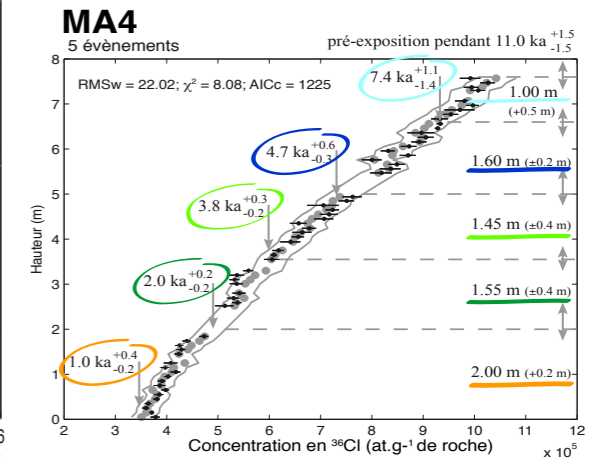
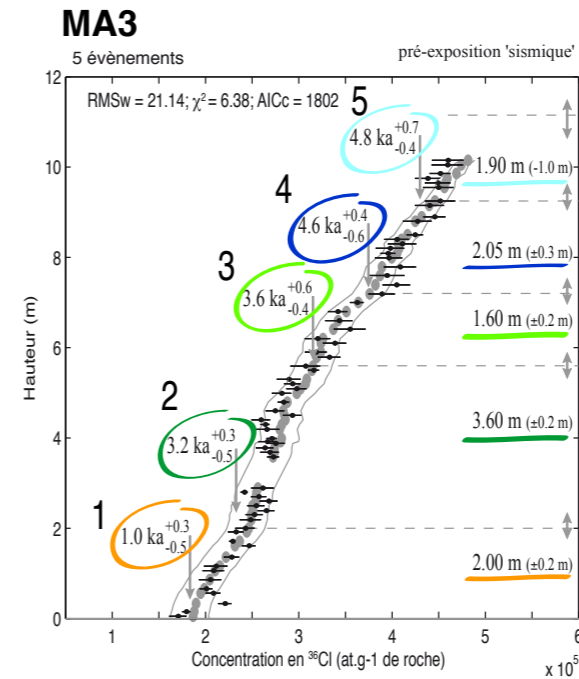
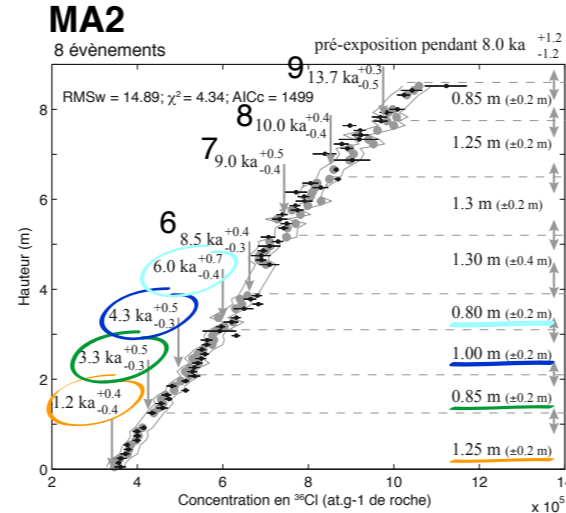
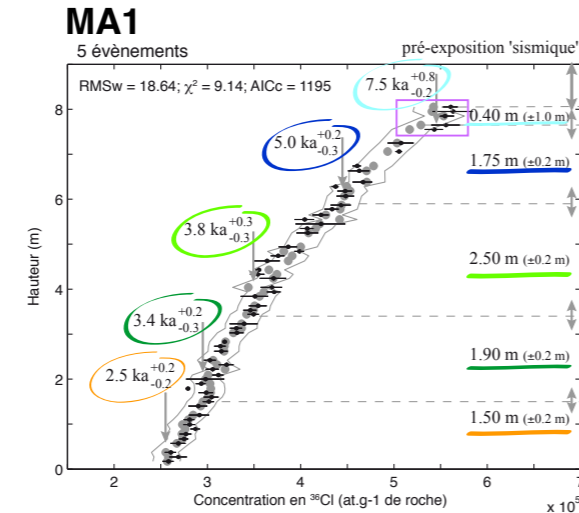
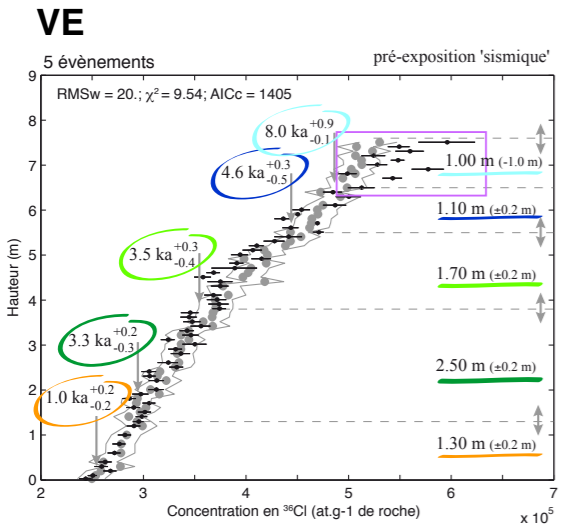
"modèles<sub>opt</sub>"



- n°5 7.6 (± 1.6) ka
- n°4 4.9 (± 0.7) ka
- n°3 4.0 (± 0.7) ka
- n°2 3.3 (± 0.7) ka
- n°1 1.2 (± 0.5) ka

"modèles<sub>max</sub>"

Mauvais ajustement des données avec le modèle pré-exposition 'sismique'



- n°5 7.2 (± 1.7) ka
- n°4 4.6 (± 0.7) ka
- n°3 3.7 (± 0.6) ka
- n°2 3.3 (± 0.6) ka
- n°1 1.0 (± 0.5) ka



# **Annexe IV**

(du chapitre 6)







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# Incremental growth of normal faults: Insights from a laser-equipped analog experiment

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### ABSTRACT

We conducted a laser-equipped analog experiment aimed at quasi-continuously monitoring the growth of a dense population of normal faults in homogeneous conditions. To further understand the way geological faults progressively gain in slip and length as they accumulate more strain, we measured with great precision the incremental slip and length changes that the analog faults sustain as they grow. These measurements show that the analog faults share common features with the natural ones. In particular, during their growth, the faults develop and maintain cumulative slip profiles that are generally triangular and asymmetric. The growth takes place through two distinct phases: an initial, short period of rapid lateral lengthening, followed by a longer phase of slip accumulation with little or no lengthening. The incremental slip is found to be highly variable in both space (along the faults) and time, resulting in variable slip rates. In particular, 'short- and long-term' slip rates are markedly different. We also find that slip measurements at local points on fault traces do not contain clear information on the slip increment repeat mode. Finally, while the fault growth process is highly heterogeneous when considered at the scale of a few slip events, it appears homogeneous and self-similar at longer time scales which integrate many slip increments. This is likely to be the result of a feedback between stress heterogeneities and slip development. The long-term scale homogeneity also implies that the long-term faulting process is primarily insensitive to the short-term heterogeneities that are rapidly smoothed or redistributed. We propose a new conceptual scenario of fault growth that integrates the above observations and we suggest that faults grow in a bimodal way as a result of a self-driven and self-sustaining process.

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## 1. Introduction

Faults generally grow, that is, accumulate slip and lengthen, through the repetition of earthquakes. Yet, how such a growth occurs is still unclear as we lack data documenting how earthquake slip events repeat on faults. Consequently, fundamental questions such as 'Do repeating earthquake slip events have similar amplitudes and lengths? Does any major slip event on a fault contribute to its further lengthening? How do repeating earthquake slips relate to cumulative displacements on long-term faults?' are still unanswered. We have only sparse information, arising partly from rare observations, and partly from models.

Paleoseismological data show that, in some cases, the paleo-slip events have followed at regular time intervals with similar amplitudes

and lengths (e.g., Rockwell et al., 2001; Klinger et al., 2003; Haibing et al., 2005). This suggests that faults may grow by accumulating constant slip without lengthening. Yet, in other cases, the paleo-earthquakes are found to have occurred at irregular time intervals and to have been dissimilar in size and location (e.g., Weldon et al., 2004; Palumbo et al., 2004; Daeron et al., 2007). This behavior is a sign of greater complexities in the fault growth process. On the other hand, the analyses of long-term faults whose cumulative displacements are the sum of hundreds to thousands of earthquakes (e.g., Taylor et al., 2004) generally provide little information on those past earthquakes. They thus fail to establish a link between repeat of earthquake slip events and fault growth.

Contrasting with the scarcity of data is our capacity of producing models and qualitative scenarios depicting how faults may grow. These are rooted in *a priori* hypotheses, such as faults behaving as simple, narrow, elastic cracks with no cohesion, linear elastic behavior of the crust, and elastic rebound theory (e.g., Scholz, 2002 for a review). The community interested in earthquakes provides numerical models that describe how earthquake slip events repeat on faults depending on the level of stress/strain being accumulated and relaxed (e.g. Shimazaki and Nakata, 1980), and on some other parameters such

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as friction on the fault plane (e.g. Tse and Rice, 1986; Rice and Ben-Zion, 1996). These models are useful as they can be directly compared with real earthquake slip data. By contrast, they are never compared with long-term slip measurements, so they do not provide any link between repeat of earthquake slip events and fault growth. On the other hand, the community more interested in overall fault behavior provides qualitative scenarios and more rarely numerical models that describe how faults may grow over long time spans ( $10^4$ – $10^7$  yr; Segall and Pollard, 1980; Walsh and Watterson, 1987, 1988; Peacock and Sanderson, 1991; Marrett and Allmendinger, 1991; Cowie and Scholz, 1992a,b; Bürgmann et al., 1994; Cartwright et al., 1995; Willemse et al., 1996; Cowie, 1998; Cowie and Roberts, 2001; Walsh et al., 2002; Kim and Sanderson, 2005). These scenarios are useful for comparison with real long-term slip measurements, but they are rarely confronted with earthquake slip data or with earthquake mechanics. They do therefore also fail to provide a link between repeat of earthquake slip events and fault growth.

Our objective is to explore an alternative field between observation of natural faults and developments of *ad hoc* models. We aim at examine and better understand how slip increments follow in space and time on a fault to make it grow. We address that topic through an analog experiment that allows for a large number of normal faults to form and grow under homogeneous conditions, while their surface lengths and slips are measured quasi- continuously with an interferometric laser system. Although this approach does not reproduce the earthquake process, it sheds light upon the way faults progressively gain in slip and length as they accommodate more strain. Note that, while many analog models have already been produced to simulate fault development (with a particular focus on normal faults; e.g. Horsfield, 1977; Faugeres and Brun, 1984;

Ackermann et al., 2001), with a few of them being monitored (Mansfield and Cartwright, 2001; Bellahsen et al., 2003; Marchal et al., 2003; Bellahsen and Daniel, 2005), this is the first time that the incremental fault growth is measured quasi-continuously, and at such a high resolution. This is also the first time the analysis is carried out with a sufficiently dense population of incremental slip events to identify common properties.

## 2. Experimental protocol

### 2.1. Experimental design and scaling

We study normal faults because their displacements can be easily measured at the model surface. Since we seek identifying general, intrinsic properties of the fault growth process, we need make the faults developing in the absence of any local complexity. Heterogeneous stress boundary conditions, variations of material properties, influence of preexisting discontinuities, etc... are as many factors – not systematic in nature – that may indeed modify the fault growth process (e.g. Bürgmann et al., 1994; Buitter and Schreurs, 2006). To avoid such possible effects, we have developed a new analog protocol that allows normal faults to be formed and grown in a homogeneous material and a homogeneous stress field. It is important to note that most available analog models do not fulfill the latter condition (details in supporting online material A). We adopt, with some modifications, the approach developed by Schreurs et al. (2002). The model is constructed on top of a thick foam bar that has been previously shortened by compression. The foam bar is then let to progressively ‘decompress’, which transfers a continuous, homogeneous extension over the width of the overlying model. For reasons explained in many

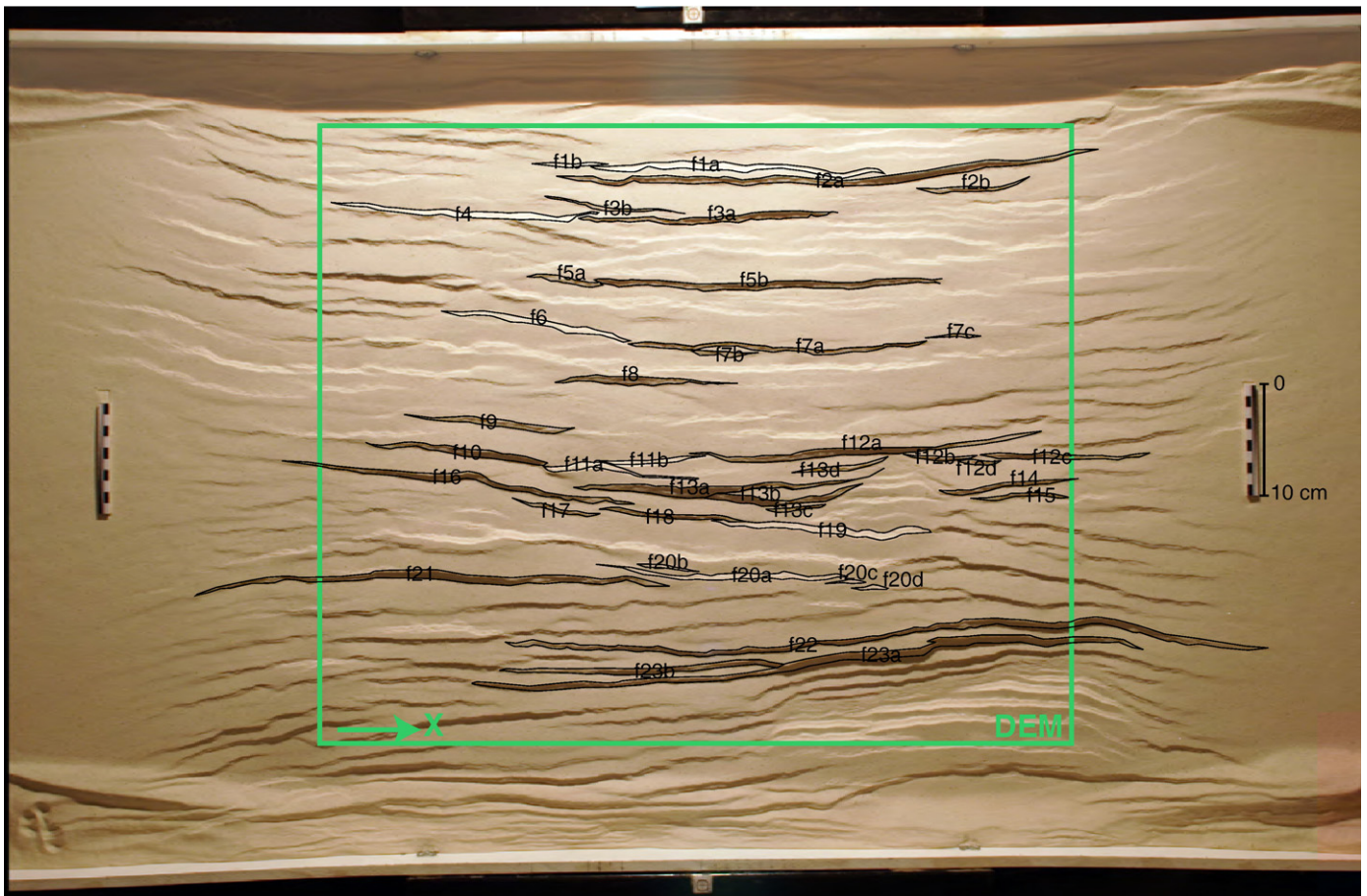


Fig. 1. Map view of the analog model at the end of the experiment. The frame is the zone under study. Analyzed faults are circled in black, with their names indicated. Synthetic and antithetic faults appear in different tones due to shading.

papers (e.g., Hubbert, 1937; Davy and Cobbold, 1991; Lallemand and Malavieille, 1992; Lohrman et al., 2003; Buiter and Schreurs, 2006), dry quartz sand is used to approximate the brittle crust. The one we use is 98.5% made of SiO<sub>2</sub>, has a mean density of 2.1, and a mean grain size of 112 μm. The scaling factor that results is 1–2·10<sup>-5</sup> (1 cm in the model represents 1–2 km in the field). Since we want to produce a large number of growing faults, the model does not include any basal material such as silicone putty that would make strain localizing onto a few large faults only (e.g. Davy et al., 1995; Bellahsen et al., 2003).

The apparatus is presented in supporting online material B. In a first step, the foam bar is compressed by 25% between two metallic jaws. Then an 8 cm-thick layer of sand is deposited on top of the compressed foam. The model is sprinkled with a thin powder (25% silica, 74% sand, and 1% black rubber dust; not a rheological interface) aimed at increasing its surface cohesion while allowing 'ground reference points' to be defined. Then the foam bar is progressively decompressed from operating back the two jaws; the decompression transfers extension to the overlying sand layer. The plastic-made sidewalls of the apparatus avoid longitudinal compression to develop in the sand cake, so that the strain field imposed at the base of the model is homogeneous in its central part (i.e., between the jaws), and decreases to zero on either side. Our study only focuses on the model central part (frame, Fig. 1).

Note that we have repeatedly performed the above experiment to ensure the strength of the results that we describe below. We have in

particular varied the thickness of the sand cake; that of 8 cm revealed to be a good compromise between the number of faults being produced and the height of their surface escarpments.

## 2.2. Measurements

Measurements were performed at each 1 mm-decompression step, hence quasi-continuously. Yet, for clarity and brevity reasons, we only describe results at each 1 cm-step. In total, we decompressed the bar by 12 such steps, making the foam still slightly compressed by the end of the experiment. The total experiment corresponds to having stretched the sand model by 20%, by successive steps of 1.7%.

First, numerical high-resolution pictures were taken with a camera at each 1 mm-extension step. The combination of these pictures allows the space-time evolution of the faults to be visualized quasi-continuously (supporting movie, online material C). Second, the model was equipped with an interferometric laser system allowing accurate (pixel size 0.3 mm) digital elevation models (DEM) of the sand surface to be constructed at each step. The laser system operates as standard Radar Interferometry (e.g., Bürgmann et al., 2000) and is described in Gravelleau and Dominguez, 2008. Twelve DEMs were constructed to monitor the evolution of the model topography at every 1 cm step of extension. From each DEM, we extracted elevation profiles following the top and base of the fault scarps. Where a fault trace had several closely spaced strands at the same along-strike

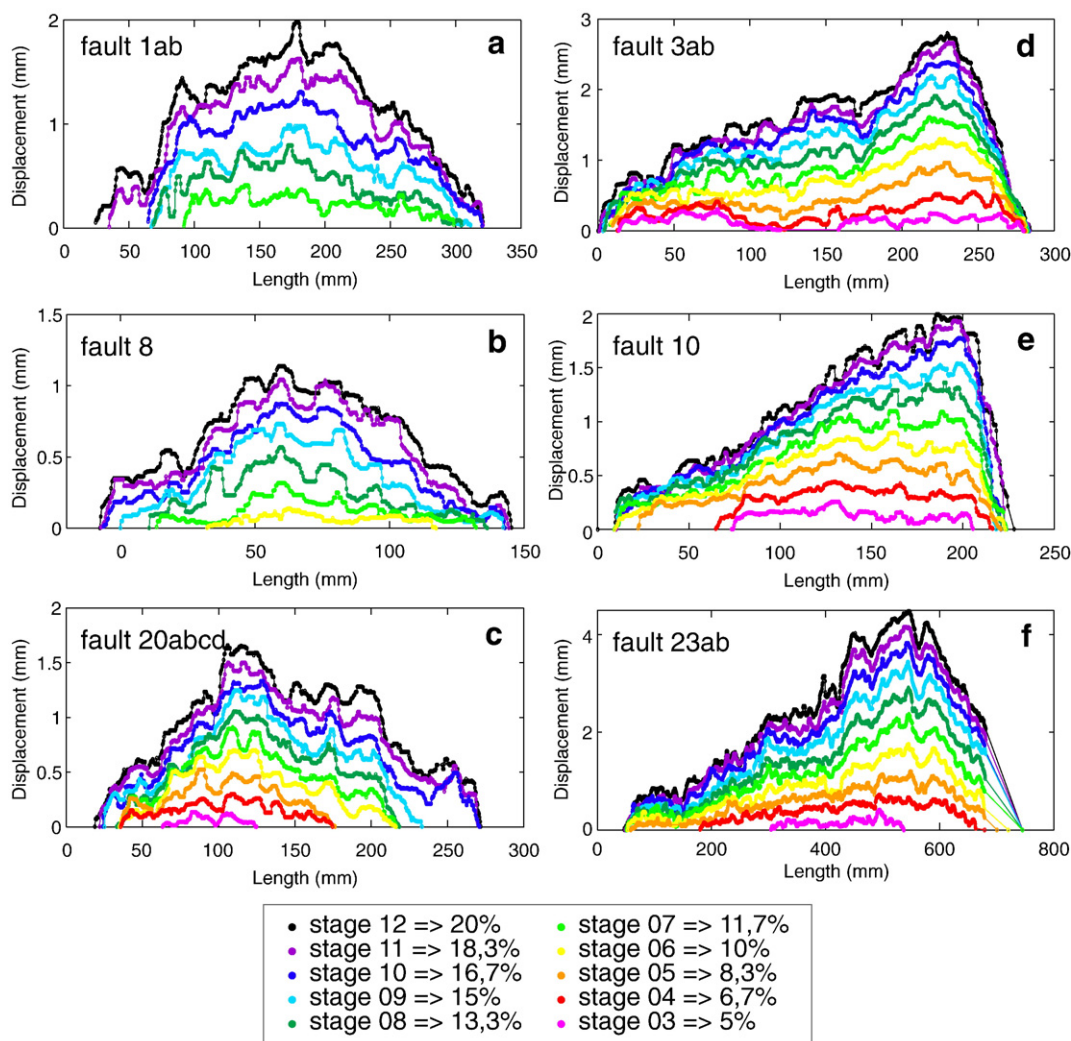


Fig. 2. Slip-length profiles on six example faults (location in Fig. 1). Naming a fault 'ab' for instance means that the fault, that contains two segments a and b, is considered as a whole. The color scale indicates the successive time steps, in percentage of extension. Stages 1 and 2 do not appear because no surface escarpment was formed at these stages.

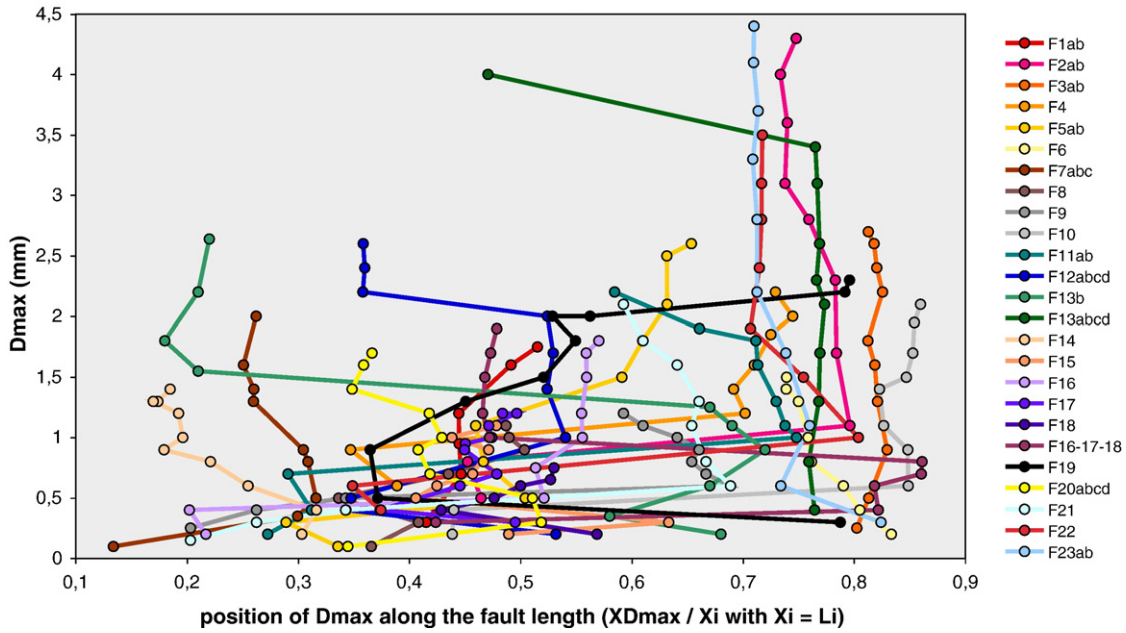


Fig. 3. Variation over time of the along-strike position of the maximum displacement  $D_{max}$ . Each point represents the position of  $D_{max}$  at each stage of the fault growth, hence is  $(XD_{max}/X_i)$ , with  $X_i$  the length of the fault at stage  $i$ , and  $X=0$  taken at the left tips of the faults. For clarity reasons, only the largest faults are represented.

position, the extracted profiles were made to include them. From the top and base profiles, we calculated the variation of vertical displacement along strike. The faults have steep scarps ( $\geq 60^\circ$ ) so that those vertical displacements approximate the actual fault slips. Results were then plotted as strike-parallel vertical projections for analysis (Fig. 2). Overall, about 100 faults formed and evolved in the interior of the central frame (Fig. 1). We restricted our analysis to the largest ones. A few faults initiated in the central frame and propagated outside of it. We considered only those that ended hardly beyond the frame. In total, we thus analyzed a population of 30 faults, synthetic and antithetic, representing a total of 250 incremental slip profiles.

2.3. Errors

Most faults have lengths that rapidly (at most, over the first third of the experiment) became greater than twice the thickness of the sand layer, which they thus cut through. We are thus confident that the 2D measurements of lengths and displacements performed at the model surface are representative of the 3D fault evolution at depth, at least over the last two thirds of the experiment.

Errors on the measurements have multiple sources. The largest likely arises from the hand- extraction of the elevation profiles, particularly during the first stages of growth when fault traces are

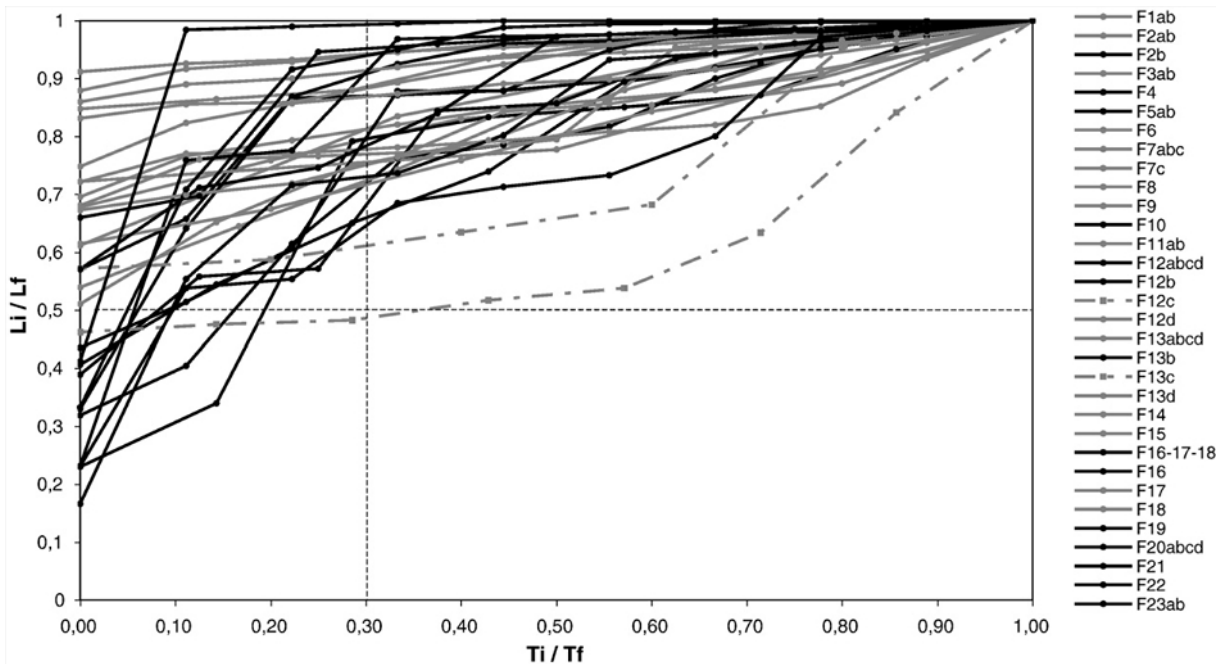


Fig. 4. Fault length increase as a function of time. On the X axis, time ( $T_i$ ) is normalized by the total duration of the fault growth ( $T_f$ ). On the Y axis, the length ( $L_i$ ) is normalized by the final total length of the faults ( $L_f$ ). See discussion in text.

tenuous. The smallest slips that we could distinctly measure in the first stages of the experiment are 0.1–0.15 mm, and happened to be on faults a few cm-long. Also, the granular nature of the sand introduces a roughness of  $\approx 0.1$  mm. The calibration of the laser introduces a periodic noise in the DEMs, having an amplitude and a wavelength of 0.25 and 10 mm, respectively. We removed that noise by slightly filtering the slip profiles (supporting online material D). From the combination of the errors above, we estimate that fault length and slip are determined with a maximum error of a few mm and of 0.1–0.15 mm, respectively, over most of the experiment but the very first

stages. In those first stages (at most, stages 1–4), fault length may be underestimated by a few cm.

### 3. Data analysis

The experiment is presented as a movie in supporting online material C. Most faults start forming surface traces soon after extension has begun. Faults initiate evenly over the central section of the model, which confirms that the extensional stress field is homogeneous. Fault traces evolve from tenuous zones of dilatance in

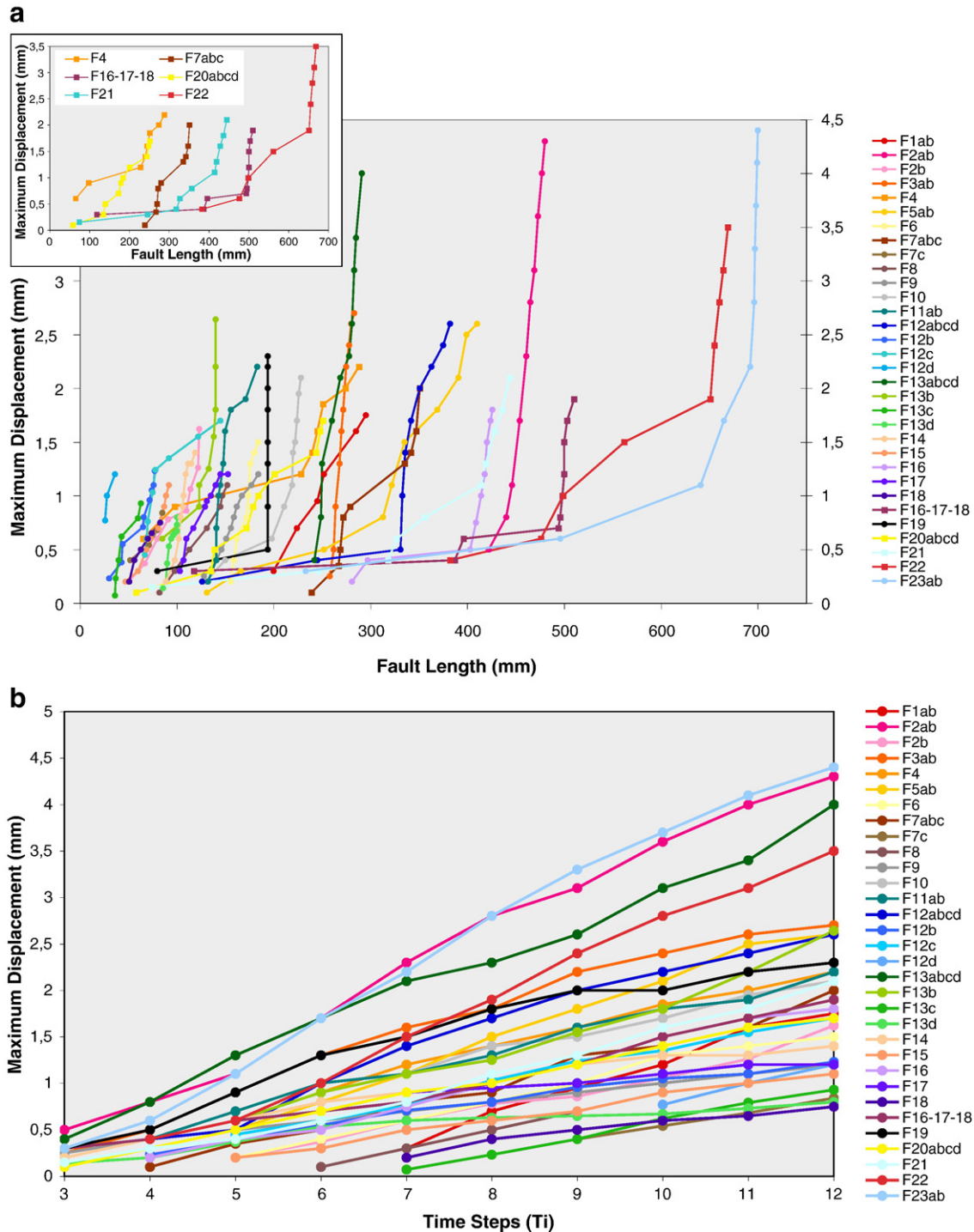


Fig. 5. a) Evolution of maximum displacement with fault lengthening. Faults possibly sustaining alternating phases of dominant lengthening and of dominant slip accumulation are highlighted in inset. b) Evolution of maximum displacement with time. See discussion in text.

the first stages, to topographic escarpments that grow higher and longer as extension increases. In the following, we define 'fault initiation' as the stage when vertical displacements become resolvable at the surface (stages 3 or 4, at most 5% extension). 'Slip' and 'displacement' are used similarly. Maximum displacement is noted  $D_{\max}$ . A 'slip increment' (or 'slip event') is the slip-length change recorded by a fault at the model surface as the model sustains an additional 1 cm-step of extension. 'Slip rate' is used to describe the displacement accumulated over a specific time span.

### 3.1. Overall shape of the slip profiles

As they grow, faults develop slip profiles that generally have a well-defined envelope shape (Fig. 2). Two types of shapes are most commonly observed, that of a symmetric triangle (Fig. 2a-b-c) and that of a strongly asymmetric triangle (Fig. 2d-e-f). A few more complex shapes are found; yet they generally arise from the evolution of triangular shapes (supporting online material E). These more complex shapes have in common to exhibit high slip gradients at both tips. In some cases, faults maintain the same overall shape as they grow (ex: Fig. 2a-b-d-f); in other cases, the shape of the slip profile changes as the fault grows: although still triangular overall, it shifts from symmetric to asymmetric (ex: Fig. 2c-e) or vice-versa. This shows that the position of the maximum displacement may vary in time along a fault (Fig. 3). While the slip profiles show any degree of asymmetry in the very first stages of growth (Fig. 3 and supporting online material F-a), they pertain to only two categories, symmetric or asymmetric, in the last stages (supporting online material F-b). Commonly, asymmetry and high slip gradient develop where a growing fault meets another, generally oblique or antithetic fault (supporting online material C).

### 3.2. Fault lengthening versus slip increase – overall behavior

Fig. 4 and supporting online material G show the evolution of fault lengths as a function of time. About 45% of the faults (black, Fig. 4) initiate with a 'short' length generally less than half their final length, then rapidly lengthen laterally in the following few stages so that, at

30% of the experiment, they have reached a length almost as long as their final one. About 50% of the faults (grey, Fig. 4) conversely initiate with a length more than half their final length, then hardly lengthen further laterally over the whole experiment. Note that, because fault lengths may be underestimated in the very first stages of the experiment, the 50% value is a minimum. Finally, two faults (dotted grey, Fig. 4) initiate with a length averaging their half final length, then hardly lengthen during most of the experiment but at the very last stages when they resume lengthening. Together these show that at least 80% of the faults achieved 70% or more of their final length in the first third of the faulting history. This demonstrates that faults grow in two distinct, successive phases. They sustain an initial phase when they essentially gain in length and this occurs rapidly; then they sustain a subsequent period when they go on growing without much lengthening.

Whether they are lengthening or not, faults accumulate vertical slip (Fig. 5a). The variation of maximum displacement versus time plots roughly as a straight line for each fault (Fig. 5b), suggesting a roughly regular vertical growth. The fan-looking distribution of the  $D_{\max}$ -Time curves further suggests that the vertical growth is more or less rapid depending on which fault is concerned; the shortest faults seem to be those most slowly increasing in slip (Fig. 5a). When variations of  $D_{\max}$  are scaled to fault size (Fig. 6), they appear roughly similar for all faults but the smallest ones. Hence, although it occurs at different rates, and although the locus of maximum slip may move over time along the faults, the vertical growth, here taken as the progressive increase in  $D_{\max}$ , is homogeneous and regular overall.

Fig. 7 depicts how the ratio of maximum displacement over length ( $D_{\max}/L$ ) varies as the faults lengthen. When faults are in a stage of rapid lengthening (Fig. 7a), their  $D_{\max}/L$  ratio keeps roughly constant, on the order of  $1-2 \cdot 10^{-3}$ . By contrast, when faults are in a stage of slip accumulation with little or no lengthening (Fig. 7b), their  $D_{\max}/L$  ratio increases as the faults accumulate more slip (up to one order of magnitude). It results that the  $D_{\max}/L$  ratios spread a broad range of values when the whole fault population is considered over the whole time of the experiment (supporting online material H). Faults with triangular profiles (red and orange) seem to have lower  $D_{\max}/L$  ratios than faults whose profiles end with high slip gradients (blue).

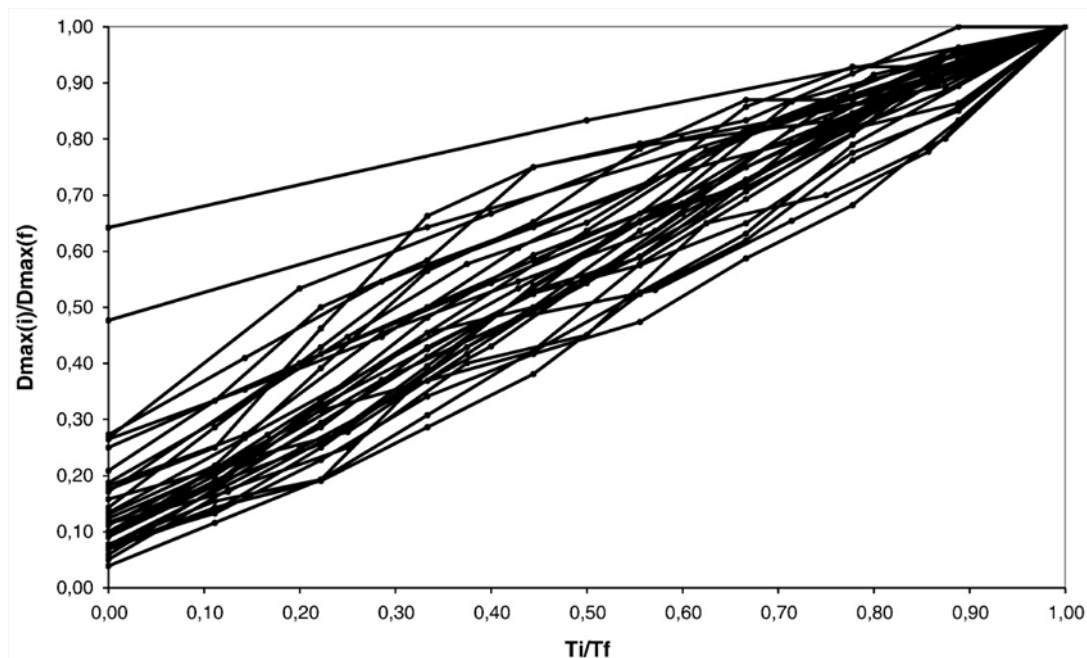
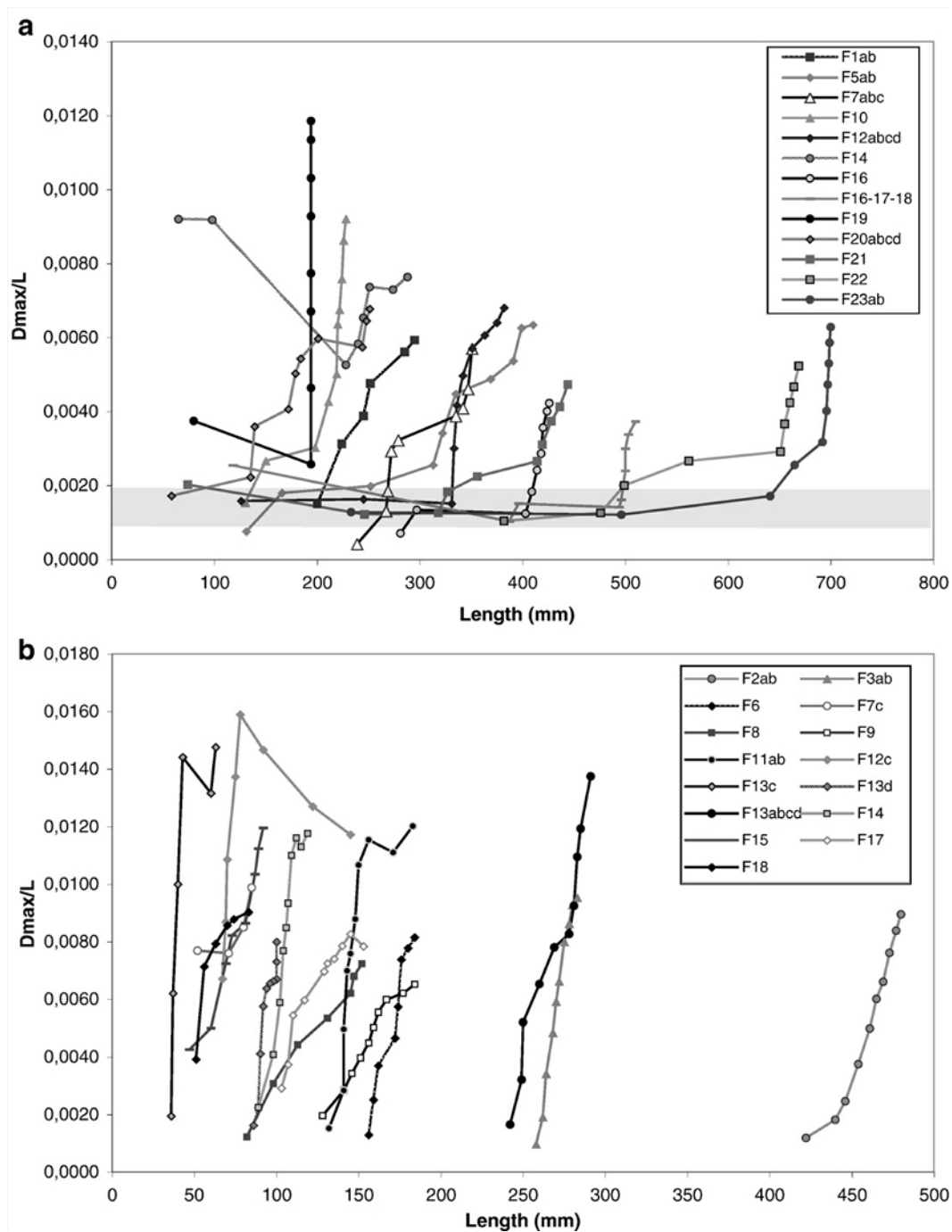


Fig. 6. Increase in vertical slip as a function of time. On the X axis, time ( $T_i$ ) is normalized by the total duration of the fault growth ( $T_f$ ). On the Y axis, the maximum slip ( $D_{\max i}$ ) is normalized by the final total maximum displacement on the faults ( $D_{\max f}$ ).



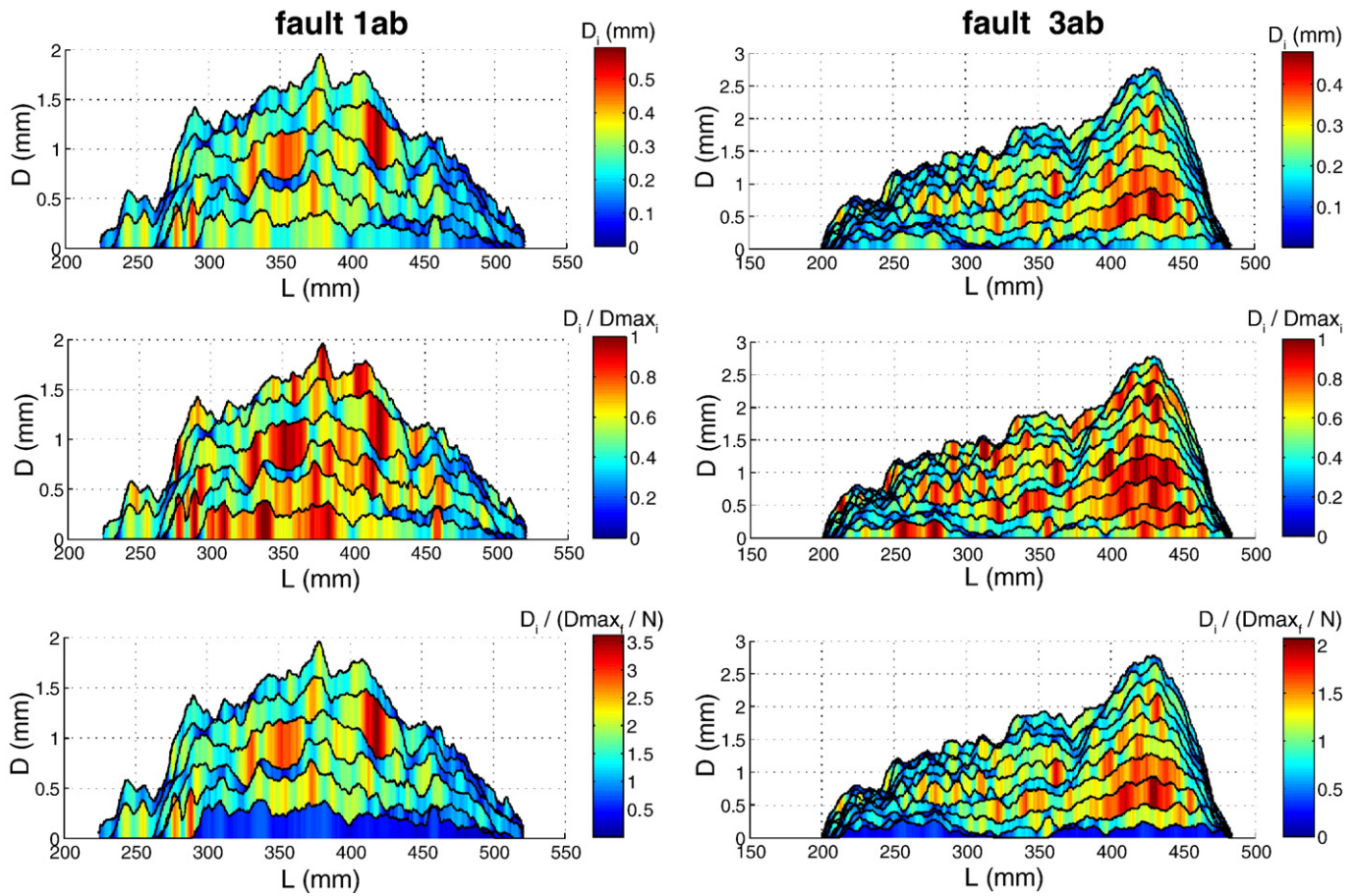
**Fig. 7.** Variation of the  $D_{max}/L$  ratio with fault lengthening. a) Faults whose growth starts by an observable, rapid lengthening phase. That occurs at roughly constant  $D_{max}/L$  ratio (grey horizontal bar). b) Faults whose initial lengthening phase is too rapid to be observed (it occurred at less than 5% of extension); those faults thus appear as mainly growing by increasing their slip. See discussion in text.

### 3.3. Fault lengthening versus slip increase – small-scale features

Additional observations can be made when fault growth is analyzed in greater details. While most faults show an initial phase of rapid lengthening, they do not extend similarly (supporting online material I). In most cases, the fault looks as lengthening continuously (at the resolution of the experiment), through the lateral propagation of one or both of its tips. In other few cases, the fault lengthens abruptly by connecting to another, initially distinct, fault segment. Generally, that segment initiates just ahead of the expanding tip of the main fault, hardly before the connection is made. In more rare cases,

the segment initiates at some distance from the expanding fault, much before the connection occurs.

Fig. 8 shows the details of the incremental growth of two example faults. The faults look like made of small patches that slip differently in both space and time. At any one time-step, the faults slip along their entire length but the slip distribution is highly heterogeneous. At any stage of growth, several patches along the fault slip by large amounts while other zones hardly or do not slip at all (top plots). Moreover, at any stage of growth, there are several, distinct areas along the fault that slip by similar large amounts (middle plots). These large slip zones are not necessarily where the final displacement is the largest,



**Fig. 8.** Space-time slip variability during fault growth. For both fault examples, top plots present the successive slip increments in absolute amplitude (mm); middle plots present the slip increments each normalized by its maximum slip value ( $D_{max,i}$ ); bottom plots present the slip increments each normalized by the average value of the maximum slip ( $D_{max,i}/N$ , with  $D_{max,i}$  the maximum slip value at the end of the experiment and  $N$  the number of slip increments).

though the zone of final  $D_{max}$  is among those slipping most. It is important to note that most of the slip patches do not coincide with any clear segment along the fault trace. Slip increase also varies in time. At any point along a fault, slip increases by variable amounts, commonly alternating from being large to being small (top and middle plots). This results in any slip gap being eventually ‘filled’ over the entire fault history. The time over which the gap is filled is variable; it is longer in zones where distinct segments have eventually connected. One can note that the incremental slip is generally either much larger (yellow to red, bottom plots) or much lower (pale to dark blue, bottom plots) than the average value that one would infer from scaling the total final slip by the number of slip increments. Finally, the figures also suggest that faults generally grow a bit slower (in term of slip increase) at both the beginning and end of their history (bottom plots).

#### 4. Discussion

Having analyzed before hundreds of natural fault slip profiles, both long-term (Manighetti et al., 2001) and co-seismic (Manighetti et al., 2005), we cannot avoid noting that the analog profiles share many common features with the natural ones. This suggests that the experimental observations are meaningful, and that those common features are inherent properties shared by natural and analog faults.

##### 4.1. Common features to natural and analog slip profiles

A first feature common to natural and analog slip profiles is their triangular, and commonly asymmetric, envelope shape. Indeed, while

it has long been thought that geological faults had elliptical or bell-shaped slip profiles in accordance with the predictions of the elastic crack theory (e.g. Pollard and Segall, 1987; Cowie and Scholz, 1992a,b; Scholz, 2002), a number of recent observations show that faults more commonly exhibit triangular slip distributions (Peacock and Sanderson, 1996; Nicol et al., 1996; Cowie and Shipton, 1998; Cartwright and Mansfield, 1998; de Jossineau et al., 2001; Manighetti et al., 2001 for normal faults; Ellis and Dunlap, 1988; Shaw et al., 2002; Davis et al., 2005 for reverse faults; Pachell and Evans, 2002 for strike-slip faults). Besides, although most available fault models assume that maximum displacement is at the fault center (e.g. Pollard and Segall, 1987; Cowie and Scholz, 1992a; Willemse et al., 1996; Cowie and Shipton, 1998; Kim and Sanderson, 2005), natural faults more commonly show asymmetric slip profiles (e.g. Peacock and Sanderson, 1991; Cartwright and Mansfield, 1998; Maerten et al., 1999; Manighetti et al., 2001; Shaw et al., 2002; Pachell and Evans, 2002; Davis et al., 2005). We make similar observations here; most analog faults grow by developing asymmetric, triangular slip profiles, basically made of a long linear section gently tapering to zero at one fault tip while ending through a high slip gradient at the other fault tip. As for natural faults, the long linear sections slant in the direction of fault lengthening (ex: faults e-f, Fig. 2). Also as for natural faults, the high slip gradients develop where the fault meets an ‘obstacle’ (or ‘barrier’) to its further lengthening (as an example, see the right tip of fault 3ab in figures 1 and 2d). That is commonly a parallel antithetic fault. Thus, as for natural faults, a slip profile remains symmetric as long as the fault has not met any barrier, and starts distorting once one fault tip at least has met an obstacle to its further lengthening. In that case, the zone of



maximum slip progressively translates toward the 'arrested' tip, where it may subsequently remain stable for a long time. Yet if the other fault tip also meets a barrier, the zone of  $D_{\max}$  may shift again laterally. Thus the position of  $D_{\max}$  may change over the fault lifetime. Furthermore, the zone where slip ends to be largest does not necessarily coincide with the locus of fault initiation, as it has been suggested (e.g. Ellis and Dunlap, 1988). It is important to note that the development of long linear sections in the slip profiles is not in keeping with the elastic crack theory (see Manighetti et al., 2001 for an extended discussion). One way to reconcile these profiles with that theory is to admit that the excess stresses that they produce on the fault plane are diffused off that plane, possibly in the form of distributed damage (Manighetti et al., 2004). We do not have the resolution to verify whether or not such diffuse deformation occurs nearby the linear slants.

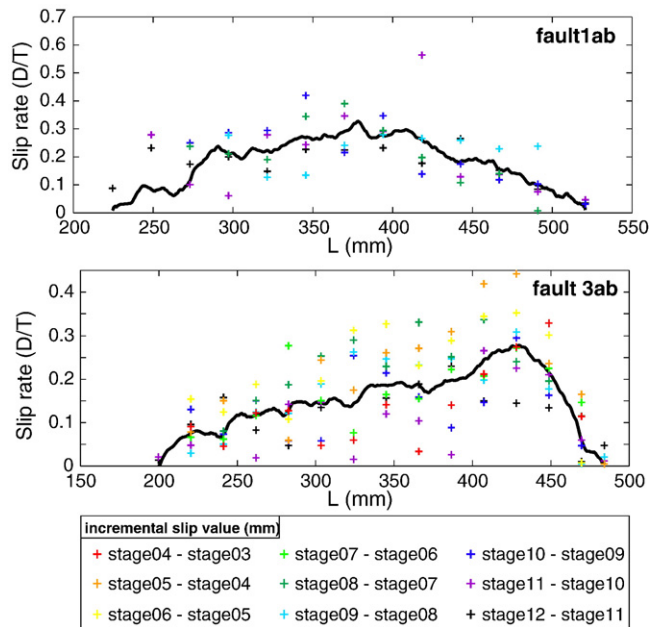
A second feature common to natural and analog faults is their bimodal mode of growth; faults are found to grow in two successive phases, an initial period of dominant, rapid lengthening, and a subsequent phase of dominant slip increase with no or little lengthening. Both phases have been observed in natural cases, but the scarcity of available data in respect to the long fault lifetimes ( $10^5$ – $10^7$  yr) has not allowed describing them in detail (Armijo et al., 1996; Morewood and Roberts, 1999; Contreras et al., 2000; Poulimenos, 2000; Goldsworthy and Jackson, 2001; Meyer et al., 2002; Walsh et al., 2002; Gawthorpe et al., 2003; Childs et al., 2003; Taylor et al., 2004; Bennett et al., 2005; Nicol et al., 2005). The numerical modeling and analog experiment of Cowie and Shipton (1998) and Mansfield and Cartwright (2001) respectively, also suggest the existence of such two phases, but the data are too few to describe them in detail. The initial lengthening phase is brief, hardly more than the third of the fault lifetime (considering the experiment duration as the 'lifetime'), while the lengthening is rapid with most faults reaching more than their half-final length by the end of that initial phase. The lengthening generally occurs through the lateral propagation of the fault tip(s) (see Marchal et al., 1998, for similar observations). We do not have the resolution to determine whether or not the fault tip propagation is a continuous process. Nor do we have the resolution to depict how some faults initiate with a length almost as long as their final one. In any case, such initial long lengths do not result from pre-existing structures being reactivated as suggested or observed in some natural cases (Meyer et al., 2002; Walsh et al., 2002; Nicol et al., 2005; Kim and Sanderson, 2005), nor from the linkage of simultaneously formed major segments, as commonly proposed (e.g., Cowie and Roberts, 2001). While they are lengthening, the faults accumulate some slip. The process is homogeneous overall, such that the  $D_{\max}/L$  ratio remains constant on each fault and similar for all faults ( $2 \pm 1.10^{-3}$ ). This suggests that faults behave as elastic cracks with constant stress drop during their initial lengthening phase. Following the initial period of rapid lengthening, the faults enter a subsequent phase during which they mainly accumulate vertical slip while hardly or no longer lengthening. Note that while many faults stop lengthening when at least one of their tips has met a barrier, that condition is not required. The phase of dominant slip accumulation is homogeneous as far as it is considered overall. The maximum displacement increases quite regularly on each fault, and in same proportions for all faults. Such a homogeneous vertical growth results in the cumulative slip profiles keeping the same overall shape, – only distorted vertically – as the faults accumulate more slip. Overall, the phase of dominant slip accumulation lasts longer than that of lengthening, generally about two thirds of the fault life. It makes the faults increasing their  $D_{\max}$  by 300% on average (and up to 700%; in respect to the value of  $D_{\max}$  at the beginning of the slip accumulation phase). The  $D_{\max}/L$  ratios increase accordingly during that phase, making the faults not behaving as elastic cracks and not having similar stress drops. By the end of the experiment, most faults still are in a phase of slip accumulation at constant length. We thus have no data to document how the faults go

on growing as they accumulate even more strain. In natural cases, it has been found that faults do not accumulate slip over a maximum threshold  $D_{\max}/L$  ratio averaging  $10^{-1}$  (e.g., Manighetti et al., 2001). This implies that faults must lengthen laterally once they have reached that maximum slip accumulation ratio. This would suggest that faults may resume lengthening by the end of a phase of dominant slip accumulation. The inset of Fig. 5a indeed shows a few fault cases that resume lengthening by the end of a slip accumulation phase. It is noteworthy that those few faults have in common to have their tips not being arrested by any clear barrier (see example faults, f6, f11ab or f19 in Fig. 1), contrary to the rest of the faults. This suggests that faults pinned at both ends may accumulate more slip without lengthening than unrestricted faults. The same conclusion has been reached for natural cases (Manighetti et al., 2001). Note that, when resuming after a slip accumulation phase, the lengthening occurs through the connection of the main fault with a smaller segment ahead of its tips(s), hence differently than over the initial lengthening period.

A third feature common to analog and natural faults is their capacity of being together homogeneous and heterogeneous; while the fault growth process is clearly homogeneous overall, it is highly irregular in detail. This duality has been suggested for a few geological faults (e.g. Manighetti et al., 2001; Bull et al., 2006), as in a few analog and numerical experiments (Cowie and Shipton, 1998; Mansfield and Cartwright, 2001). The balance between regularity and irregularity roots in the time scale. When the fault growth is considered over a long time span that includes many slip increments, it appears homogeneous; the slip profile maintains the same envelop shape; the slip and length increases are regular overall. In contrast, when the fault growth is considered over time scales that are short in relation to the entire fault history, it appears highly heterogeneous; the slip profile shape exhibits saw-tooth patterns; the slip and length increases are markedly variable in space and time.

#### 4.2. Implications on $D_{\max}/L$ 'scaling'

The bimodal growth of the faults causes their maximum displacement and length to change by various proportions over time. As recognized by Mansfield and Cartwright (2001) and Walsh et al. (2002), those changes produce irregular growth paths in  $D_{\max}$ - $L$  space, that diverge from the roughly linear relationship that has been derived from available natural  $D_{\max}$ - $L$  data ( $D_{\max}=k \cdot L$  with  $k \approx 10^{-2}$ ; e.g., Schlische et al., 1996; Manighetti et al., 2001; Davis et al., 2005). The question then arises of whether or not a scaling of the form  $D_{\max}=k \cdot L$  may be used to prescribe the way natural faults are growing (e.g. Cartwright et al., 1995; Mansfield and Cartwright, 2001). Our experiment suggests that the answer roots in the time scale over which the fault growth process is analyzed. When that process is considered over time scales that are short in relation to the entire fault history, the scaling does not apply and data exhibit a broad range of scatter (Fig. 5a). This has been shown to be true in natural cases: on the shortest time scale which is that of individual earthquakes, displacement and length of ruptures are clearly not related by a unique function (Manighetti et al., 2007). By contrast, when the fault growth is considered over a long time span including a large number of slip increments, the linear scaling likely applies, as revealed by the natural data. Thus, for natural faults whose growth history is unknown, the  $D_{\max}$ - $L$  measurements may be difficult to interpret. Manighetti et al. (2001) have shown that the  $D_{\max}/L$  ratio of a fault depends on the overall shape of its slip profile: profiles unrestricted at both tips (symmetric triangles) have the lowest ratios reaching at most  $2 \cdot 10^{-2}$ ; those restricted at one tip (asymmetric triangles) have higher ratios with a threshold value of  $3$ – $4 \cdot 10^{-2}$ ; those pinned at both tips have the highest values up to  $10^{-1}$ . This suggests that the combined knowledge of a fault slip profile shape and of its corresponding  $D_{\max}/L$  ratio may indicate the current 'growth mode' of the fault: dominant slip accumulation if the measured ratio is far



**Fig. 9.** Comparison of 'long-term' and 'short-term' slip rates. For both fault examples, the black curve is the long-term slip rate profile that one would infer from dividing the final cumulative slip profile by the total number of growth time steps. Conversely, the color crosses represent the short-term slip rates that one would infer from dividing each incremental slip value by one single time lapse. Colors indicate, as before, the succession in time.

below the threshold value corresponding to its profile shape; dominant lengthening if that ratio is close to the threshold value. Though more variable, our analog data are in keeping with that inference (supporting online material H).

**4.3. Implications on space-time slip and slip rate variability**

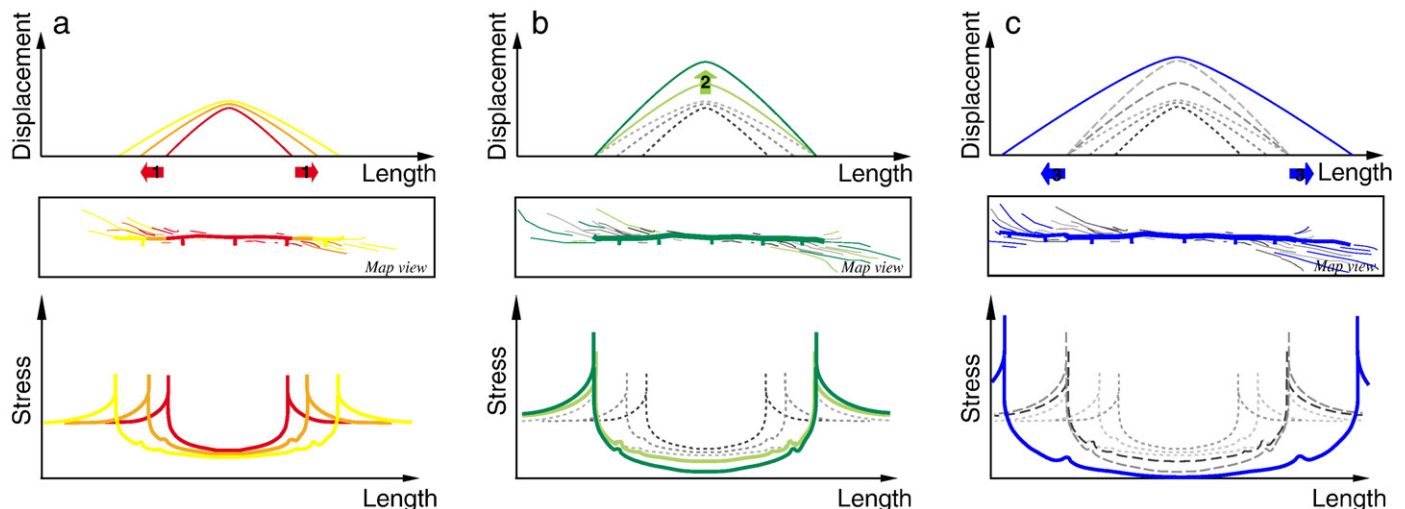
Determining the rates at which faults slip is a critical issue for these rates indicate the degree of activity of the faults (e.g. Nicol et al., 1997). Numerous studies have thus been conducted in the last 30 yr to determine the past and present motion rates on the major active faults worldwide. Most, however, have restricted their analysis to a single slip rate determination at one spot of the fault length (e.g. Wright

et al., 2004; Frankel et al., 2007). The underlying assumption is that the slip rate does not vary much along fault strike. Our experiment suggests that this may be incorrect. At any instant of the fault evolution, the displacement varies markedly along the fault length; so thus would the slip rate. Studies providing several slip rate estimates along a fault actually reach the same conclusion (e.g. Tapponnier et al., 2001; Taylor et al., 2004; Bull et al., 2006). Most studies that report slip rates also implicitly admit that fault rate does not vary much either in time (e.g. Wright et al., 2004). Our experiment shows that this may be incorrect. Fig. 8 shows that slip varies widely in time at one given spot of the fault length. Fig. 9 is even clearer. In each plot, the black curve is the mean slip rate that one would infer from measuring and dating the total cumulative slip profile. It is thus equivalent to the long-term slip rate profile. The color dots are the slip rates that one would infer from measuring the slip produced at each step of the growth, with the step duration being known. The black curve clearly results from highly variable slip rates being averaged. This shows that, at any spot of the fault length, the slip rate averaged over a long time span is far different, by an order of magnitude, from the slip rate averaged over a short time span. The incremental rate is particularly low both in the very first and very last stages of the fault growth. The experiment thus demonstrates that the slips and slip rates on faults are markedly variable in both space and time. A similar conclusion has been reached in the rare studies that analyzed the natural fault slip behavior at different space and times scales (e.g. Weldon et al., 2004; Chevalier et al., 2005; Bull et al., 2006). One consequence is that the common assertion that earthquake ruptures showing a similar slip amplitude at one local spot of a fault are characteristic (e.g. Nishenko and Bulland, 1987; Rockwell et al., 2001; Klinger et al., 2003; Haibing et al., 2005) is likely erroneous.

**4.4. Proposed scenario of overall fault growth**

Fig. 10 presents a scenario of overall fault growth that includes the major observations described above. That scenario is only qualitative for modeling the physics of faulting and earthquakes is not the scope of the paper. Also, it only describes the general features of the growth process for a fault that would ideally remain isolated throughout its evolution. Integrating the small-scale complexities of the fault evolution is a further step of work.

In a first phase (Fig. 10a), the fault rapidly propagates laterally while accumulating slip in proportion so that, overall, the fault behaves as an elastic crack. At this stage, the lengthening likely occurs continuously.



**Fig. 10.** A qualitative scenario of fault growth compatible with the observations. The scenario is drawn for an ideal, isolated fault, thus having a triangular, symmetric slip profile. Top plots present the slip distribution on the growing fault. Middle plots present the fault in map view. Bottom plots presents the stress distribution resulting in the 'faulted volume' (fault plane+damage zone). The dotted grey lines remind the previous stages. See the text for more details.

While lengthening, the fault develops a roughly triangular slip profile. As that profile develops, the medium holding the linear slants sustains extended damage, possibly in the form of multiplied micro-cracking (Marder and Fineberg, 1996; Manighetti et al., 2004). Large triangular zones of damage having their apex centered on the point of fault initiation are requested to maintain a triangular profile (Manighetti et al., 2004). The damage zone accommodates the excess stresses that result from the triangular shape of the slip profile, so that stress on the fault plane remains about constant, compatible with the elastic crack behavior and the self-similarity.

The fault then stops propagating, generally as it starts interacting with nearby faults. Once arrested, the fault continues growing by accumulating more slip (Fig. 10b). As it does so, the overall shape of its slip profile is vertically distorted while the stress concentration at its tips increases. The increased amount of stress accommodated within the damage zone likely makes some of the damage cracks growing larger. Some of these cracks may eventually form well-expressed segments ahead of the main fault tips. Since excess stresses are accommodated within the damage zone, the stress on the fault plane still remains about constant.

Once slip, hence stress, accumulation is such that the strength of the damage zone ahead of the fault tips is overcome, we hypothesize that the fault resumes lengthening, likely by connecting to one of the largest damage cracks (Fig. 10c). Note that the  $D_{\max}/L$  ratios of the connecting fault and segment are unlikely to be similar, as commonly assessed (e.g., Cowie and Roberts, 2001). As the fault lengthens, its  $D_{\max}/L$  ratio drops, while its slip profile remains triangular overall. Since the fault length increases, the damage zone widens. New damage structures are thus created, while previous structures either continue accommodating stress and strain, or become inactive and possibly offset by the lengthening fault.

We suggest that faults basically grow through alternating phases of dominant lengthening and dominant slip accumulation. The initial phase of rapid lengthening that occupies the first third of the fault lifetime is likely different however than any subsequent phase of further lengthening; that initial phase occurs through fault tip propagation, while subsequent lengthening rather occurs through segment connection. Interestingly, the bimodal growth that we describe in Fig. 10 is 'self-driven' and 'self-sustaining': as a fault accumulates some slip, it damages the adjacent medium, particularly at its tips. That damage in turn helps the fault accumulating more slip (without lengthening), for it diffuses the high stresses that would otherwise result on the fault plane, particularly at its tips. As the fault accumulates more slip, stresses increase within the damage zone, so that a few individual fault segments may eventually form ahead of the growing fault. At some level of slip accumulation, the strength of the damage zone is overcome and the fault lengthens by connecting to the segment(s) previously produced ahead of its tips. The process then repeats. It is thus self-driven since the fault itself produces the features that contribute to both its steady growth at constant length (the damage zone) and its further lengthening (the segments ahead). It is also self-sustaining since, as the fault grows longer, the stresses to be diffused in the damage zone grow higher, making the process unable to stop (in the absence of any local complexity, such as 'external' barriers). One implication is that the segments that eventually connect to form a larger fault, do not initiate randomly nor simultaneously, as commonly assumed (Cowie and Roberts, 2001; Davis et al., 2005; Walsh et al., 2003; Kim and Sanderson, 2005; Nicol et al., 2005). Besides their length is somehow dictated by the 'strength' of the medium (actually, its apparent stiffness) where the fault is growing.

Note that growth in isolation conditions as described above must be seldom. Yet, the encounter of a fault with a barrier does not change the general process. The barrier only increases the strength of the zone ahead of the fault tips, and this may increase the duration of the slip accumulation phase (hence more slip is accumulated), while distorting the slip profile from symmetric to asymmetric.

## 5. Conclusions

Our laser-equipped analog experiment led to the formation and growth of a dense network of normal faults for which we could measure the slip and length changes at successive short time lapses. This provides a unique opportunity to visualize and quantify the fault growth process and to examine how slip increments follow on a fault to make it grow. The analog faults are found to exhibit a number of features similar to those observed on geological faults. This suggests that these features may be intrinsic properties of faults, which thus should be reproduced by fault models. This furthermore suggests that the experiment is meaningful, with its results likely to apply to faults of any slip mode. The quasi-continuous observation of the fault growth that we perform allows us to examine these properties with far more detail than what is possible on natural faults. The major properties that we have identified are the following:

- Slip distributions have linear sections almost as long as the entire faults, and those sections generally persist during the fault evolution. This observation is not in keeping with the elastic crack theory. It suggests that extended damage of the medium is likely to result from the growth of a fault.
- Faults start growing in two successive phases: an initial, short (first third of the fault lifetime) period dominated by rapid lengthening, followed by a subsequent, longer (last two thirds of the fault lifetime) phase of mainly slip increase at roughly constant length.
- Faults are likely to go on growing through the alternation of phases of dominant lengthening and phases of dominant slip accumulation. Additional data are however necessary to confirm this hypothesis. Such bimodal growth is likely to be 'self-driven' and 'self-sustaining'.
- Faults grow overall homogeneously even though the individual slip events are characterized by large slip and stress heterogeneities. This suggests that there is a feedback between stress heterogeneities and slip development. The consequence of this is that the long-term faulting process is primarily insensitive to the short-term heterogeneities that are rapidly smoothed or redistributed.

One implication of this duality between heterogeneous and homogeneous behavior is that the short- and long-term slips and slip rates on natural faults are likely to be markedly different. Successive slip events are likely to also significantly differ in both slip amplitude and distribution. Together these results call for great caution when interpreting the available slip data on faults and ruptures, particularly when doing so in terms of 'characteristic earthquakes', as commonly done.

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## Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at doi:10.1016/j.epsl.2008.06.042.

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## Supporting online material A – Schlagenhauf et al, EPSL, 2008.

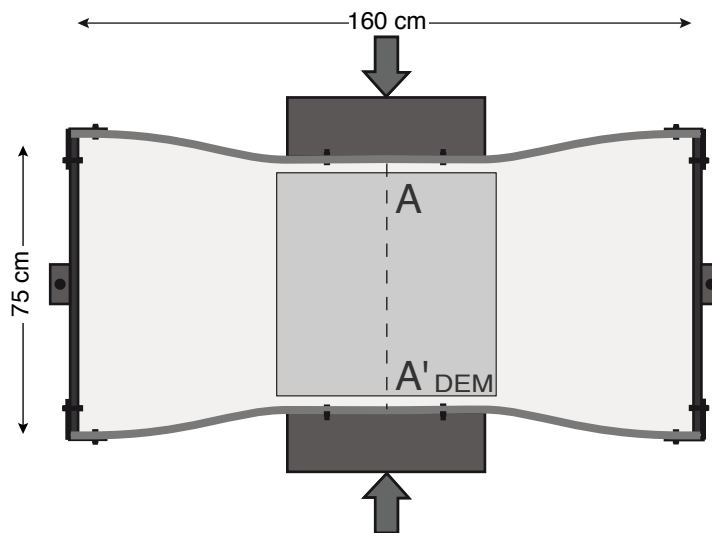
### Brief description of the extensional analog models produced.

In the majority of the analogue models published so far, the extension is generated by gravity: a basal plate is tilted, so that the ductile overlying material sustains a tangential component of the gravity force that transmits extension to the sand above (e.g. Vendeville and Cobbold, 1987; Childs et al., 1993; Marchal et al., 1998, 2003; Mansfield and Cartwright, 2001; Medwedeff and Kranz, 2002; Acocella et al., 2005). In that case, the fault initiation is diachronic, while most faults are synthetic. Other experiments produce extension by imposing a sliding discontinuity within the model: one side of the model is fixed while the other side rests on a basal plate that is progressively pulled apart. In that case, fault initiation is also diachronic; at any stage, new faults initiate above the sliding discontinuity while faults previously formed stop slipping (e.g. Basile et al., 1992; Brun and Tron, 1993; Brun and Nalpas, 1996; Acocella et al., 1999). Finally, some other models produce extension by stretching a basal rubber sheet; that stretching transmits a quite homogeneous extension to the overlying materials (e.g. McClay and Ellis, 1987; Higgins and Harris, 1997; Clifton et al., 2000; Le Calvez and Vendeville, 2002; McClay et al., 2002; Bellahsen et al., 2003, 2005; Hus et al., 2005). Yet, as the rubber sheet stretches, compressional stresses are induced perpendicular to the stretched zone. That compression interacts with the extension, producing a heterogeneous stress field when large amounts of deformation are imposed.

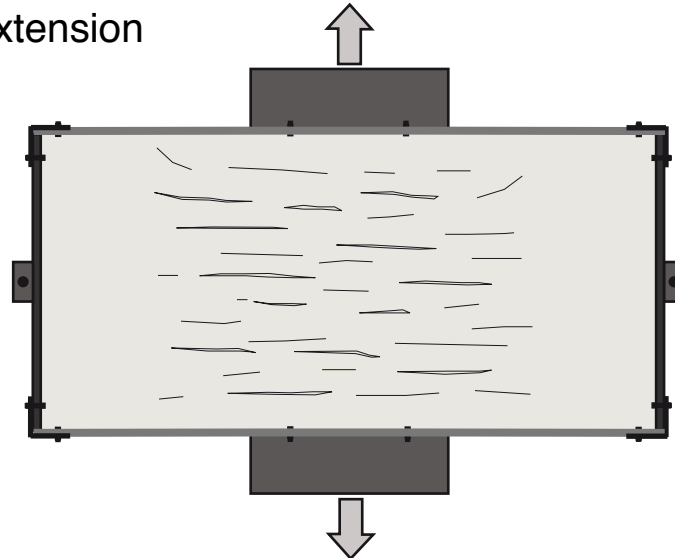
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(a) Compression of the foam and spreading of the sand (map view)



(b) Stage 2 : extension



(c) vertical section AA'



- Dry sand (8 cm thick)
- Steel bar
- Foam plate (4 cm thick)
- Viscous liquid

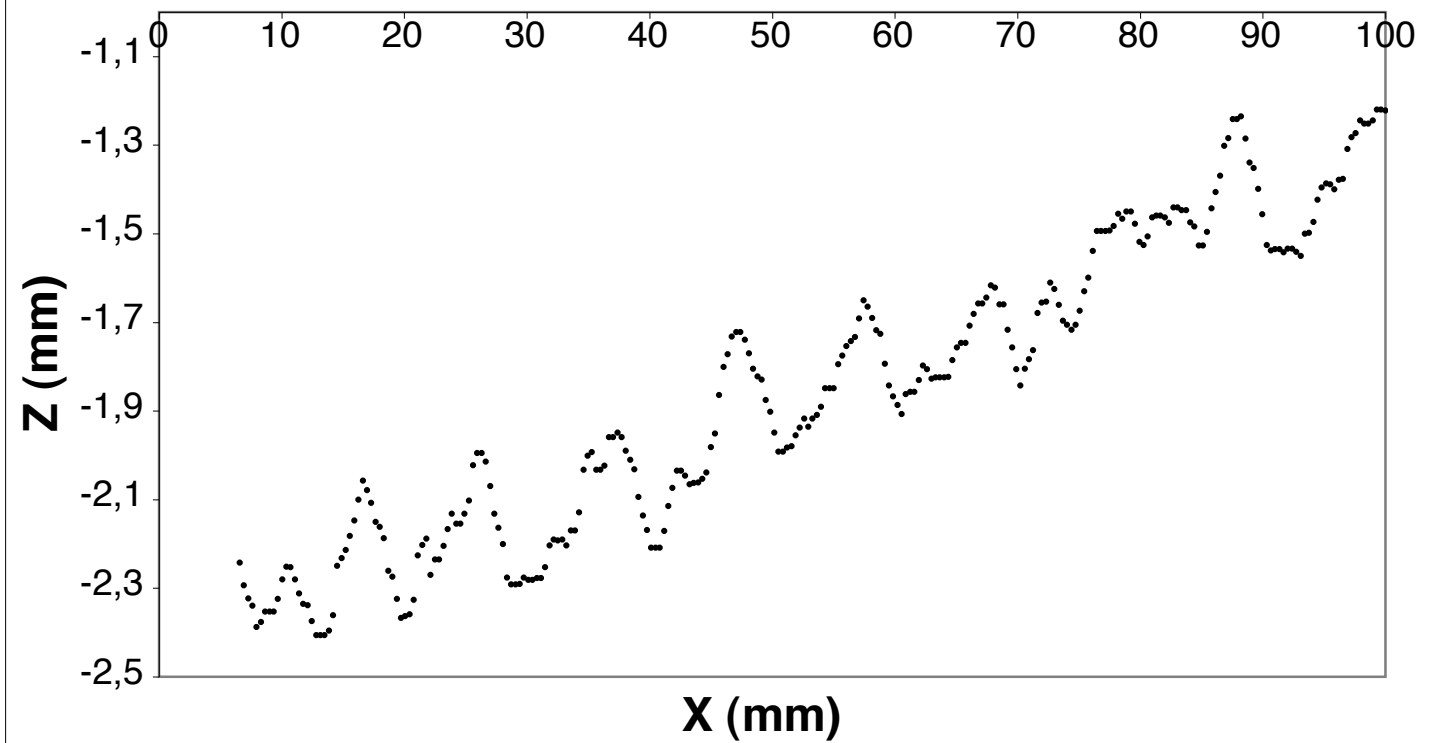
**Supporting online material B - Schlagenhauf et al., EPSL, 2008**

View of the experiment. The apparatus is 160 x 75 cm<sup>2</sup>. The 4 cm-thick foam plate is tightly enclosed within a frame having its two shortest sides rigid and its two longest sides flexible (Komacel plastic). The foam and frame are sealed together with silicone. The foam rests on a basal rigid plate. A thin layer of viscous liquid between the foam and the basal plate allows the foam to slide homogeneously with no buckling. Further details are given in text.

**Supporting online material C - Schlagenhauf et al., EPSL, 2008**

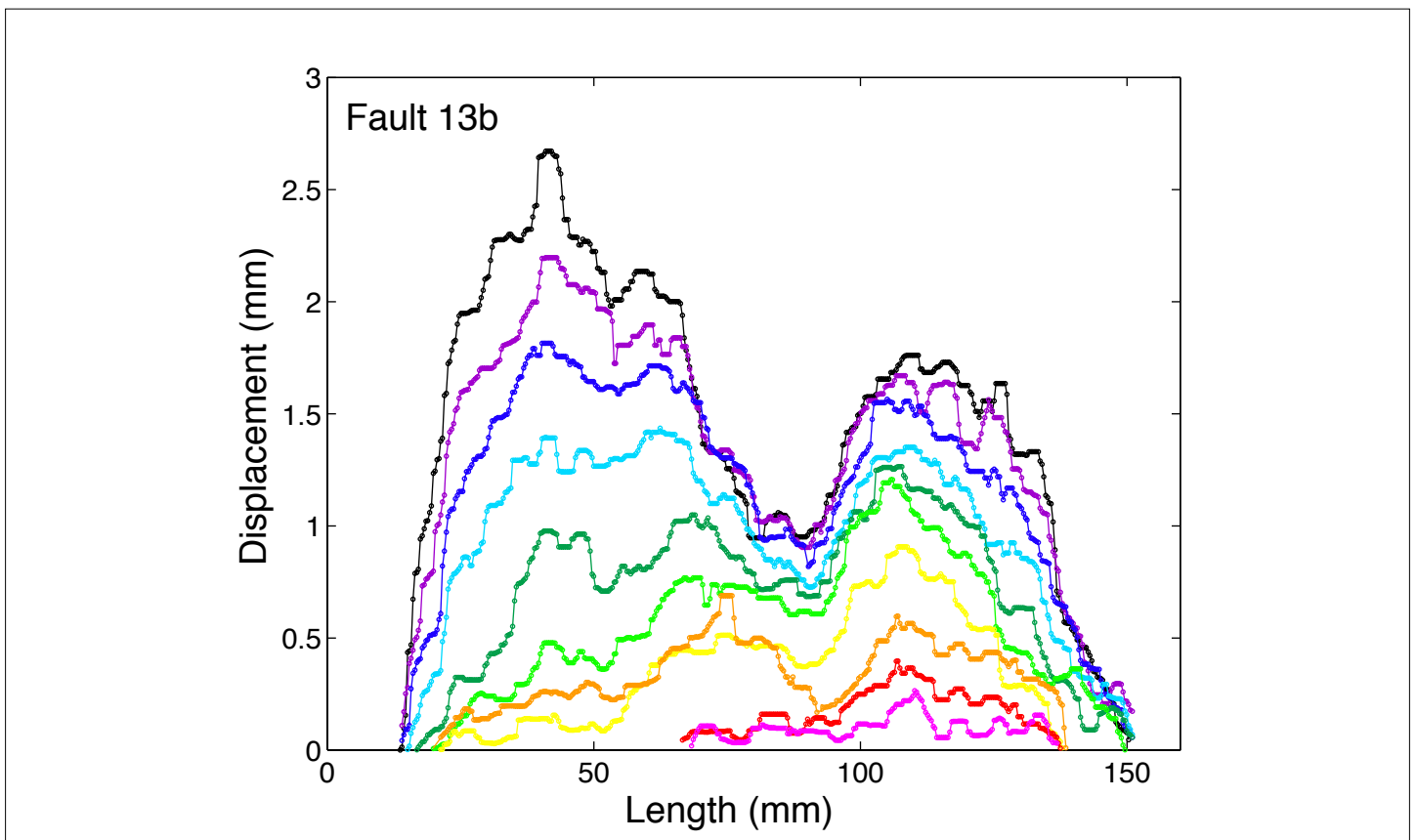
Movie, in map view, of the model from its initial stage with no extension to its final stage with 20 % of extension. Faults initiate evenly throughout the model from the very first stages of the experiment. They then go on growing over the entire duration of the experiment.





**Supporting online material D - Schlagenhauf et al., EPSL, 2008**

Example of the laser-induced noise. The noise is periodic, with a mean amplitude of 0.25 mm, and a wavelength of 10 mm that includes about 19 points of measurement. We have thus filtered the elevation profiles by averaging the data by sets of 19 points.



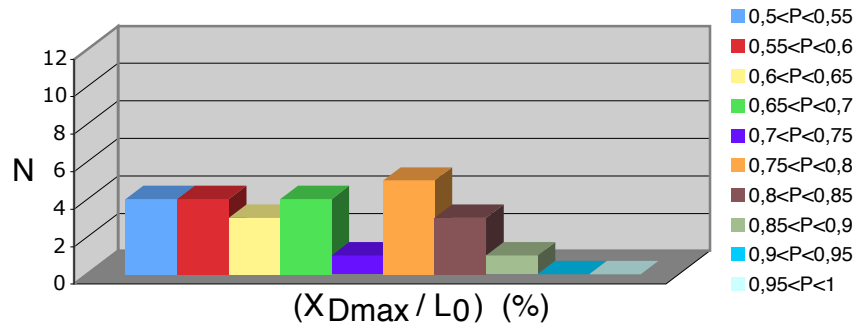
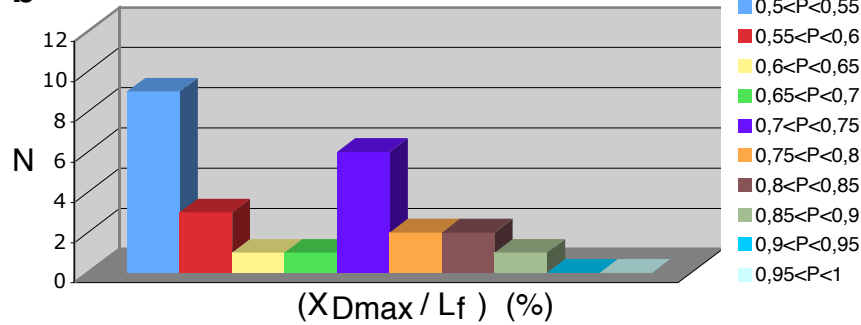
**Supporting online material E - Schlagenhauf et al., EPSL, 2008**

Example of a fault having a ‘complex’ slip profile shape. That shape however arises from the evolution of a more simple triangular shape. The final profile exhibits high slip gradients at both tips.

**a**

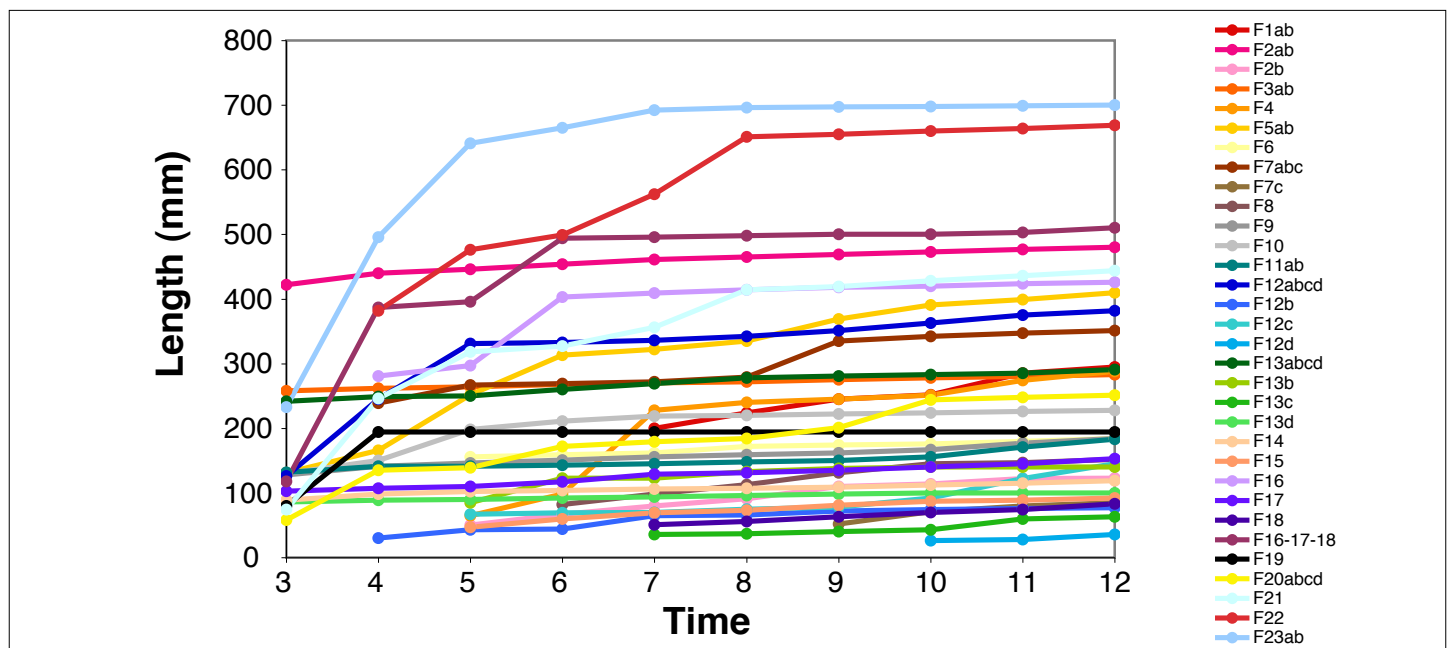
- Annexe IV - A. Schlagenhauf, thèse, 2009

IV-18

**b**

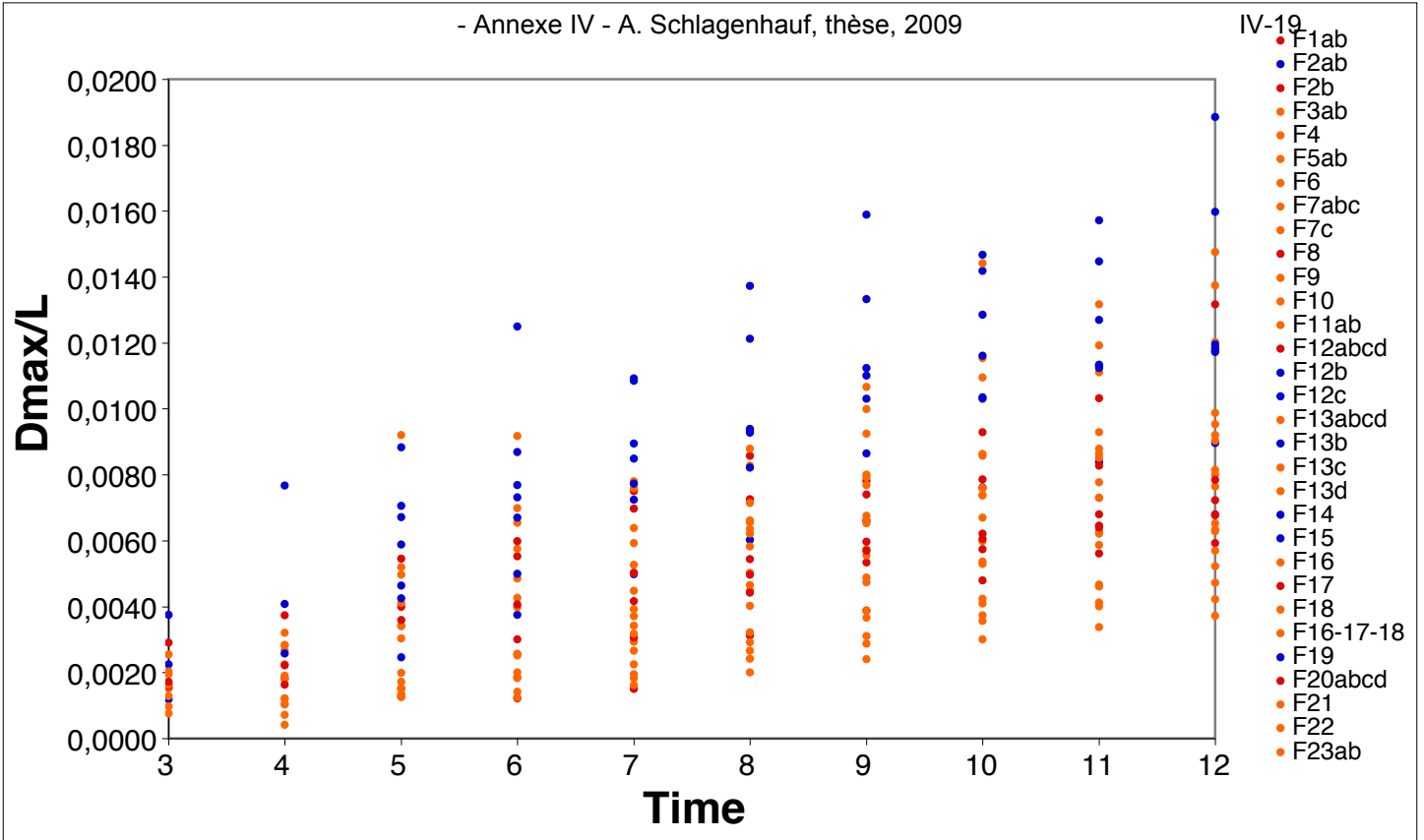
### Supporting online material F - Schlagenhauf et al., EPSL, 2008

a) Position of the maximum displacement along the faults at their initial stage of growth (which generally coincides with stage 3 or 4). The X axis plots the position of  $D_{max}$  at the first stage of growth ( $X_{Dmax}$ ) divided by the initial fault length ( $L_0$ ). b) Position of the maximum displacement along the faults at their final stage of growth (stage 12 for all faults). The X axis plots the position of  $D_{max}$  at the final stage of growth ( $X_{Dmax}$ ) divided by the final fault length ( $L_f$ ). In both figures, the position is referred to the fault half-length.



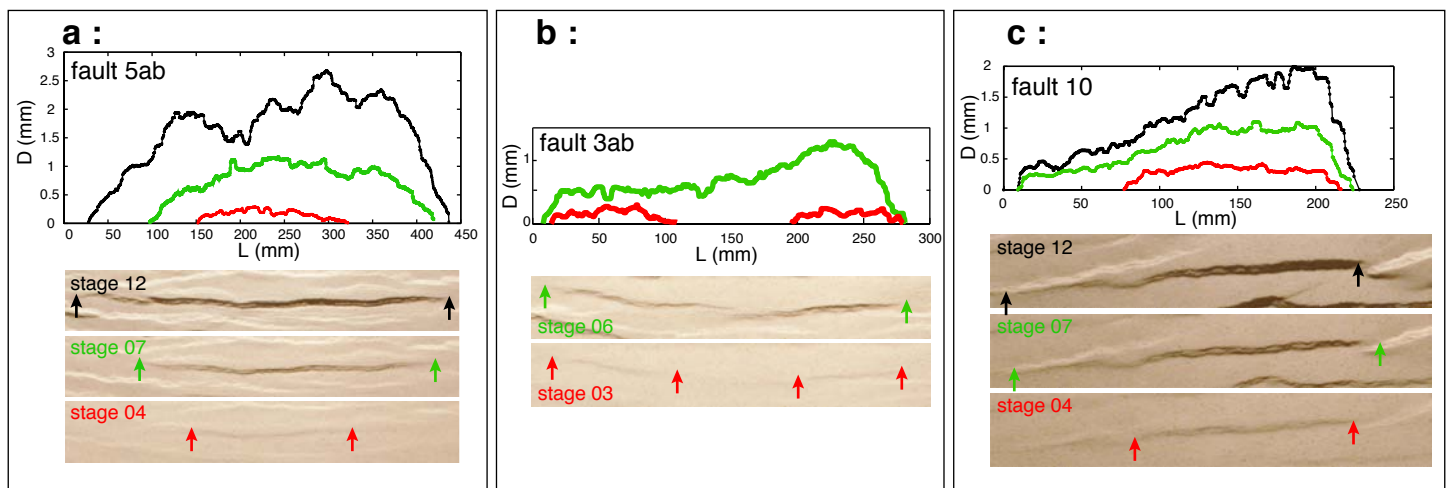
### Supporting online material G - Schlagenhauf et al., EPSL, 2008

Evolution of fault length as a function of time. More than half of the faults reach their final length in the very first stages of the experiment (under less than 5% of extension), before any vertical slip is visible at the model surface.



### Supporting online material H - Schlagenhauf et al., EPSL, 2008

Evolution of the  $D_{max}/L$  ratio as a function of time. The envelop shape of the slip profiles is indicated by colors, red for triangular and symmetric profiles, orange for triangular and asymmetric profiles, and blue for profiles with more complex shapes having high slip gradients at both tips. The later generally accumulate more slip than the former.



### Supporting online material I - Schlagenhauf et al., EPSL, 2008

Illustrating the different modes of fault lengthening. In all figures, top plots are the growing slip profiles, while bottom plots show the growing faults in map view. Colors refer to the various stages of growth. a) Both tips of the fault propagate laterally, as if continuously; the slip profile keeps roughly symmetric. b) Two initially distinct fault segments propagate toward each other and eventually meet and connect to form a larger-scale fault. c) The larger fault further lengthens by connecting to a segment formed due ahead of one of its tips. Note that the fault is rapidly arrested at both tips as it meets two antithetic faults. High slip gradients develop at the arrested tips.





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# Identifying past earthquakes on an active normal fault (Magnola, Italy) from the chemical analysis of its exhumed carbonate fault plane

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### ABSTRACT

A normal fault scarp exhumed by repeated strong earthquakes is made of a series of rupture zones that were exposed, thus weathered, over significantly different time spans. We show that such differential weathering can be detected in the chemical content of the fault scarp rocks, and its signature used as a base to decipher the past earthquake history of the fault. We focus on the Magnola normal fault, Central Italy, whose Holocene seismic slip history has already been determined by Palumbo et al. (ESPL, 225, 163–176, 2004) from *in situ* <sup>36</sup>Cl cosmic ray exposure dating of the fault limestone scarp surface. Five major earthquakes were found to have occurred over the last 12 ka, with slips of 1.5–3 m and recurrence times of 0.7–3.1 ka. We analyze the major and trace element concentrations of 15 carbonate samples collected from base to top of the 10 m-high Magnola Holocene scarp, next to the previous sampling done by Palumbo et al. [Palumbo, L., Benedetti, L., Bourlès, D., Cinque, A., Finkel, R., 2004. Slip history of the Magnola fault (Apennines, Central Italy) from <sup>36</sup>Cl surface exposure dating: evidence for strong earthquake over Holocene. Earth Planet. Sci. Lett. 225, 163–176.]. We find that most element concentrations decrease upscarp at a rate averaging 5%/m. This decrease is attributed to leaching and re-precipitation of purer calcite that increase with exposure time. Superimposed to the overall leaching, concentration peaks are found at the transition zones separating the earthquake ruptures. These concentration peaks likely result from enrichment of the scarp sections that remained stuck in the 30–50 cm-thick impurity-rich upper soil during the quiescence periods that separated the earthquakes. Because the rare earth elements (REE) are among those most significantly enriched at the earthquake transition zones, they are the best chemical markers of past large seismic events. We finally propose a first-order model that reproduces adequately the observations. Our preliminary work thus suggests that geochemical analyses of preserved, seismically exhumed limestone fault scarp rocks may help deciphering the seismic slip history of a fault.

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## 1. Introduction

If a complete catalog of the major earthquakes ( $M > 6$ ) that have occurred on active faults over the last few tens of thousands years was available, we would know how such earthquakes have repeated in the past, so that anticipating how similar strong earthquakes may occur in the future would become possible. Unfortunately such catalog does not exist and never will. As a result, our only chance to learn how major earthquakes have repeated in the past is to analyze their surface geological record, if any. Yet methods to analyze that record are

currently scarce and costly. The most common approach, paleoseismology, consists in digging few m-deep trenches across active faults (e.g., Daëron et al., 2005; Haibing et al., 2005; Kumar et al., 2006), which often reveals logistically difficult. Besides, while paleoseismology allows identification and dating of the few last major earthquakes on the faults, it more rarely allows determining the displacements produced by those earthquakes. Another, more recent approach focuses on faults having an escarpment at surface, and consists in determining the exposure time of that escarpment through the measurement of its content in *in situ* cosmogenic <sup>36</sup>Cl (e.g., Zreda and Noller, 1998; Mitchell et al., 2001; Benedetti et al., 2002, 2003; Palumbo et al., 2004). The method is powerful, yet extremely costly.

Our aim is to explore a new method that may help recovering past earthquake histories in a simpler and less expensive way than presently possible. We focus on normal faults, for those are the only ones to provide a complete geological record of their past major ruptures. Indeed, the action of a large earthquake on a normal fault is to uplift its footwall, so that a portion of the fault plane becomes

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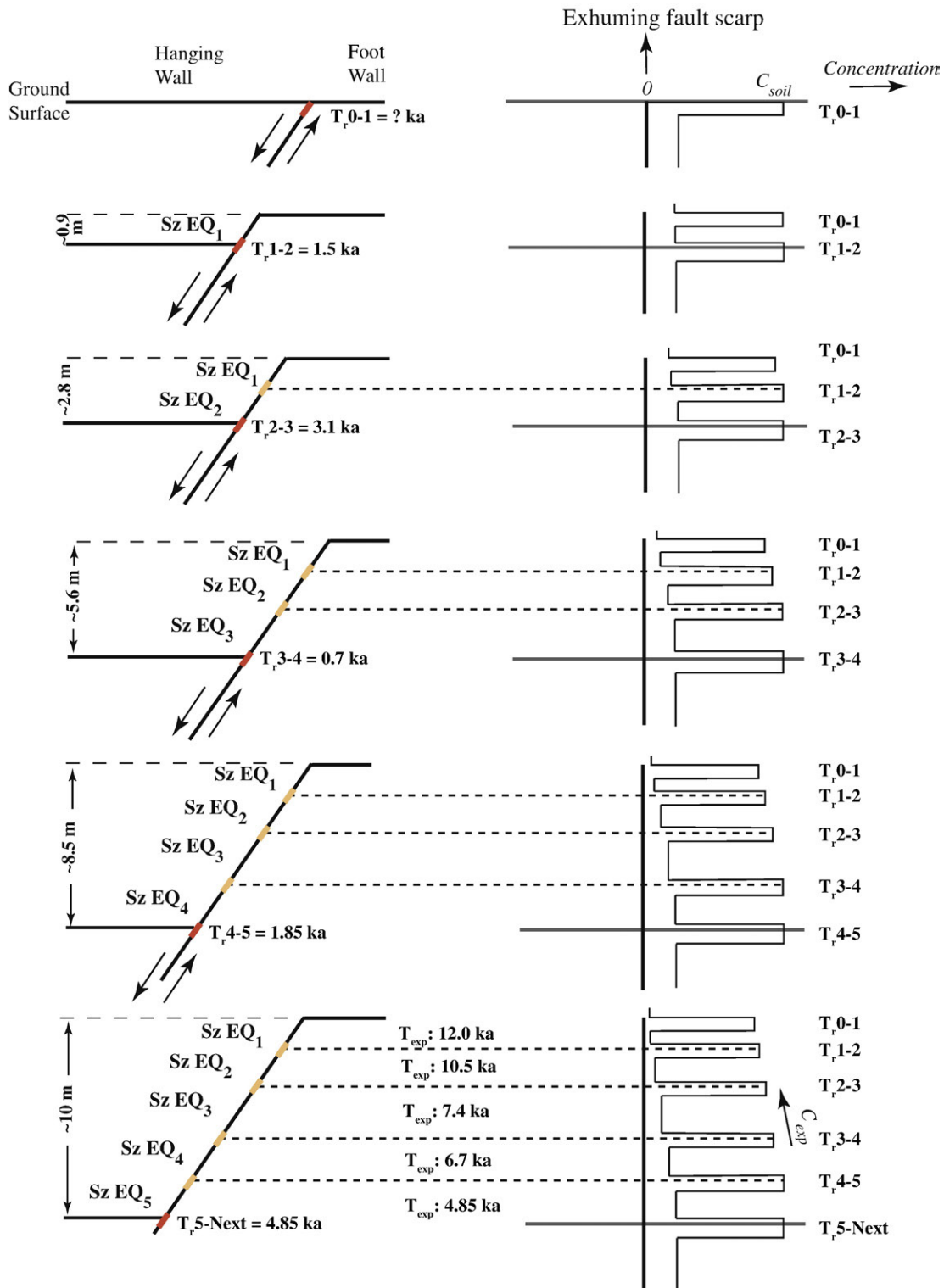
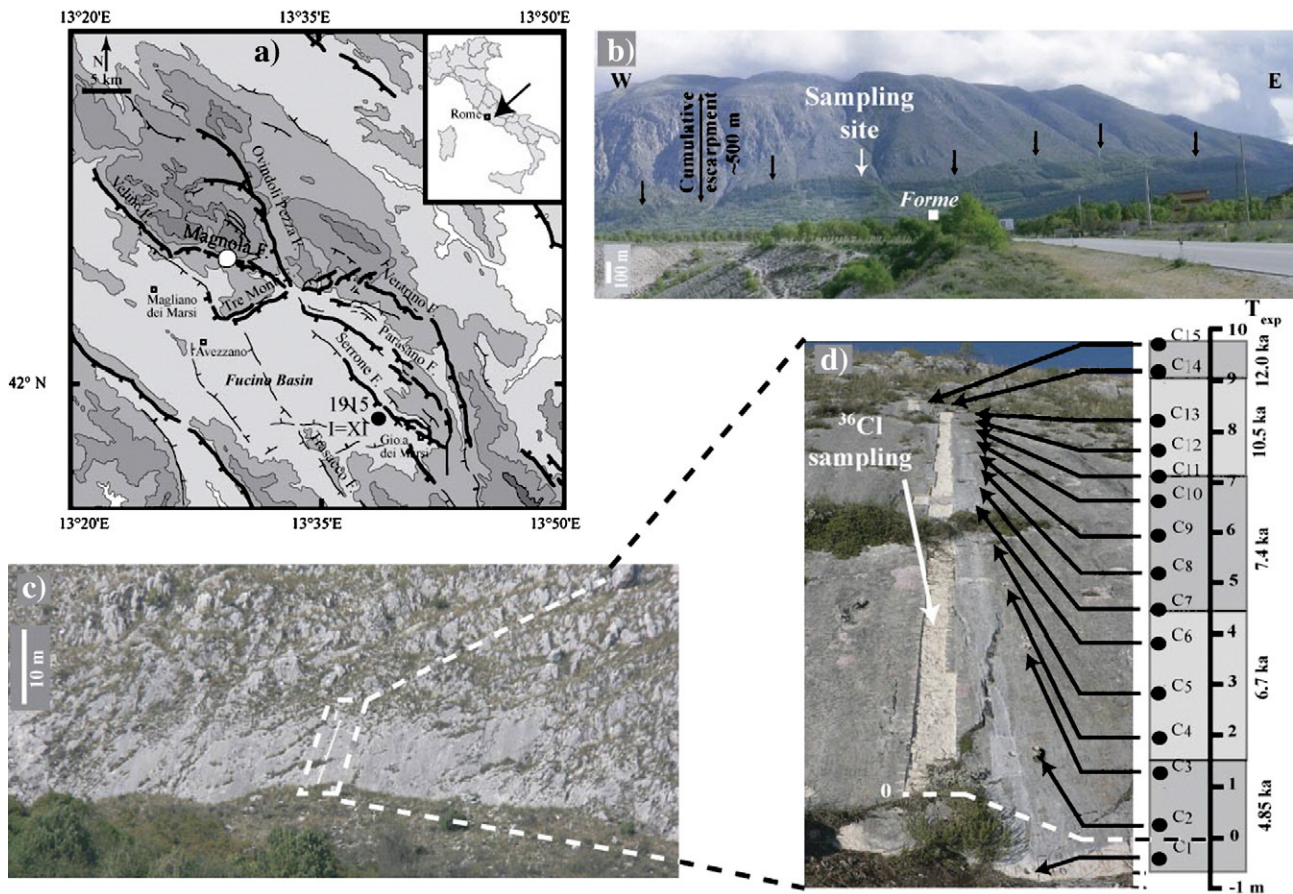


Fig. 1. Formation of an exhumed normal fault scarp through the repetition of major earthquakes (left panel), and upscarp evolution of chemical element concentrations observed in present study (right panel). Sz EQ<sub>n</sub> for slip zone exhumed by earthquake  $n$ . The five major earthquake-exhumed slip zones identified by Palumbo et al. (2004) on the Holocene Magnola fault scarp are reported, with the corresponding ages of exhumation ( $T_{exp}$ ). The transition zones between earthquakes  $n$  and  $n'$  are indicated (color patches and horizontal dashed lines) with the corresponding time of residence ( $T_{r, n-n'}$ ) of the scarp rocks in the upper soil. See text and Fig. 9 for explanation of the right panel (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

exposed to the atmosphere (Fig. 1, left panel). Hence, as large earthquakes repeat on a normal fault, the fault plane is progressively exhumed. The earthquake succession eventually leads to the formation of a high topographic escarpment, all the higher as the earthquake history is longer (Fig. 2b). When the offset rocks are

limestones, the youngest parts of such escarpments are commonly preserved in the form of ~10 m high scarps running at their base (Fig. 2b-c). Those scarps basically represent the section of the fault plane that has been exhumed by the few last major earthquakes on the fault.



**Fig. 2.** Sampling site on the Magnola fault. a) General tectonic setting (redrawn from Palumbo et al., 2004). Only major active faults are reported, all normal. Grey colors are for topography. The 1915 Fucino earthquake is indicated. The dot on the Magnola fault locates the sampling site. b) Northward view of the Magnola cumulative escarpment. The Holocene limestone scarp is indicated with black arrows. The Forme village is indicated. c) Closer view of the Holocene limestone scarp at the site of sampling. The white line is the vertical groove sampled by Palumbo et al. (2004) for  $^{36}\text{Cl}$  exposure dating. d) Closer view of the  $^{36}\text{Cl}$ -sampled site, together with the position of the cores (shown by the black arrows) sampled and analyzed in the present study. The upscarp positions of the cores are reported relative to the bottommost sample of Palumbo et al. (2004), taken as the zero. The grey zones indicate the major exhumed slip zones identified by Palumbo et al. (2004), with their exposure ages ( $T_{\text{exp}}$ ). Note that the scarp surface shows noticeable differences from base to top. The very base of the scarp appears as a white, 20 cm high strip that contrasts with the grey color of the rest of the scarp. On its bottom first 4–5 m, the scarp surface is smoothed with little evidence of disruption and weathering. Further above, it becomes rougher and locally dissected by karstic fissures, vertical flutes, and pits whose density increases with height. Above a height of  $\approx 10$  m the surface of the scarp is dismantled into chaotic blocks.

We focus on one of these limestone scarps, that of the Magnola fault in Central Apennines, Italy. Magnola is one of the main faults of the Fucino system, part of which ruptured in a devastating,  $M=7$  earthquake in 1915 (Fig. 2a; Odone 1915, Michetti et al., 1996, Galadini and Galli 1999). The Holocene activity of the Magnola fault is expressed by a continuous limestone scarp that runs at the base of its 500 m-high cumulative escarpment (Fig. 2b–c). The Magnola fault is an ideal target to explore a new method for past earthquake recovering, since its Holocene seismic slip history is already known. This slip history was derived from top to base-continuous *in situ*  $^{36}\text{Cl}$  cosmic ray exposure dating of the scarp surface, at one spot of the fault length (Palumbo et al., 2004, Fig. 2c–d). The  $^{36}\text{Cl}$  concentration distribution was best explained by five successive major earthquake exhumations over the last  $\approx 12$  ka, with slips between 1.5 and 3 m, separated by 1–3.1 ka-long periods of quiescence (Fig. 2d). Though powerful to recover seismic slip histories, cosmic-ray exposure dating is extremely costly, thus cannot be used as a routine.

We here explore a new approach that may help overcoming these limitations. As a normal fault breaks in a large earthquake, a section of its plane is suddenly exhumed and exposed to supergene weathering; rain, wind, snow, vegetation, bacterial activity, etc... The longer the surface is exposed, the more intense may be the weathering. A normal fault scarp is thus made of a series of rupture zones having been

exposed, thus weathered, over significantly different time spans. We suspect that such differential weathering may be detected in the chemical content of the fault scarp rocks, and its 'signature' used as a base to identify the various rupture zones on the scarps.

We thus investigate the chemical content of 15 carbonate rock samples collected from the base to the top of the Magnola Holocene scarp, where Palumbo et al. (2004) earlier determined the seismic slip history of the fault (Fig. 2d). We analyze major and trace element concentrations in the samples, and examine the concentration changes as a function of scarp height, hence of exposure time. We find that the chemical composition of the host-rock primarily varies as a function of scarp height, with significant step-changes correlating with the major seismic slip events recognized by Palumbo et al. (2004). This preliminary work thus suggests that simple geochemical analyses of seismically exhumed fault rocks may help recovering the seismic slip history of the fault.

## 2. Seismotectonic Holocene history of the Magnola fault

The NW-striking, 15 km-long Magnola is one of the largest active normal faults of the Fucino system, which dissects the central part of the Apennine chain (Piccardi et al., 1999). Its 500 m-high cumulative escarpment is underlined by a continuous,  $\approx 10$  m-high,  $50^\circ$ -steep

**Table 1**  
Major and trace element concentrations measured in host rocks

	Average concentration (ppm)	Minimum uncertainty (%)	Average Uncertainty (%)	Maximum Uncertainty (%)	C1	C2	C2*	C3	C3'	C4	C4*	C5	C6
				Elevation(m)	-0.35	0.3	0.3	1.25	1.25	1.95	1.95	2.85	3.85
<i>Trace elements (ppm)</i>													
Li	0.356	0	5	11	0.413	0.303	0.311	0.374	0.349	0.384	0.371	0.400	0.380
Sr	144	0	2	3	156	163	161	147	152	137	140	138	135
Y	1.614	1	2	6	2.12	1.41	1.39	1.98	2.00	2.26	2.24	1.25	1.96
Zr	0.485	0	3	6	0.858	0.519	0.527	0.594	0.559	0.522	0.521	0.429	0.562
Nb	0.036	2	6	8	(0.150)	0.039	0.038	0.041	0.044	0.026	0.029	0.021	0.037
Cs	0.024	0	5	10	(0.059)	0.020	0.021	0.032	0.029	0.025	0.025	0.015	0.026
Ba	4.68	1	3	4	2.88	3.78	3.87	3.03	3.16	(16.4)	(16.6)	2.70	3.33
La	0.575	0	3	6	0.885	0.558	0.551	0.785	0.813	0.887	0.891	0.400	0.718
Ce	0.362	1	3	6	0.686	0.382	0.385	0.481	0.489	0.431	0.460	0.257	0.399
Pr	0.091	0	3	11	0.159	0.084	0.085	0.130	0.136	0.139	0.140	0.060	0.113
Nd	0.401	0	2	6	0.670	0.364	0.365	0.546	0.567	0.597	0.605	0.271	0.526
Sm	0.080	1	7	14	0.137	0.066	0.067	0.118	0.120	0.114	0.128	0.044	0.105
Eu	0.019	0	9	14	0.028	0.018	0.018	0.029	0.025	0.030	0.028	0.014	0.024
Gd	0.109	1	4	12	0.175	0.101	0.100	0.135	0.153	0.165	0.164	0.074	0.136
Tb	0.017	2	6	11	0.028	0.015	0.016	0.023	0.022	0.025	0.027	0.011	0.021
Dy	0.117	0	5	14	0.170	0.098	0.105	0.158	0.151	0.180	0.154	0.090	0.143
Ho	0.028	1	5	12	0.040	0.026	0.026	0.037	0.037	0.041	0.043	0.023	0.034
Er	0.088	1	5	11	0.127	0.079	0.078	0.116	0.112	0.127	0.132	0.072	0.106
Yb	0.077	4	8	12	0.108	0.073	0.066	0.092	0.099	0.121	0.106	0.064	0.091
Lu	0.013	2	7	16	0.017	0.012	0.011	0.016	0.016	0.016	0.015	0.010	0.016
Pb	0.282	0	3	6	0.264	0.236	0.246	0.228	0.227	0.300	0.293	0.255	0.262
Th	0.026	4	10	19	0.050	0.026	0.021	0.043	0.046	0.028	0.033	0.014	0.031
U	0.324	0	3	7	0.299	0.281	0.281	0.307	0.294	0.261	0.266	0.346	0.327
<i>Major element (%)</i>													
Al <sub>2</sub> O <sub>3</sub>	0.005	–	–	22	(0.024)	0.002	–	0.007	0.006	–	–	–	0.003
CaO	56.349	–	–	4.5	55.57	57.45	–	56.23	56.59	55.41	–	56.66	55.23
MgO	0.422	–	–	4	0.425	0.401	–	0.440	0.437	0.449	–	0.425	0.436
Na <sub>2</sub> O	0.013	–	–	9	0.013	0.010	–	0.013	0.012	0.015	–	0.013	0.013
K <sub>2</sub> O	0.004	–	–	5	0.006	0.003	–	0.005	0.004	0.004	–	0.003	0.004
Fe <sub>2</sub> O <sub>3</sub>	0.010	–	–	3	0.015	0.011	–	0.011	0.010	0.011	–	0.010	0.011
MnO	0.005	–	–	7	0.004	0.005	–	0.005	0.005	0.003	–	0.005	0.004
TiO <sub>2</sub>	0.001	–	–	10	0.003	0.001	–	0.002	0.002	0.002	–	0.001	0.002
P <sub>2</sub> O <sub>5</sub>	0.032	–	–	7	0.035	0.024	–	0.036	0.039	0.039	–	0.022	0.043

For trace elements, measurements and associated uncertainties are expressed in ppm, while they are in % for major elements. Measurements considered as outliers (see text) are in brackets. Measured concentrations close to 3–4 times the blank values are in italic. Hyphens indicate concentrations too low to be measured. Concentrations measured on standards BIR, BR24, UBN and BHVO-2 are indicated for information. Note that no data are reported for Rb, Ta and Hf because concentrations in all samples were always below the detection limit.

scarp, that has been shown to be the youngest, Holocene section of the seismically exhumed fault plane (Palumbo et al., 2004). Since the fault offsets Meso-Cenozoic carbonates, its slip surface, i.e., the fault plane, is made of carbonates as well. More specifically, the Holocene section of the fault plane appears as a 1–15 cm-thick layer of tough, recrystallized calcite locally embedding small limestone clasts, whose outer, exposed surface has been smoothed into a flat plane by the repeated fault slips (Fig. A-a in the Appendix; see description of similar faults in Agosta et al., 2007; Agosta and Aydin, 2006). The smooth fault plane overlies a thicker zone (a few m) of densely fissured, brecciated carbonates. The overall fault structure is thus similar to that of most active faults.

Palumbo et al. (2004) have determined the Holocene seismic slip history of the Magnola fault, at one central spot of its length (Fig. 2a). The slip history was traced using the distribution of in situ <sup>36</sup>Cl concentration versus height of the Holocene scarp, distribution best explained by five successive major earthquake exhumations having occurred 12,000, 10,500, 7400, 6700, and 4850 years ago (Fig. 1 and 2d).

Thus, at the site sampled by Palumbo et al. (2004), the Holocene scarp surface consists of a uniform and smooth layer of calcite which, from top to base, includes at least five distinct sections, each 1.5 to 3 m high, that have been exposed to supergene weathering over 12, 10.5, 7.4, 6.7 and 4.85 ka. We analyze the chemical

composition of rock samples taken from each of the five sections, and examine whether or not the composition varies as a function of exposure time.

### 3. Sampling and chemical processing

Using a portable driller, we collected fifteen cores (C1–15; Fig. A-b in the Appendix) along a vertical profile less than a meter away from the sampling groove of Palumbo et al. (2004) (Fig. 2c–d). The difficulty of the sampling (made by roping down the scarp) and the toughness of the carbonates forced us to limit the number of collected samples. The cores were sampled every 0.5–1 m from base to top of the scarp. Their upscarp position is defined with respect to the lowest sample collected by Palumbo et al. (2004).

The cores are overlain with a ~100 μm-thick dark coating that markedly contrasts with the white color of the host rock (Fig. A-b in the Appendix). We captured high-resolution images of the coating by Scanning Electron Microscopy (SEM) and made qualitative chemical analyses of that coating using a coupled Electron Dispersive Spectroscopy (EDS) Microprobe. Because the coating is too thin to be isolated from the underlying rock sample, we could not perform any detailed and accurate mineralogical analyses (such as X-ray diffraction). However, we report semi-quantitative major element analyses performed using a non-destructive μX-Ray Fluorescence spectrometry



C7	C7'	C8	C9	C10	C11	C11*	C12	C13	C13*	C14	C15	BIR	BR24	UBN	BHVO-2
4.45	4.45	5.25	5.95	6.65	7.15	7.15	7.65	8.2	8.2	9.2	9.7				
0.324	0.325	0.259	0.149	0.358	0.383	0.365	0.438	0.337	0.378	0.449	0.388	3.49	6.68	28.80	4.84
150	147	151	139	130	130	133	159	135	135	137	146	116	6.84	588	413
1.56	1.53	1.57	1.12	1.72	1.55	1.65	1.03	0.79	0.78	1.51	2.35	18.0	2.57	31.1	29.3
0.464	0.448	0.376	0.411	0.433	0.490	0.499	0.382	0.348	0.347	0.400	0.501	16.0	3.34	311	186
0.033	0.031	0.019	0.022	0.025	0.034	0.031	0.028	0.009	0.015	0.025	0.026	0.593	0.045	41.8	19.8
0.018	0.018	0.027	0.030	0.026	0.024	0.022	0.014	0.005	0.005	0.014	0.028	0.007	11.3	0.461	0.101
3.02	2.89	2.59	4.35	2.34	1.95	1.90	1.98	1.79	1.82	3.00	(16.9)	6.58	26.00	385	131
0.545	0.547	0.498	0.317	0.577	0.456	0.474	0.324	0.223	0.209	0.522	0.907	0.610	0.331	33.3	15.1
0.351	0.330	0.328	0.233	0.344	0.311	0.314	0.275	0.156	0.148	0.344	0.440	1.94	0.812	72.6	37.8
0.084	0.086	0.070	0.051	0.091	0.075	0.075	0.048	0.033	0.030	0.083	0.141	0.374	0.119	9.36	5.32
0.367	0.374	0.322	0.214	0.409	0.331	0.332	0.198	0.143	0.135	0.375	0.674	2.41	0.603	39.3	24.4
0.082	0.073	0.066	0.044	0.081	0.076	0.075	0.039	0.031	0.027	0.074	0.125	1.10	0.217	8.20	6.07
0.020	0.018	0.015	0.010	0.021	0.017	0.020	0.008	0.006	0.005	0.021	0.032	0.514	0.078	2.45	2.01
0.102	0.103	0.085	0.072	0.108	0.095	0.102	0.063	0.041	0.042	0.112	0.164	1.842	0.311	7.18	5.98
0.014	0.015	0.015	0.011	0.017	0.015	0.013	0.007	0.006	0.007	0.015	0.027	0.370	0.060	1.08	0.952
0.110	0.111	0.111	0.075	0.120	0.110	0.114	0.062	0.047	0.046	0.113	0.180	2.53	0.400	5.55	5.18
0.026	0.027	0.025	0.020	0.029	0.025	0.029	0.015	0.011	0.012	0.027	0.043	0.580	0.094	1.04	0.987
0.088	0.078	0.079	0.059	0.094	0.083	0.093	0.055	0.044	0.043	0.074	0.122	1.73	0.289	2.65	2.55
0.070	0.075	0.074	0.047	0.077	0.076	0.069	0.046	0.035	0.037	0.077	0.109	1.65	0.293	2.05	1.99
0.013	0.013	0.015	0.008	0.013	0.013	0.013	0.008	0.006	0.007	0.013	0.017	0.254	0.046	0.293	0.284
0.200	0.206	0.227	(0.755)	0.240	0.280	0.282	0.222	0.202	0.214	0.261	0.290	2.86	12.73	3.47	1.55
0.024	0.021	0.017	0.019	0.026	0.029	0.028	0.015	0.012	0.011	0.022	0.032	0.030	0.064	4.73	1.25
0.316	0.300	0.288	0.349	0.392	0.427	0.432	0.268	0.218	0.235	0.525	0.262	0.012	0.063	1.21	0.408
-	-	-	-	0.002	0.004	-	-	-	-	0.001	0.001	-	-	-	-
56.87	55.89	56.91	55.87	55.55	57.08	-	57.31	56.89	-	55.45	57.07	-	-	-	-
0.422	0.430	0.358	0.372	0.389	0.443	0.464	0.418	0.418	0.418	0.425	0.463	-	-	-	-
0.014	0.013	0.010	0.008	0.013	0.015	0.018	0.014	0.014	0.014	0.016	0.017	-	-	-	-
0.004	0.004	0.003	0.003	0.003	0.004	0.003	0.003	0.002	0.002	0.004	0.004	-	-	-	-
0.009	0.009	0.011	0.010	0.011	0.013	0.007	0.007	0.007	0.007	0.011	0.011	-	-	-	-
0.006	0.006	0.005	0.006	0.006	0.006	0.005	0.006	0.006	0.006	0.010	0.004	-	-	-	-
0.001	0.001	0.001	0.001	0.001	0.001	0.001	0.001	0.001	0.001	0.001	0.002	-	-	-	-
0.029	0.031	0.022	0.047	0.036	0.033	0.025	0.017	0.017	0.017	0.032	0.040	-	-	-	-

( $\mu$ XRF, Eagle<sup>®</sup> III EDAX; energy of 20 keV). These analyses were obtained on the coatings (spot diameter of 300  $\mu$ m) and on the underlying rock samples (spot diameter of 30  $\mu$ m) at increasing distances from the surface (first analysis at 0.1 cm depth, then every 0.5 cm from surface to bottom; Fig. A-b in the Appendix). Results are given in Table A (Appendix). Because measurements were homogeneous in the deepest parts of the cores (i.e., below 0.1 cm), only few cores were completely analyzed.

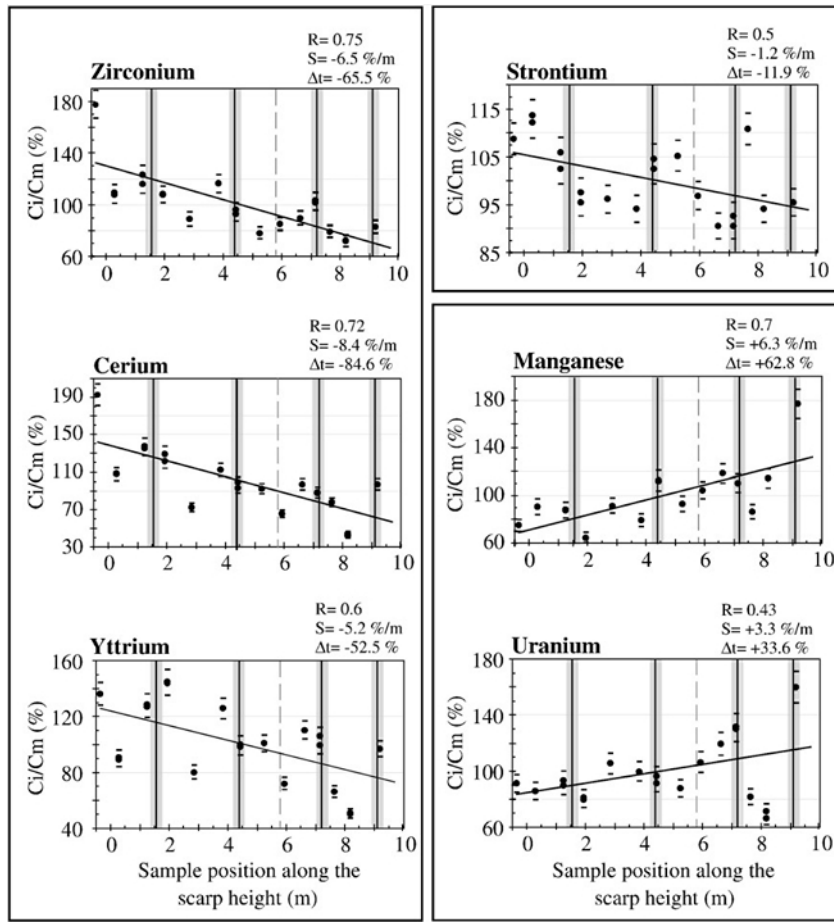
Accurate major and trace element contents were measured in the host rock of all cores (Table 1). The first 2 mm of each core was removed to eliminate the surface coating. Measurements were done on sub-samples collected in the first 1.5 cm of the coating-free cores (Fig. A-b in the Appendix). Trace elements were measured by ICP-MS (VG<sup>®</sup> PQ turbo +) using the procedure described by Marini et al. (2005) but slightly modified to eliminate the carbonates before whole rock dissolution. Measurements were corrected for machine drift using a multispikes (Be, As, In, Tl and Bi) and calibration of the signal was done using repeated analyses of the reference material BHVO-2. The precision on measurements can be evaluated from the analysis of the international rock standards UBN and BIR-1 and the in-house standard BR24. Because UBN and BIR-1 have very low trace element contents, the precision is not very good but usually better than 10% (Table 1). Analytical uncertainties are estimated by comparing results obtained on multiple data acquisition for the same sample and complete duplicate analyses of selected samples (Table 1). Those uncertainties are  $\approx$  5% for most elements, lower

than 15% for some elements and reach higher values ( $\approx$  20%) for the few elements whose concentration is hardly higher than 4 times the blank values (Th, Lu). In the following, we propagate the maximum uncertainty obtained for each trace element to the whole series of measurements of that element. Major element concentrations were determined from aliquots of the solution used for trace element analyses (Table 1, concentrations expressed as oxides). Analyses were carried out using a ICP-AES (Perkin Elmer<sup>®</sup> Optima 3300 DV). Calibration was done using Fluka<sup>®</sup> standard solutions for atomic spectrometry. The analytical uncertainties are estimated from three repeated measurements of each sample. The maximum uncertainties range between 3 and 10% for all elements except Al<sub>2</sub>O<sub>3</sub> whose concentrations are very low. Here again, we propagate the maximum uncertainty obtained for each element to the entire set of measurements of that element.

In Table 1, all measured concentrations, even when they approach the detection limit, are listed. In brackets are the concentrations that are more than double the average concentration of the entire sample collection; we consider these values as outliers, and as such, we ignore them in all calculations presented below.

#### 4. Data analysis

The Magnola scarp rock consists of a very homogeneous, low magnesium carbonate (56 $\pm$ 0.6 % CaO and 0.40 $\pm$ 0.05% MgO; Table 1).



**Fig. 3.** Concentrations of Zr, Ce, Y, Sr, Mn, and U, as a function of scarp height. In all plots, X is the on-fault position of the samples increasing upward and Y is the concentration of sample *i* (Ci) normalized by the mean concentration (Cm) averaged over the whole sample collection, in percent. Solid vertical lines represent the transitions between the major earthquake slip zones identified by Palumbo et al. (2004), with their position uncertainties in grey (uncertainties on the on-scarp position of the transitions are ±25 cm on average according to the Fig. 5c of Palumbo et al., 2004). The thin dotted vertical line locates a possible additional earthquake suggested by Palumbo et al. (2004). The slanted straight lines are the regression curves calculated from all measurements (but those in C15 sample). When several measurements were done for the same sample, only the average was used for the regression calculation, but all measurements are plotted in the figure. Concentrations are given in Table 1 and regression parameters shown on the plots are given in Table C-b (Appendix).

The trace element compositions define patterns parallel to that of Reference Shales (Sholkovitz, 1988; Fig. B-a in the Appendix), but with much lower and relatively constant contents related to the high carbonate content (Plank and Langmuir, 1998). The seawater origin of the Magnola carbonates is well preserved in the samples, as demonstrated by the systematic negative Ce anomalies (Fig. B-b in the Appendix).

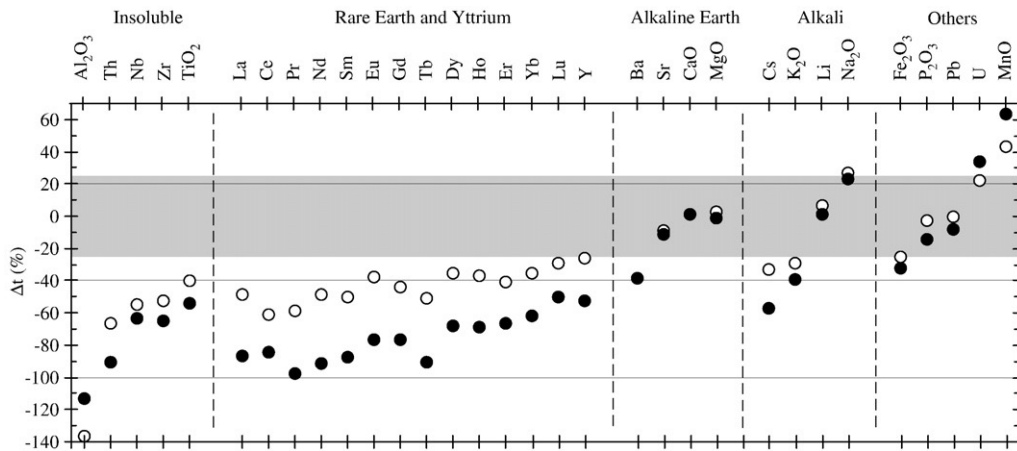
#### 4.1. Evolution of host rock chemical content as a function of scarp height

To examine how the major and trace element concentrations measured in the host rock vary as a function of scarp height, hence as a function of exposure time, we built a series of figures that we will first describe before analyzing them.

First, because of the large number of measured elements, we calculate regression coefficients between element pairs to highlight possible common behaviors. The regression coefficients are reported in Table B-a (Appendix). Expected affinities appear clearly, such as those linking the REE together and to Yttrium, and those linking the insoluble elements together. Thus, in the following, we may discuss the behavior of the REE-Y and insoluble elements by restricting our analysis to only one well-constrained element representative of each group.

Second, we built a series of plots (Fig. 3 and Fig. C-a in the Appendix) documenting the behavior of each element as a function of scarp height. In all these plots, the element concentration in each

sample (Ci) is normalized to the average concentration of that element in the whole sample collection (Cm), and plotted as a function of scarp height. Error bars on normalized concentrations are calculated using the maximum uncertainties given in Table 1. The transitions separating the major earthquake slip zones identified by Palumbo et al. (2004) are reported as vertical lines, with their on-scarp position uncertainties (grey bands). Then, for each element, a linear regression is calculated as a function of scarp height (straight line). The regression parameters are reported in Table C-a in the Appendix. Finally, Fig. 4 shows the total base-to-top concentration change values (Δt) obtained for all elements (empty dots). Maximum uncertainties on concentration measurements being ±25% in the worse case, any number falling within the ±25% stippled band indicates no significant concentration variation with scarp height. Figs. 3, 4 and Fig. C-a in the Appendix show that most element concentrations vary as a function of scarp height. The total concentration changes from base to top are significant with values up to 65% (Fig. 4, empty dots; Al<sub>2</sub>O<sub>3</sub> not well constrained). Though irregular in detail, the overall base to top evolutions are generally satisfactorily averaged by linear trends with regression coefficients up to 0.6 (Table C-a in the Appendix). When the topmost and oldest C15 sample is excluded from the calculations, those coefficients significantly increase (up to 0.75) (Table C-b and Fig. C-b in the Appendix). We suspect that C15 is not representative of the original scarp composition because of the significant mechanical disruptions that alter the crest of the scarp where C15 was sampled. We thus



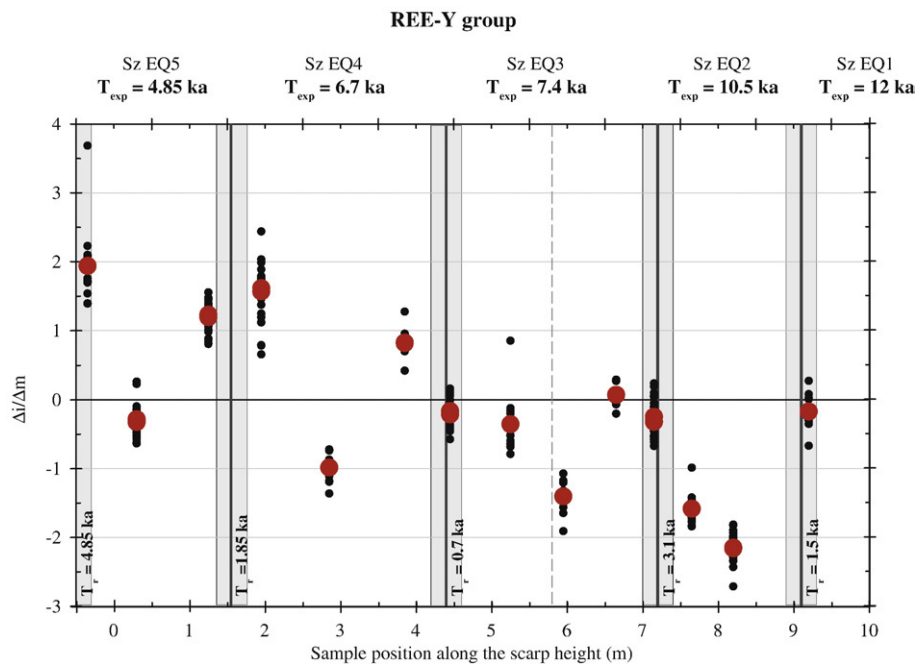
**Fig. 4.** Overall concentration changes from base to top of the scarp. Elements are listed by chemical affinities. The Y-axis is the total concentration change ( $\Delta t$ , in %) inferred from the regression calculations (Table C in the Appendix). Empty and solid dots are  $\Delta t$  values from Tables C-a and C-b in the Appendix, respectively. Changes smaller than  $\pm 25\%$  (grey horizontal field) are considered insignificant.

exclude C15 from the analysis. The recalculated parameters are shown in Figs. 3, 4 (solid dots) and Table C-b in the Appendix.

Three distinct overall behaviors are distinguished. About 22% of the elements (Pb, Sr, Ca, Mg, P, Na, Li) do not show any significant variation along the scarp height; their concentrations change by less than 25% from base to top. Note that among these 'stable' elements are those most abundant in the carbonate host rock. The vast majority of elements (72%) display a significant concentration decrease, between 30 and 100%, as one goes from the base to the top of the scarp. These 'decreasing' elements include the REE, both light and heavy (LREE and HREE, respectively), together with Yttrium, the insoluble elements, some of the alkaline (K, Cs) and alkaline Earth elements (Ba), iron (Fe), and aluminum (Al; though not well-constrained). The decrease of the LREE is larger than that of the HREE. Finally, the remaining 6% of the

elements, namely Mn and U, show an opposite behavior, with concentrations increasing by 30–60% from base to top of the scarp.

A closer inspection of the concentration versus height plots (Fig. 3 and Fig. C-b in the Appendix) shows that, though the element concentrations follow a roughly linear trend when considered overall, they also display peaks that significantly depart from the average trend. A way to quantify those peaks is to measure them with respect to the average concentration of the considered element in the whole scarp ( $C_m$ ); the peak amplitude thus is the difference  $\Delta i$  between  $C_i$  (concentration in sample  $i$ ) and  $C_m$ . Since we want to compare the peak amplitudes for a large number of elements, we need to normalize them so that they eventually fall in a similar range of values. We choose here to normalize each peak amplitude  $\Delta i$  by the mean absolute value calculated from all  $\Delta i$  (for a given element,



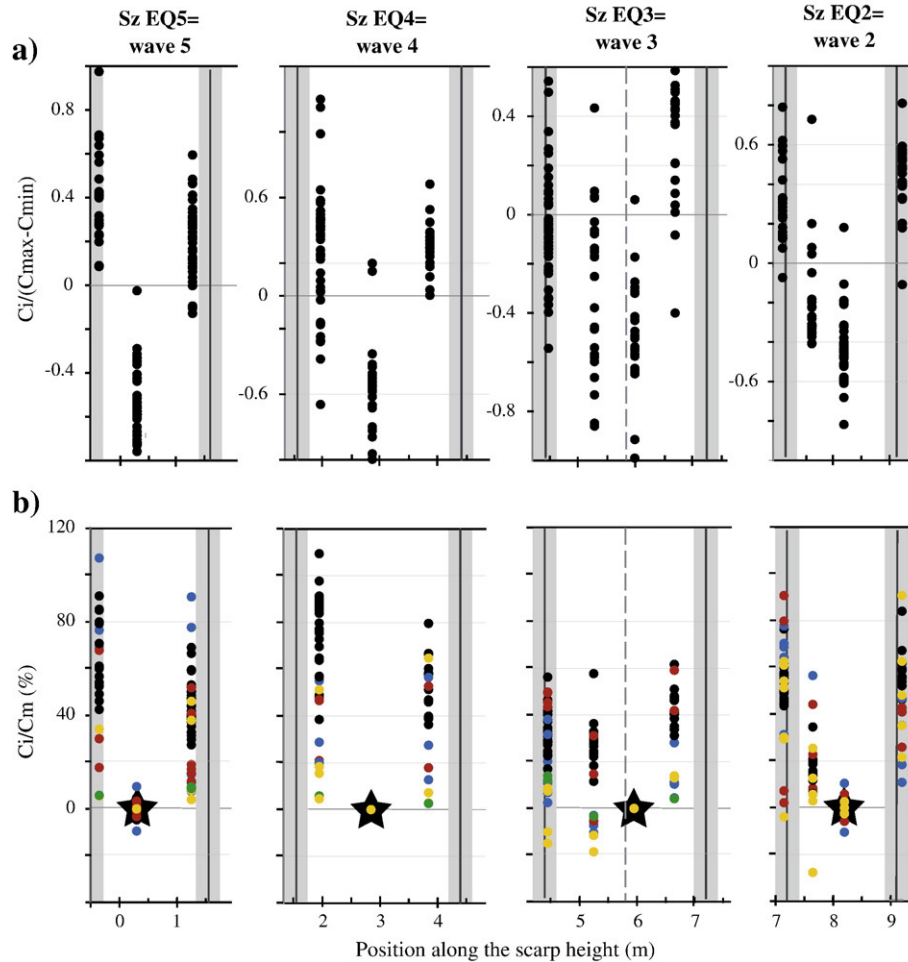
**Fig. 5.** Concentration variability of the REE-Y group as a function of sample position along the fault scarp. Transitions between slip zones are shown as in Fig. 3. The Y-axis measures the difference  $\Delta i$  between a given concentration in core  $i$  ( $C_i$ ) and the mean concentration over the whole core collection ( $C_m$ ), normalized to the mean absolute value of those differences ( $\Delta m = \sum |\Delta i| / N$  with  $N$  the number of measurements); Y thus plots  $\Delta i / \Delta m$  (that is  $(C_i - C_m) / \Delta m$ ). The black dots are individual element measurements while the red dots correspond to the average behavior of all the REE-Y. Sz EQ $i$  for slip zone exhumed by earthquake  $i$ .  $T_{exp}$  is exposure time while  $T_r$  is residence time in upper soil (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

$\Delta m = \sum |\Delta i| / N$  with  $N$  the number of measurements). Fig. 5 shows the result for the REE-Y group. It first confirms that the REE and Y (black dots) behave in a similar manner. Their peak amplitudes can thus be averaged (red dots). Plotted as a function of scarp height, the red dots define sort of a 'wavy' curve, with maxima 'crest-to-trough' amplitudes up to 100% of the mean concentration  $C_m$  (see Fig. 3 and Fig. C-b in the Appendix for values in %). The crest values decrease overall as one goes upscarp, to reach negative values at its top. All concentration maxima are found at or close to the transition zones that separate the successive earthquake slip surfaces, while all concentration minima are roughly in the middle of the slip surfaces. Fig. 6 shows that the vast majority of analyzed elements have the same 'wavy behavior' as the REE-Y group. Fig. 6a examines whether the various elements show a similar variation pattern within a given slip zone. It is thus designed to highlight the shape of the 'waves'. Note however that data are not dense enough to constrain the precise shape of the so-called 'waves', so that the term 'wave' is only used in a general sense. Fig. 6b shows the amplitude of the concentration changes, hence of the 'waves', depicted in Fig. 6a. The way these two figures are built is described in the caption. Fig. 6a shows that most element concentrations vary similarly both within each slip zone and from one slip zone to the other.

In any slip zone, the concentrations draw a concave-up 'wave' centered at about the half-slip zone height, with maxima at or close to both edges of the slip zone. Most elements show that pattern repeated in all slip zones (REE-Y group, Zr, Th, K, Ti, Nb), or in most of them (Cs, U, Mg, Na, P, Pb, Fe, Mn, Li). Fig. 6b shows that in any slip zone the crest-to-trough amplitudes of the waves are large for most elements, reaching up to 100% of  $C_m$ . The elements with the largest variations (>40%  $C_m$ ) in all or most slip zones include the REE-Y group, Th, Cs, and P.

#### 4.2. Evolution of the coating composition with scarp height

The SEM images of the core surfaces (Fig. 7) suggest that the roughness of the exposed fault plane increases as one goes upscarp; while striated small spots are preserved in the bottommost sample (Fig. 7j), none is found further upward where the sample surfaces rather show an increasing density of microscopic pits (Fig. 7a–c–e–g). At a smaller scale, the pitted samples appear coated with a large number of spherical and/or flattened features  $\sim 5 \mu\text{m}$  in size, most entrapped in pits (Fig. 7f), and with  $>10 \mu\text{m}$ -long filaments (Fig. 7d–h). The shape and size of the former features suggest that they may be algae (Smith et al., 2000), while the filaments may be rhizines of



**Fig. 6.** Concentration oscillations within seismic slip zones. In both a and b, each slip zone is treated independently so that each plot includes only the data pertaining to this slip zone. a) Common pattern of the concentration oscillations. Elements are in black. The X-axis and transition zones are as in Fig. 3. Sz EQ*i* for slip zone exhumed by earthquake *i*. The Y-axis plots the concentrations  $C_i$  normalized by the crest-to-trough amplitude of the wave ( $C_{\max} - C_{\min}$ ), with respect to the regression curve calculated from all measurements (but C15). The concentrations vary with the same pattern in each slip zone. The detailed list of elements plotted is given in Table E (Appendix). b) Amplitudes of the concentration oscillations. Colors distinguish the elements using their chemical affinities: blue, insoluble; black, REE-Y; green, alkaline earth; red, alkali; yellow, others. The X-axis and transition zones are as in Fig. 3. The Y-axis plots the concentrations  $C_i$  normalized by the mean concentration  $C_m$ , with the value considered as the trough of the wave being forced to zero for all elements (star; note that it is not necessarily  $C_{\min}$ ). The amplitudes of the concentration oscillations are large for most elements. Their values are reported in Table E (Appendix) (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

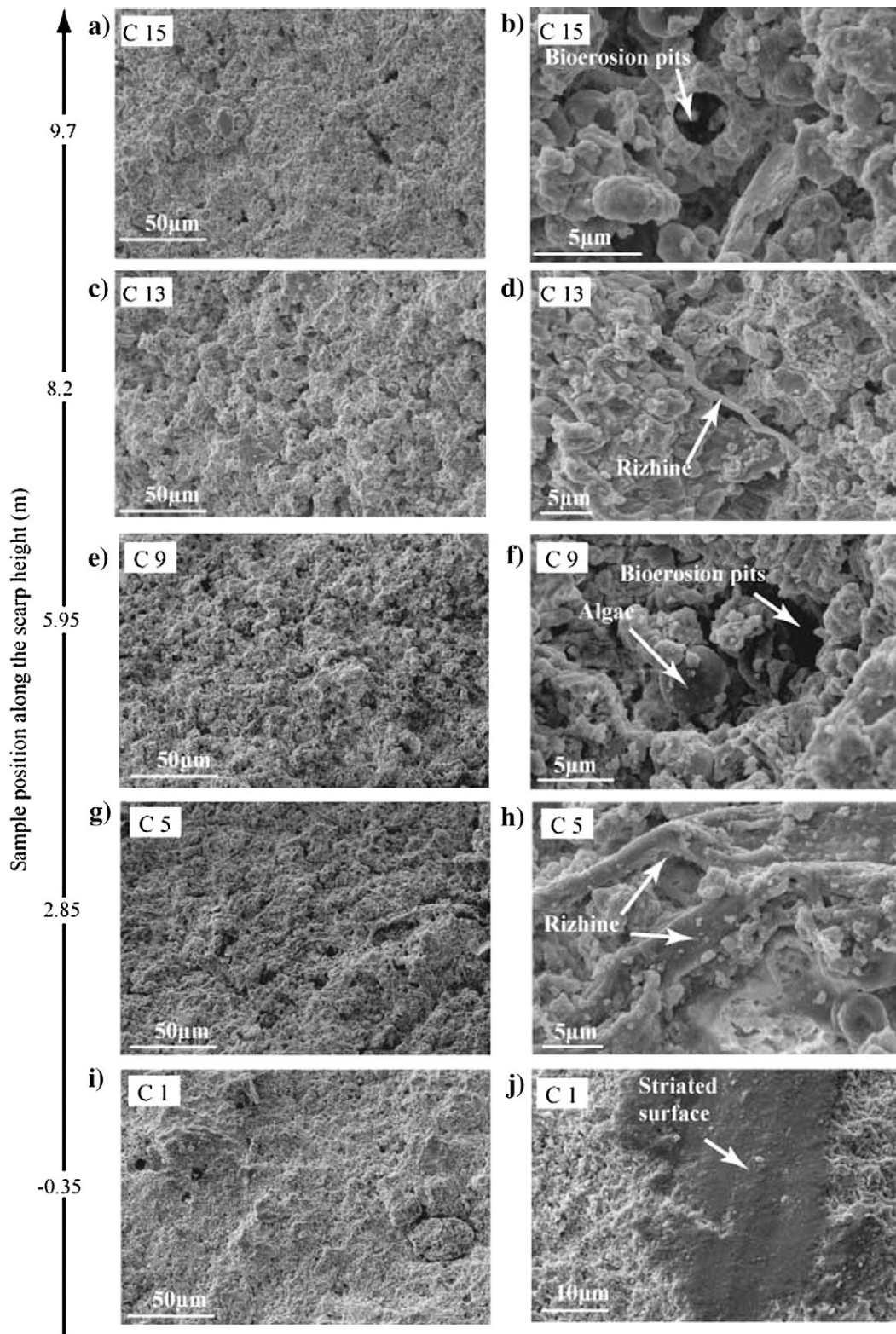


Fig. 7. SEM images of the core surfaces. Panels a, c, e, g, and i show general views of the surfaces, while panels b, d, f, h and j show close-up.

lichens. The close proximity of the rounded algae and the pits (Fig. 7f), the peculiar shape of the pits – as rounded as the algae (Fig. 7b–f), and the apparent penetration of the rhizines into the rock (Fig. 7d–h), combine to demonstrate the presence of organic compounds on (epilithic) and within (endolithic) the rocks. The EDS analyses of the surface coating reveal the presence of non-carbonated particles, such

as scarce micrometer-scale iron sulphides (Fig. D-a in the Appendix) and decamicrometer-scale aluminum-rich silicates (Fig. D-b in the Appendix). These particles contrast with the carbonate composition of the host rock suggesting that part of the coating material is not derived from the host rock; external agents thus likely contribute to the coating formation. The semi-quantitative major element analyses

performed by  $\mu$ XRF on the coatings are given in Table A (Appendix) and are markedly different from the values obtained for the same elements in the host rocks. Even though the coatings consist mainly of CaO, they also contain a significantly higher proportion of SiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub>. In Table D (Appendix), we provide correlation coefficients between the various elements measured in the coating, and the coefficients highlight significant relationships between SiO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>, and Fe<sub>2</sub>O<sub>3</sub>, but also between K<sub>2</sub>O and SO<sub>3</sub>. Together with the microscopic observations and the EDS analyses, these correlations confirm that the coating is partially made of aluminum-rich silicates and iron-sulphides. These could be a combination of various minerals including clays and jarosite. When plotted as a function of scarp height, the major element contents of the coatings do not show any specific pattern or trend, and most of them remain constant despite local variability.

## 5. Discussion

### 5.1. Overview of carbonates chemical weathering processes

The carbonate weathering processes have been studied in considerable detail, particularly because of the ubiquity of calcite in the sub-surface environments (e.g., Sweeting, 1966; Trudgill, 1985; Smith et al., 2000; Plan, 2005). The most critical property of calcite is its highly reactive nature; the calcite bulk equilibrium is weak and many factors may disrupt it, leading to calcite dissolution or degradation.

Calcite dissolution occurs under the action of agents such as rain, moisture, acidic fluid circulations, and/or biogenic features (e.g., Trudgill, 1985; Budd et al., 1995; White and Blum, 1995). Any fluid containing carbon dioxide has the ability, once in contact with calcite, to dissolve it, even at the microscale. Significant fluid circulations occur at shallow depth (<1–2 km), where the carbonates are buried before they are possibly exhumed and those fluids are harmful to limestone, particularly in the sub-surface where they reside and flow through an acidic soil. Indeed, the shallow soil, and particularly its 30–50 first cm, is a zone of intense acidification due to the high density of microorganisms, humus material and plant roots. Such materials produce large amounts of CO<sub>2</sub> through processes such as root respiration, oxidation of organic material and microbial respiration (e.g., Reardon et al., 1979; Paul and Clark, 1989; Fierer et al., 2003), which lead to relatively low, thus acidic soil–water pH. Additional bacterial acidification also commonly occurs (Le Métayer-Levrel et al., 1999). Altogether these processes contribute to dissolve the carbonates. Once exposed to free air, the carbonates experience additional chemical weathering under supergene agents such as rain (pH ≈ 5.7), moisture or biogenic features (such as lichens, bacteria, etc...; Danin et al., 1983; Chen et al., 2000; White, 2000; Burford et al., 2003). This weathering leads to further dissolution of the calcite. Because biotic weathering and internal fracturing may also occur and contribute to open paths into the rock, the calcite dissolution front may go deep into the host rock, at least several cm (e.g., Le Métayer-Levrel et al., 1999).

Once calcite dissolution has occurred, dissolution products may either re-precipitate or be leached away. In the former case, while re-precipitating, the calcite may incorporate impurity elements in its new lattice. Elements most favorably incorporated into a re-precipitating calcite lattice are those with an ionic radius similar to, and a valence compatible with, that of Ca<sup>2+</sup>, because they can substitute to Ca and produce minimum spatial disorder in the calcite lattice. Those elements include Sr, Pb, Mg and P (e.g., Mucci and Morse, 1983; Pingitore and Eastman, 1986; de Leeuw, 2002), and possibly Co, Cd, Ba, Zn (e.g., Reeder et al., 2001; de Leeuw, 2002). Under reducing conditions, U and Mn can also be incorporated into a calcite lattice. Indeed, Mn<sup>2+</sup> can occupy an on-center position within the calcite lattice, with spatial disorder not significantly different from the host Ca<sup>2+</sup> (Cheng et al., 2001; Sturchio et al. 1998). The uptake of uranium

by calcite is more complex and depends on the uranium valence. In oxidized surface environments, uranium exists as U<sup>6+</sup> and U<sup>4+</sup>, U<sup>6+</sup> being usually much more soluble and mobile than U<sup>4+</sup> (Langmuir, 1997; Elzinga et al., 2004). U<sup>6+</sup> behaves roughly as Mn and is easily incorporated in the calcite structure (Sturchio et al., 1998; Reeder et al., 2001; Kelly et al., 2003). U<sup>4+</sup> forms in reducing conditions, which prevail in the presence of organic matter and/or of specific bacteria (Lovley et al., 1991; Anderson and Lovley, 2002; Ohnuki et al., 2005). In that case, the U<sup>4+</sup> ion occupies the Ca position with minimal distortion and has thus a stable location in calcite (Lovley et al., 1991; Sturchio et al., 1998; Reeder et al., 2001). U<sup>4+</sup> can therefore be integrated within the calcite structure for very long times. Usually, the incorporation of U into calcite is coupled with that of Na (Kelly et al., 2003; Sturchio et al., 1998), though in smaller proportions. Any disruption of the calcite equilibrium leading to calcite dissolution and re-precipitation in the presence of the elements listed above likely results in some element–calcite solid solution formation, so that these elements are incorporated into the re-precipitated calcite lattice and remain in the final material. On the other hand, this is not the case of other elements such as Al, REE and high-field strength elements (Nb, Th, Zr, Ti), which are not easily accommodated in the calcite structure and are most probably trapped at grain boundaries. These elements are released during calcite dissolution, but are not accepted in the re-precipitating structure. The net result is a loss of these elements in the final calcite. Note that, in such carbonate environment, the process that governs the trace element behavior differs from that most commonly found in alteration profiles such as lateritic profiles where the elements least affected by mobility are the insoluble high field strength elements (Braun et al., 1993). In addition, some elements (such as Fe and REE) are preferentially complexed by organic matter (e.g., Zhaoliang et al., 2006; Pourret, 2006), and/or are taken up by specific bacteria (e.g., Takahashi et al., 2005). The dissolved REE can easily be assimilated onto certain bacteria cell walls by adsorption on carboxylate and phosphate sites (e.g., Takahashi et al., 2005). As a result of interaction between host rock and organic matter or bacteria, trace elements incompatible with the calcite structure may thus be progressively removed from the host rock to form organic complexes. Taken together, the combined processes of calcite dissolution–(re)precipitation and leaching eventually result in the crystallization of a purer calcite within the host rock, and in the removal of a range of impurity elements that become free to circulate in the medium around.

### 5.2. Implications on the Magnola scarp rock chemical contents

The weathering processes described above account well for the upscarp concentration changes that we document.

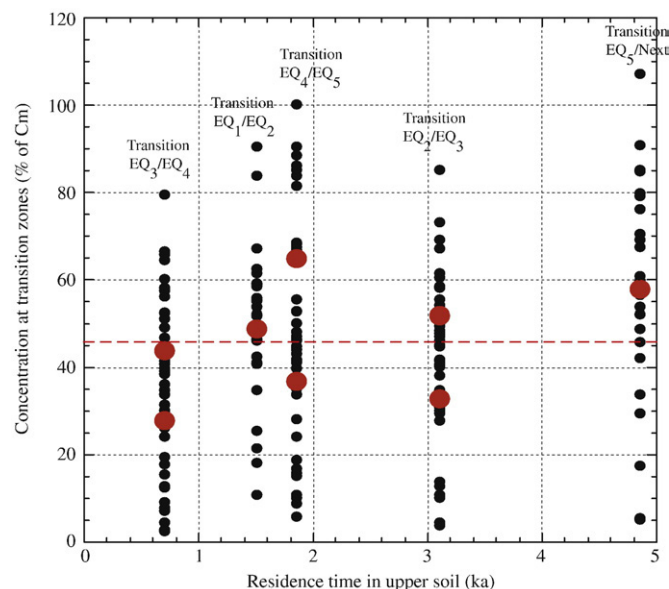
First, it has to be noted that the scarp rocks that we analyze have formed at depth, likely in the first kilometer below the ground surface (the total slip accommodated by the fault is less than 1 km). The exposed slip surface that we study is actually a 1–15 cm-thick layer of tough cemented rocks. Such architecture is common to most active normal faults offsetting carbonates (e.g., Agosta and Aydin, 2006). The cemented slip zone is generally thought to have formed as a result of mixed processes such as comminution of the original rock due to intense shearing during repeated earthquake slip events, dissolution/precipitation of the host rock minerals, various chemical reactions between rock minerals and external agents, etc... (e.g., Sibson, 1977; Chester and Logan, 1986). The high content in calcite crystals and the low Mg concentration of the Magnola fault carbonates suggest that the scarp rocks have initially formed at depth through repeated events of dissolution–precipitation (e.g., Busemberg and Plummer, 1989). We thus expect the carbonate scarp rocks to have had a homogeneous composition before they were exposed to the free air. In any case the trace elements pattern of the initial host limestone is preserved and uniform in the cemented slip zone at the scale of the scarp (Fig. B in the Appendix). More than 70% of the analyzed elements have their

concentration decreasing upward along the scarp height and therefore as a function of exposure time. The rate of decrease is high, averaging 5% per meter, and leads to concentration reductions up to 100% in the topmost samples that have been exposed for the longest time. This shows that the limestone scarp is increasingly degraded as exposure time increases, with that degradation consisting in the progressive leaching of most chemical elements. It is likely that such leaching occurs concomitant with a steady-state dissolution–(re)precipitation of the calcite. Such steady-state dissolution–(re)precipitation process is consistent with the observed upscarp stability of the elements that easily substitute to Ca in the calcite lattice (Pb, Sr, Mg and P). The REE–Y elements are particularly affected by the upscarp leaching. This may be because these trace elements occupy weak positions (such as grain boundaries) within the carbonate assemblages, and as such, can be easily mobilized, even in the absence of any lattice destruction. The LREE are leached more efficiently (–8.6%/m) than the HREE (–6.9%/m), possibly because endolithic bacteria that are known to live in carbonate host rocks retain LREE less efficiently than HREE (Takahashi et al., 2005).

In contrast to most trace elements, Mn, U and possibly Na concentrations increase as a function of exposure time. It could be related to the capacity of Mn, U and Na to enter the calcite structure. Yet, as we highlighted the presence of organic matter on the core surfaces, it is more likely that the upscarp increase in uranium concentration results from interaction between the rock and organic material present on or next to the surface.

Superimposed to the overall concentration changes discussed above, an additional signal – concentration peaks at the transition zones separating the major slip events – is identified in most concentration–scarp height profiles (Fig. 5 and 6). The concentrations of most elements increase by up to 100% (compared to the average value Cm) within a 0.5–1 m thick zone that brackets the earthquake transitions. As shown in Fig. 1, the transitions separating the different exhumed fault sections are specific because they remain at the interface between soil and atmosphere for the long periods of time (commonly several thousands of years) that separate the major earthquakes. As seen before, the soil–atmosphere interface, and more specifically the first 30–50 cm of the soil, is a zone of intense acidification, extremely harmful to limestone. Not only the large amounts of acidic fluids that reside through the soil enhance calcite dissolution, but also the high enrichment of those fluids in numerous chemical impurities contributes to contaminate the re-precipitating calcite by introducing various trace elements in its structure and/or at its grain boundaries. We thus suggest that the concentration peaks observed at the transition zones between earthquakes result from the scarp rock calcite having re-precipitated in an impurity-rich upper soil. The impurities that accumulated in the soil and in the fluids, colloids and particles, likely have their origin in the weathering of the footwall section of the scarp already exhumed above the soil. The REE–Y being the elements most affected by leaching of the exposed scarp, they are expected to be highly concentrated in the upper soil and residing fluids (e.g., Duddy, 1980; Dia et al., 2000; Tauton et al., 2000; Tyler, 2004; Zhaoliang et al., 2006; Pourret et al., 2007). This is consistent with those elements having the largest concentration peaks at the transition zones. Because a colluvial layer (that can be locally several 10 cm thick; e.g., Palumbo et al., 2004) progressively builds up at the base of the scarp as the scarp degrades, the thickness of the acidic layer can locally increase and extend the zone of concentration peaks slightly ‘above’ the actual transition zone (Fig. 1, right panel).

The scenario that we suggest implies that the concentration peaks measured at the transition zones built up during the quiescence periods that separated the major earthquakes. According to Palumbo et al. (2004), those periods, which we call ‘residence times in the upper soil’, range from 0.7 to 4.85 ka (Fig. 5). Fig. 8 shows the concentration peaks (in %) at the transition zones, as a function of residence time in the upper soil (data in Table E in the Appendix). The black dots



**Fig. 8.** Concentration peaks in transition zones as a function of residence time in upper soil. Concentrations (in % of Cm) are measured relative to the trough of the wave identified in each seismic zone (see Fig. 6b). Data from Table E (Appendix). Individual elements are shown with black dots while red symbols correspond to the average concentrations of all elements. Transition zones between earthquakes are indicated. Note that concentration peaks are shown by one or two values for they are defined by one or two values bracketing the transition zones (see Figs. 5 and 6). The horizontal dotted-line indicates the average, steady-state concentration at the transition zones.

correspond to individual elements while the red symbols show their average behavior. The concentration peaks differ depending on the considered element. Yet the average values are roughly constant, about 46% (dotted-line), regardless of the residence time. This constancy suggests that the concentration peaks build up rapidly in the soil (concentrations reach the average peak value in less than 0.7 ka), where they are subsequently preserved (concentration peaks remain constant regardless of residence time). The elements with concentration peaks higher than the average value are primarily the REE and Y (Table E in the Appendix). We could not really evaluate whether the coatings follow a similar pattern because the layer is too thin to be analyzed by the same technique as the host rocks. However, our qualitative major element study shows that the coatings have chemical compositions different from that of the host rock and partly derive from external contributions. We suggest that the chemical composition of the coating essentially reflects the local, intense chemical exchanges that occur between the rock surface and the biologic compounds that reside on that surface. Micro-scale disruptions of the rock surface may also favor trapping of exogenous materials such as the observed Al, Si, Fe, and S-bearing particles (Potter and Rossmann, 1977; Dorn, 1998, 2007). The coating chemical content is thus constantly remobilized and modified by external agents, so that it probably does not retain the signal observed in the ‘preserved’ host rock.

### 5.3. Inferring earthquake slip history from the chemical content of fault scarp rocks

Our study shows that:

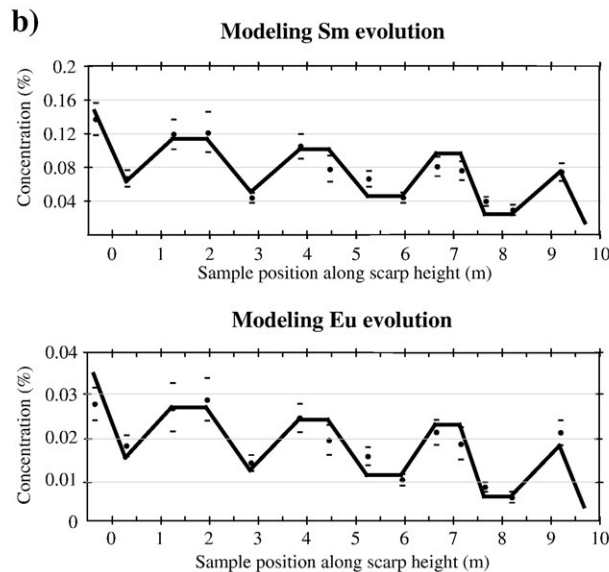
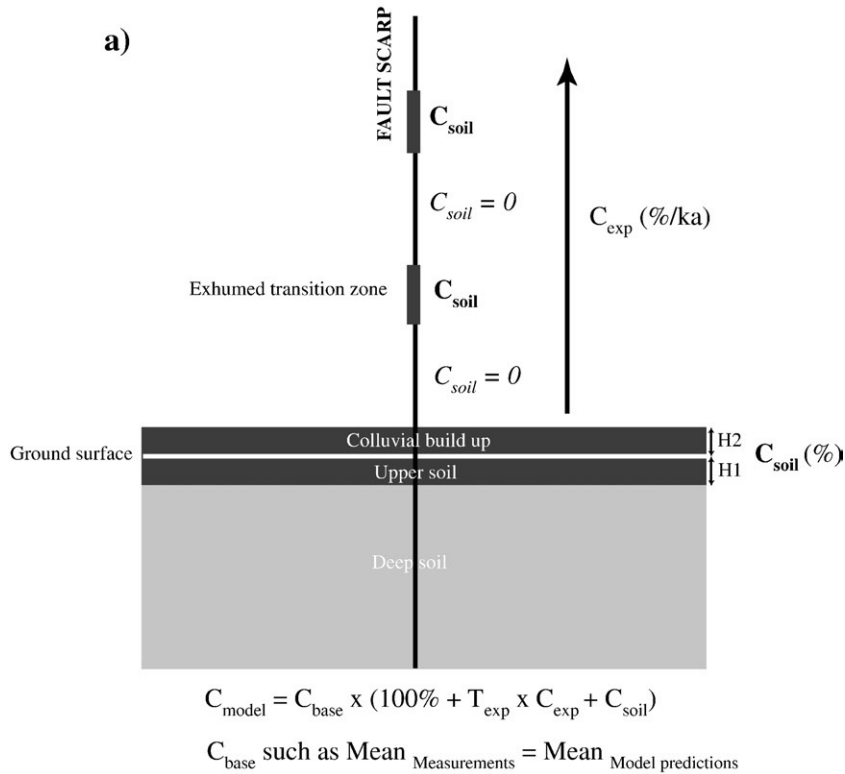
- The concentrations of most major and trace elements decrease as a function of exposure time of the fault scarp surface, at a rate averaging 5%/m and 4.5%/ka. The REE and Y have the highest decrease rate ( $\approx 8\%/m$  and  $7.2\%/ka$ ; Table C-b in the Appendix).
- Concentration peaks are found at the transition zones between seismically exhumed scarp sections. They most likely result from enrichment processes occurring while the rocks are long standing in the upper soil.

- Those enrichments in the soil occur at a fast rate since steady-state concentration peaks are reached in less than 700 years.

We suggest that these observations can be generalized to unravel the recent (<10–15 ka) earthquake slip history of a normal fault having a well-preserved limestone scarp at the ground surface. The approach would consist in:

- Sampling the rocks from the base to the top of the preserved fault scarp; a sampling performed every 30–50 cm should be sufficient to capture the major earthquake ruptures; samples need not be thicker than 2 cm;

- Measure the host rock Y and REE contents, which are the best markers, and report them as a function of scarp height;  
 - Concentration peaks should be found that locate the transition zones between the past major earthquake ruptures; the concentration-height plots may thus allow identifying the major earthquake rupture zones on the scarp, while measuring the amplitude of the slip produced at surface by those earthquakes (at the site of study);  
 - The age of the identified earthquakes might then be obtained by exposure dating of a few samples only taken from each identified rupture zone.



**Fig. 9.** First-order model proposed to reproduce the observations. a) Basic parameters of the model. The fault scarp is sketched in cross-section. The evolution of concentrations as earthquake slip events repeat is sketched in Fig. 1 (right panel). In the absence of data to document it, the transition from  $C_{soil}$  to C in deeper soil is assumed to be sharp. b) Model predictions (broken line) compared to measurements (dots) for Sm and Eu. The fit coefficients  $R$  are reported in Table F (Appendix).



It is noteworthy that faults with a fresh scarp that exposes the preserved fault plane, generally are active normal faults that offset limestone rocks in a temperate environment; other rock types and/or conditions do not allow fault planes to be preserved very long at the surface. This suggests that the approach that we have developed may be applicable to any fault exhibiting a preserved scarp surface.

#### 5.4. A possible model to recover the seismic slip histories

To strengthen and generalize our approach, we built a first-order numerical model presented in Fig. 9. The model input parameters are (Fig. 9a): 1) the existence of a 0.5–1 m thick zone of enrichment at the soil–atmosphere interface (half extending in the upper soil [H1], half in colluvial buildup at the scarp base [H2]). Note that we consider a rather large transition zone to allow realistic uncertainties on the transition positions; 2) a constant concentration peak  $C_{\text{soil}}$  (%) in the transition zones; 3) a constant concentration decrease rate  $C_{\text{exp}}$  (%/ka) in the exposed scarp surface. Note that we ignore how the concentrations evolve from the upper to the deeper soil; in the absence of data, we thus simply hypothesize a sharp transition between  $C_{\text{soil}}$  and  $C$  in the deeper soil (Fig. 1, right panel). The model calculates the concentration  $C_{\text{model}}$  of an element as a function of scarp height, with the formula  $C_{\text{model}} = C_{\text{base}} \times (100\% + T_{\text{exp}} \times C_{\text{exp}} + C_{\text{soil}})$ , with  $C_{\text{base}}$  a reference chosen such that the average measurement concentration equals that of the model values, while  $T_{\text{exp}}$  is the exposure time of the sample. The earthquake slip history is that inferred by Palumbo et al. (2004), as discussed before. Fig. 9b shows the calculated values for two elements, Sm and Eu, while Table F (Appendix) reports the fit coefficient  $R$  between measurements and model predictions for all REE-Y elements. We used a value of  $C_{\text{soil}}$  of 53%, which corresponds to the average concentration peak for the REE-Y group (Table E in the Appendix), while  $C_{\text{exp}}$  is  $-7.2\%/ka$ , the average decrease rate for the REE-Y group (Table C-b in the Appendix). The model shown by the thick line in Fig. 9b fits well the measured concentrations shown by the dots and their error bars, with  $R$  coefficients averaging 91%, thus much better than a simple linear regression (Table F in the Appendix, right column). This suggests that our simple model reproduces quite well the observations. Provided that the sampling is dense enough along the scarp to precisely locate the transition zones and that a few exposure ages are available, the model might be inverted to recover part of an ignored earthquake slip history.

## 6. Conclusions

So far, identification of past major earthquakes on faults has been a difficult and costly task, reachable through only two techniques, paleoseismology and cosmogenic exposure dating. Though powerful, those methods are costly. In the present study, we have developed a new method that may help deciphering the recent (last 10–15 ka) earthquake history of a normal fault having a well-preserved limestone scarp at the ground surface, in a simpler and less expensive way (an order of magnitude cheaper) than presently possible.

We showed that the chemical content of young limestone fault scarp rocks contains quantitative information on the scarp exhumation history. The element concentrations vary as the result of two parameters: (1) the time during which rocks are exposed to the air, and (2) the time during which rocks are in contact with the upper soil (30–50 first cm). The former correlates with the age of the earthquakes, while the later correlates with their recurrence time. The exposure time governs the duration, hence the intensity, of the various weathering processes that affect the exposed carbonate surface. In our case, steady-state microscopic calcite dissolution–precipitation seems to be the dominant process, resulting in progressive leaching of most elements, at a roughly constant rate of 5%/m or 4.5%/ka. On the other hand, the residence time of the scarp

rocks at the soil–atmosphere interface governs the duration, hence the intensity, of the weathering processes that occur in the upper soil. Calcite dissolution–precipitation seems again the dominant process. Yet only the narrow scarp section in contact with the impurity-doped upper soil is affected by that process, while other parts of the escarpment are not. Narrow concentration peaks thus result that precisely locate the paleo-positions of the upper soil, hence of the transition zones between the successive major earthquakes.

We conclude that simple geochemical analyses may help recovering the seismic slip history of a fault having a preserved escarpment. The only requirement is that sampling is dense enough from base to top of the scarp for concentration–scarp height curves to be determined with sufficient detail. REE and Y seem to be the best tracers. Local concentration peaks along the scarp height should correspond to the location of transition zones between major earthquake ruptures. Knowing the position of the transition zones should help determining the amplitude of the slips produced by the successive earthquakes. We propose a first-order model that reproduces well the observations. Provided that a fault scarp is sampled densely enough for the concentration–scarp height curves to be precisely determined, the model might be inverted to recover, through iterative adjustments, the best-fitting earthquake slip exhumation history on the fault. This would significantly limit the number of required dating of the fault scarp.

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## Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at [doi:10.1016/j.epsl.2008.03.059](https://doi.org/10.1016/j.epsl.2008.03.059).

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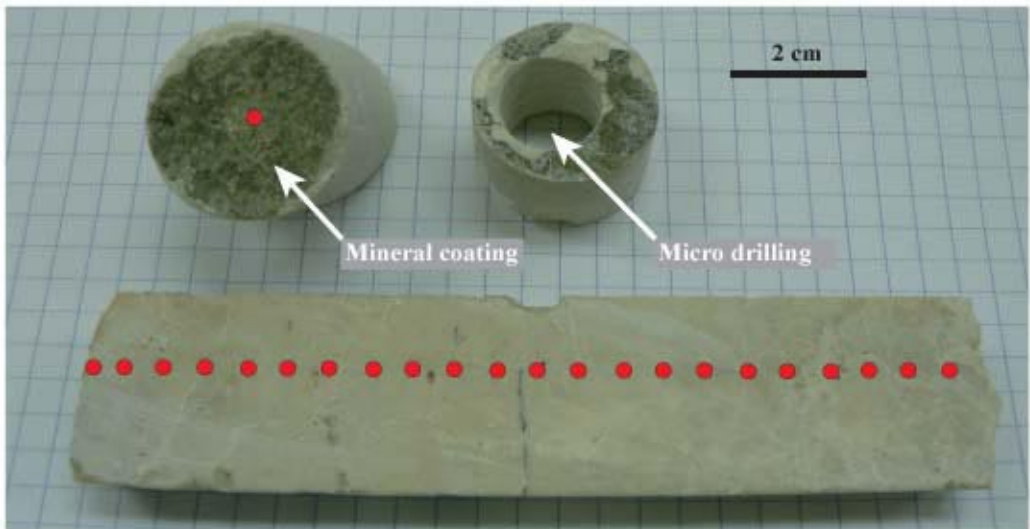
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a)



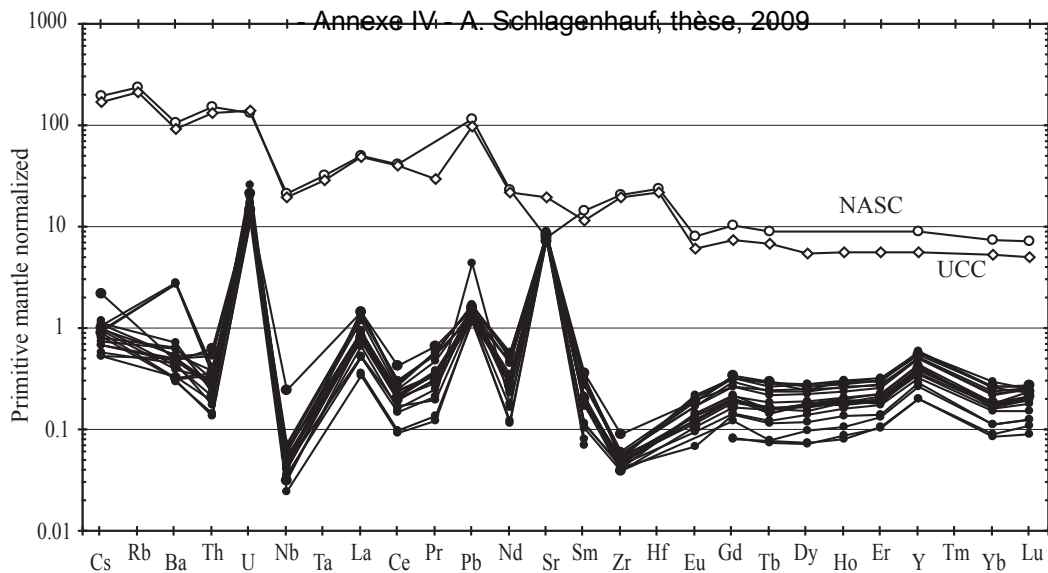
b)



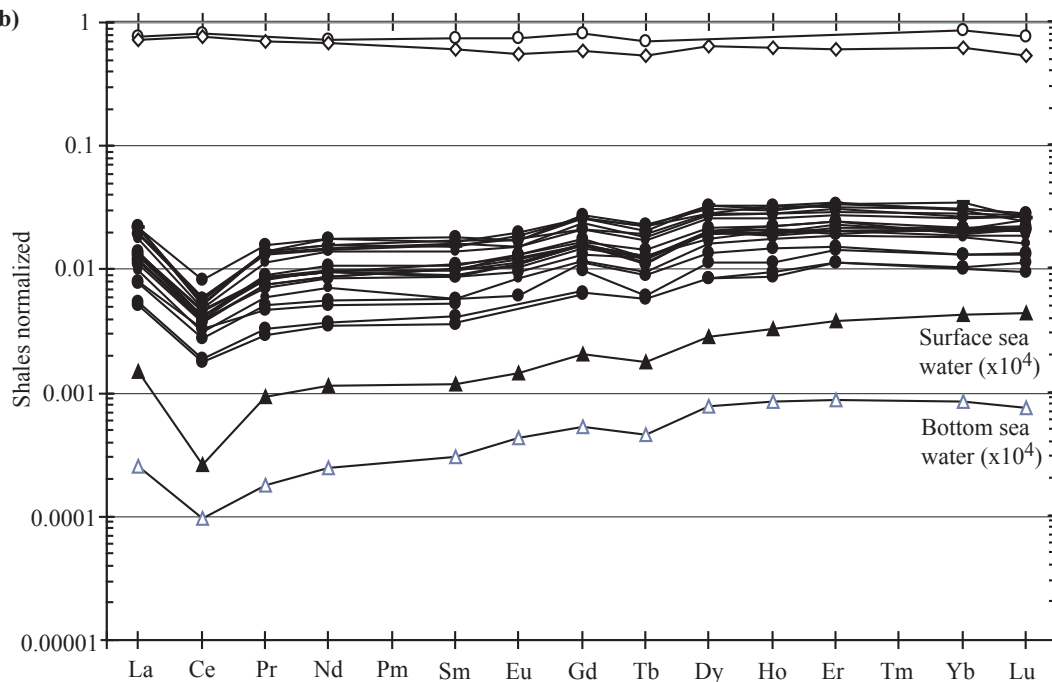
a) Westward view in cross-section of the Magnola fault scarp, showing the 1-15 cm-thick carbonate-made slip zone, overlaying carbonate breaches.

b) View of one sampled core. The top left piece shows the surface coating with a dark color while the half-core at the bottom shows the light color of the inner rock (host rock). The red dots locate the semi-quantitative measurements done with  $\mu$ XRF. Major and trace element concentrations were measured in micro-samples drilled within the host rock, as shown in the top-right figure.

a)



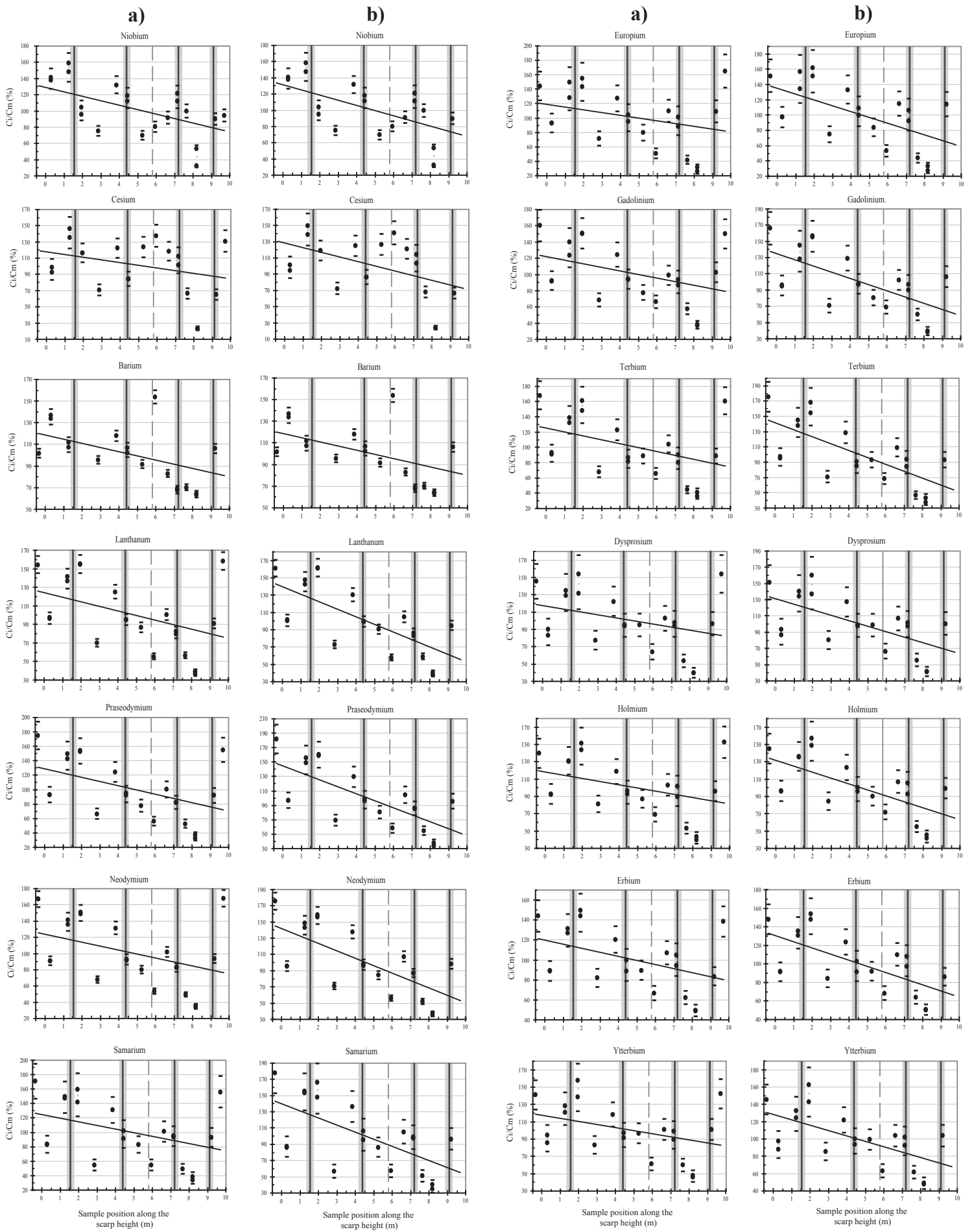
b)



Trace element and Rare Earth patterns for the Magnola samples.

a) The concentrations are normalized to primitive mantle (Hofmann, 1988). The black dots show data from this study while the empty dots show the pattern of the North American Shale Composite (Gromet et al., 1984; Condie, 1993) and the empty diamonds, the pattern of average Upper Continental Crust (Mc Lennan, 2001). The Magnola host rocks exhibit the same general pattern as the two references, but the absolute concentrations are lower by two orders of magnitude. This suggests that the Magnola carbonates have a chemical signature close to that of other sediments, but with an overall depletion that can be related to the high carbonate content (e. g., Plank and Langmuir, 1998). Significant positive anomalies in Ba, U, Sr, and Y, and a negative anomaly in Ce are also highlighted.

b) The REE concentrations are normalized to Reference Shales (Sholkovitz, 1988). Symbols as in a, with solid and empty triangles for surface and bottom seawater, respectively ( $\times 10^4$ ; Alibo and Nozaki, 1999). The Magnola REE patterns closely resemble that of seawater, though concentrations are five to six orders of magnitude higher. This shows that the seawater origin of the Magnola carbonates is preserved in the analyzed samples.

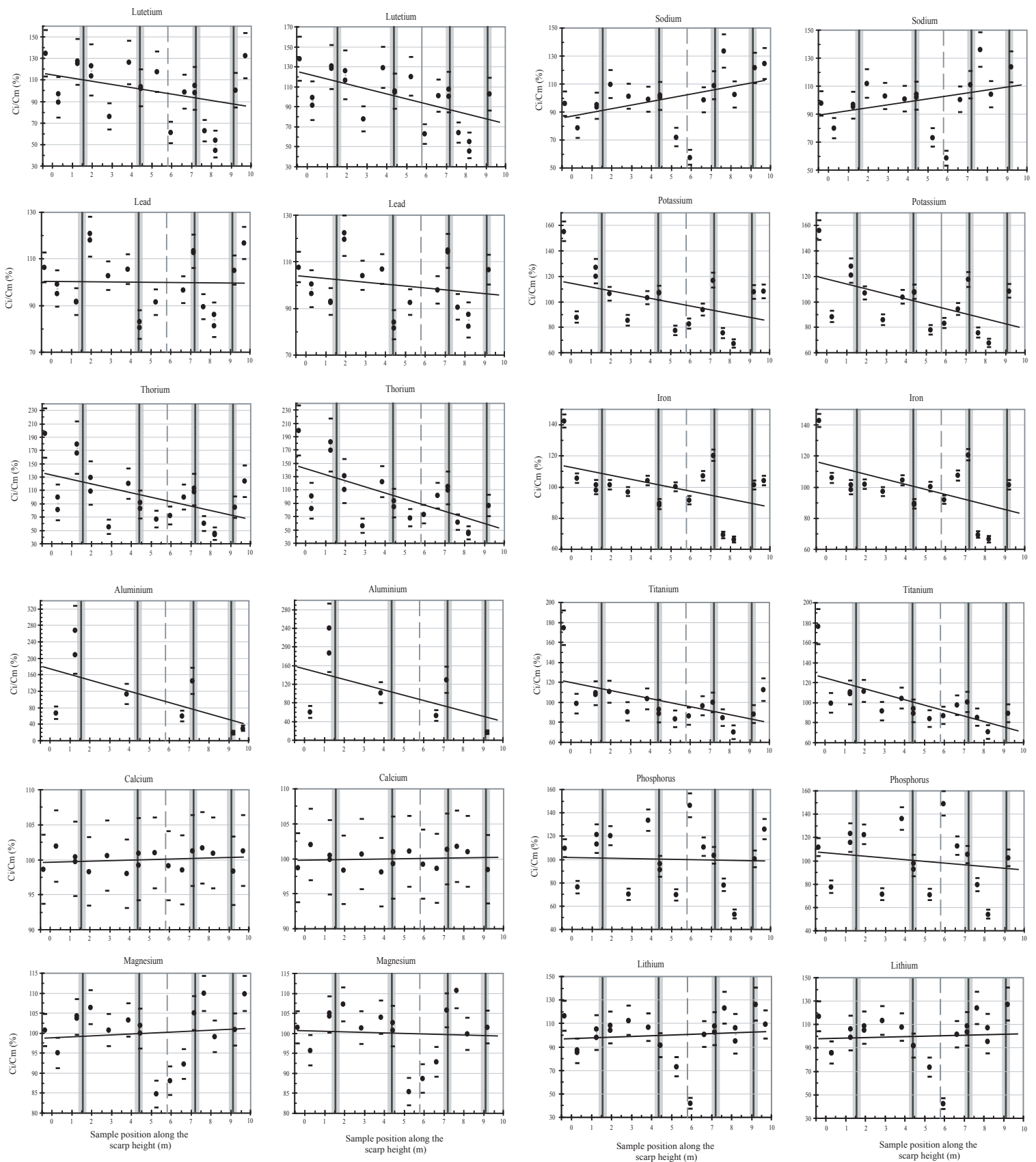


Carcaillet et al., Figure C - Electronic Supplement (caption in following)

a)

- Annex IV - A. Schlagenhauf, thèse, 2009

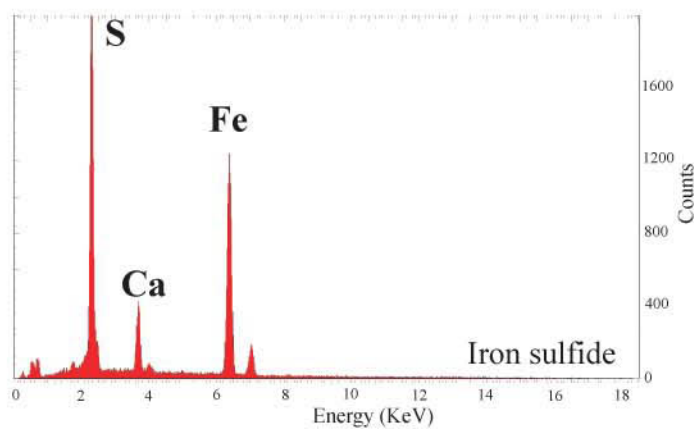
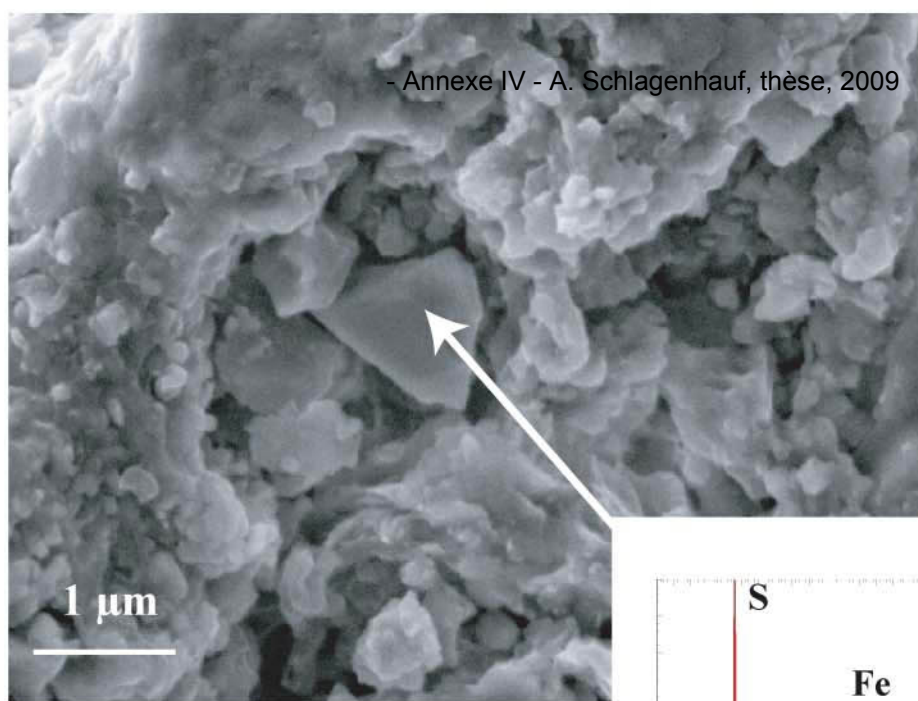
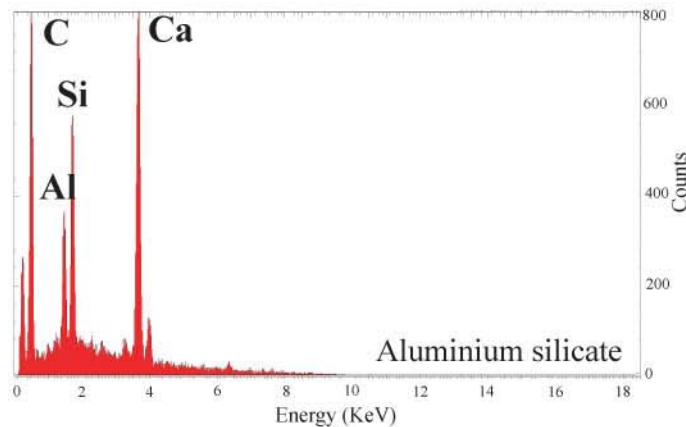
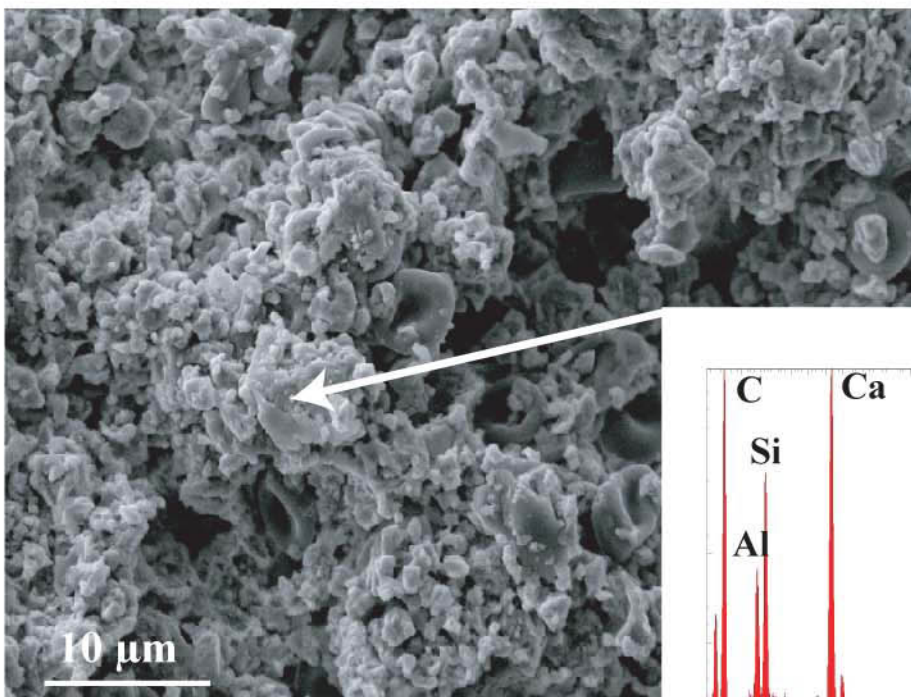
IV/187



Element concentrations vs. sampling height for all elements not presented in Figure 3. In all plots, X is the on-fault position of the samples, in meters, increasing upward; Y is the concentration of sample  $i$  ( $C_i$ ) normalized by the mean concentration ( $C_m$ ) averaged over the whole sample collection, in percentage. Solid vertical lines represent the transitions between the major earthquake slip zones identified by Palumbo et al. (2004), with their position uncertainties in grey. The thin dotted vertical line locates a possible additional earthquake suggested by Palumbo et al. (2004). The slanted straight lines are the regression curves calculated from all measurements. When several measurements were done for the same sample, only the average was used for the regression calculation, but all measurements are plotted in the figure.

a) Regressions are calculated using all measurements listed in Table 1. Regression parameters are reported in Table C-a.

b) Regressions are calculated as in a) but without sample C15. Regression parameters are reported in Table C-b.

**a)****b)**

SEM images and EDS analyzes of minerals embedded in the coating. Photographs on the left are SEM images of the coating in core C9, while plots show EDS analyzes of the two minerals shown by the arrows.

Annexe IV - A. Schlagenhauf these, 2009

Samples	Elevation (m)	Depth in core (cm)	Na <sub>2</sub> O (%)	MgO (%)	Al <sub>2</sub> O <sub>3</sub> (%)	SiO <sub>2</sub> (%)	P <sub>2</sub> O <sub>5</sub> (%)	SO <sub>3</sub> (%)	K <sub>2</sub> O (%)	CaO (%)	TiO <sub>2</sub> (%)	MnO (%)	Fe <sub>2</sub> O <sub>3</sub> (%)
C1	-0.35	<i>Coating</i>	<i>1.79</i>	<i>0.49</i>	<i>1.18</i>	<i>2.77</i>	<i>1.83</i>	<i>0.46</i>	<i>0.52</i>	<i>90.38</i>	<i>0.12</i>	<i>0.03</i>	<i>0.47</i>
		0.1	2.39	0.82	0.18	0.44	2.06	0.3	0.24	93.42	0.01	0.01	0.14
		0.5	2.37	0.72	0.03	0.08	2.02	0.25	0.24	94.1	0.03	0.02	0.14
		1	2.81	0.85	0.25	0.27	2.21	0.38	0.26	92.8	0.05	0.01	0.11
		1.5	3.1	0.9	0.2	0.18	2.15	0.36	0.25	92.69	0.05	0.02	0.09
		2	2.43	0.67	0	0.07	2.01	0.25	0.27	94.15	0.02	0.01	0.12
		2.5	2.46	1.02	0.27	0.28	2.21	0.4	0.25	92.74	0.24	0.01	0.11
		3	2.67	0.83	0.36	0.3	2.17	0.36	0.28	92.84	0.04	0.03	0.12
		3.5	2.39	0.82	0.18	0.44	2.06	0.3	0.24	93.42	0.01	0.01	0.14
		4	2.59	0.7	0.18	0.26	2.28	0.38	0.29	93.13	0.07	0.01	0.11
		4.5	2.83	0.72	0.15	0.27	2.14	0.44	0.33	92.96	0.06	0.01	0.1
		5	0.59	0.84	0.19	0.26	2.22	0.38	0.28	94.82	0.21	0.01	0.2
		5.5	2.93	0.95	0.25	0.24	2.19	0.38	0.28	92.61	0.04	0.01	0.11
		6	2.34	0.95	0.19	0.29	2.07	0.29	0.28	93.4	0.01	0.02	0.17
		6.5	2.34	0.95	0.19	0.29	2.07	0.29	0.28	93.4	0.01	0.02	0.17
		7	2.6	0.87	0.3	0.3	2.19	0.41	0.26	92.86	0.03	0.02	0.15
		7.5	2.5	0.8	0.16	0.22	2.16	0.37	0.27	93.35	0.03	0.02	0.12
		8	2.37	0.63	0	0.07	2	0.25	0.25	94.21	0.03	0.01	0.18
8.5	2.69	0.86	0.25	0.31	2.16	0.42	0.28	92.87	0.07	0.01	0.09		
9	2	0.63	0.1	0.23	2.09	0.28	0.27	94.22	0.03	0.01	0.13		
C2	0.30	<i>Coating</i>	<i>1.75</i>	<i>0.49</i>	<i>1.18</i>	<i>2.77</i>	<i>1.83</i>	<i>0.46</i>	<i>0.52</i>	<i>90.38</i>	<i>0.12</i>	<i>0.03</i>	<i>0.47</i>
		0.1	1.86	0.52	0.11	0.25	1.95	0.42	0.32	94.32	0.06	0.02	0.17
		0.5	2.81	0.74	0.16	0.1	1.9	0.29	0.3	93.49	0.05	0.03	0.13
		1	2.49	0.73	0.2	0.22	1.97	0.36	0.3	93.37	0.1	0.03	0.23
		1.5	2.22	0.75	0.2	0.22	1.96	0.4	0.39	93.56	0.1	0.07	0.15
		2	0	0.47	0.1	0.22	1.96	0.34	0.21	96.54	0.04	0.01	0.09
		2.5	2.89	0.63	0.17	0.26	1.94	0.37	0.28	93.2	0.06	0.03	0.17
		3	1.65	0.55	0.05	0.07	1.9	0.32	0.3	94.94	0.03	0.03	0.15
		3.5	1.32	0.66	0.18	0.19	1.97	0.37	0.28	94.73	0.03	0.02	0.25
		4	2.06	0.62	0.16	0.2	1.96	0.35	0.29	94.1	0.07	0.03	0.15
4.5	1.74	0.64	0.16	0.31	2.3	0.38	0.31	93.86	0.06	0.04	0.2		
C3	1.25	<i>Coating</i>	<i>1.7</i>	<i>0.49</i>	<i>1.48</i>	<i>4.2</i>	<i>1.81</i>	<i>0.45</i>	<i>0.52</i>	<i>88.35</i>	<i>0.2</i>	<i>0.03</i>	<i>0.8</i>
C4	1.95	<i>Coating</i>	<i>1.04</i>	<i>0.28</i>	<i>2.39</i>	<i>7.2</i>	<i>1.9</i>	<i>0.6</i>	<i>0.64</i>	<i>84.79</i>	<i>0.25</i>	<i>0.04</i>	<i>0.88</i>
		0.1	2.24	0.73	0.15	0.09	1.91	0.44	0.3	93.9	0.06	0.03	0.16
		0.5	2.2	0.81	0.21	0.09	1.9	0.33	0.29	93.92	0.1	0.02	0.15
		1	1.97	0.57	0.05	0.08	1.92	0.34	0.27	94.56	0.07	0.03	0.13
		1.5	1.85	0.59	0.17	0.24	1.98	0.41	0.31	94.13	0.12	0.02	0.15
		2	1.71	0.69	0.17	0.18	1.97	0.36	0.34	94.28	0.07	0.02	0.21
		2.5	2.14	0.48	0.06	0.08	1.93	0.3	0.27	94.43	0.08	0.07	0.17
		3	2.24	0.69	0.07	0.09	1.91	0.35	0.3	94.08	0.07	0.04	0.16
		3.5	1.9	0.6	0.15	0.24	2.01	0.36	0.35	94.13	0.06	0.04	0.16
		4	1.98	0.62	0.12	0.2	1.96	0.35	0.37	94.14	0.09	0.03	0.14
4.5	1.53	0.61	0.13	0.24	1.95	0.39	0.3	94.64	0.02	0.18	0.08		
5	1.99	0.66	0.13	0.31	1.83	0.33	0.28	94.14	0.06	0.04	0.23		
C5	2.85	<i>Coating</i>	<i>1.09</i>	<i>0.04</i>	<i>1.56</i>	<i>3.94</i>	<i>1.85</i>	<i>0.55</i>	<i>0.63</i>	<i>89.62</i>	<i>0.14</i>	<i>0.04</i>	<i>0.56</i>
C6	3.85	<i>Coating</i>	<i>0.94</i>	<i>0.5</i>	<i>3.75</i>	<i>7.53</i>	<i>2.05</i>	<i>0.88</i>	<i>0.71</i>	<i>82.74</i>	<i>0.17</i>	<i>0.03</i>	<i>0.7</i>
C7	4.45	<i>Coating</i>	<i>1.45</i>	<i>0.34</i>	<i>1.45</i>	<i>4.39</i>	<i>1.79</i>	<i>0.43</i>	<i>0.49</i>	<i>88.91</i>	<i>0.13</i>	<i>0.03</i>	<i>0.6</i>
C8	5.25	<i>Coating</i>	<i>1.02</i>	<i>0.08</i>	<i>1.86</i>	<i>4.44</i>	<i>1.86</i>	<i>0.52</i>	<i>0.53</i>	<i>88.81</i>	<i>0.16</i>	<i>0.04</i>	<i>0.66</i>
		0.1	1.81	0.44	0	0.03	1.86	0.27	0.28	95.07	0.05	0.04	0.16
		0.5	1.45	0.78	0.08	0.29	1.93	0.27	0.34	94.52	0.1	0.02	0.21
		1	1.77	0.56	0.09	0.08	1.98	0.34	0.28	94.68	0.06	0.06	0.1
		1.5	1.74	0.42	0.13	0.21	2.03	0.38	0.3	94.54	0.09	0.03	0.13
		2	1.92	0.53	0.14	0.19	2.02	0.38	0.36	94.07	0.18	0.03	0.2
		2.5	1.96	0.58	0.06	0.08	1.97	0.37	0.26	94.51	0.02	0.02	0.18
		3	1.38	0.58	0.11	0.27	1.91	0.28	0.22	94.97	0.02	0.03	0.23
		3.5	2.03	0.61	0.2	0.5	1.89	0.28	0.32	93.86	0.07	0.03	0.2
		4	1.75	0.5	0.16	0.23	2	0.41	0.34	94.33	0.11	0.03	0.14
		4.5	2.18	0.72	0.25	0.24	1.98	0.28	0.27	93.77	0.06	0.04	0.21
		5	2.28	0.73	0.22	0.26	2.02	0.37	0.31	93.49	0.09	0.03	0.19
		5.5	2.21	0.64	0.25	0.3	2	0.27	0.28	93.76	0.08	0.02	0.18
6	2.07	0.56	0	0.05	1.92	0.28	0.26	94.62	0.07	0.02	0.16		
6.5	1.94	0.65	0.18	0.26	2.04	0.4	0.29	93.97	0.1	0.02	0.16		
7	1.97	0.79	0.16	0.07	1.87	0.3	0.25	94.3	0.08	0.04	0.18		
7.5	2.09	0.71	0.14	0.19	2.01	0.37	0.27	93.97	0.06	0.03	0.15		
8	1.98	0.75	0.17	0.19	2.05	0.4	0.3	93.89	0.07	0.04	0.18		
C9	5.95	<i>Coating</i>	<i>1.7</i>	<i>0.34</i>	<i>0.96</i>	<i>2.28</i>	<i>1.94</i>	<i>0.48</i>	<i>0.42</i>	<i>91.38</i>	<i>0.11</i>	<i>0.03</i>	<i>0.34</i>
C10	6.65	<i>Coating</i>	<i>1.95</i>	<i>0.38</i>	<i>0.67</i>	<i>1.81</i>	<i>1.93</i>	<i>0.4</i>	<i>0.41</i>	<i>91.82</i>	<i>0.11</i>	<i>0.03</i>	<i>0.51</i>
C11	7.15	<i>Coating</i>	<i>1.75</i>	<i>0.26</i>	<i>1.03</i>	<i>2.83</i>	<i>1.86</i>	<i>0.51</i>	<i>0.47</i>	<i>90.71</i>	<i>0.1</i>	<i>0.03</i>	<i>0.46</i>
C12	7.65	<i>Coating</i>	<i>1.95</i>	<i>0.43</i>	<i>0.26</i>	<i>0.49</i>	<i>1.85</i>	<i>0.37</i>	<i>0.34</i>	<i>93.99</i>	<i>0.06</i>	<i>0.03</i>	<i>0.24</i>
C13	8.20	<i>Coating</i>	<i>2.09</i>	<i>0.59</i>	<i>1.26</i>	<i>3.29</i>	<i>1.86</i>	<i>0.62</i>	<i>0.63</i>	<i>88.97</i>	<i>0.14</i>	<i>0.04</i>	<i>0.52</i>
		0.1	3.31	0.94	0.26	0.24	2.17	0.42	0.34	92.15	0.04	0.02	0.12
		0.5	2.88	0.72	0.15	0.23	2.17	0.43	0.32	92.91	0.05	0.02	0.13
		1	3.26	0.92	0.23	0.2	2.13	0.39	0.32	92.39	0.04	0.01	0.12
		1.5	3.07	0.68	0.25	0.28	2.21	0.42	0.38	92.51	0.1	0.02	0.1
		2	3.14	0.87	0.2	0.24	2.15	0.43	0.31	92.48	0.06	0.02	0.11
		2.5	2.92	0.73	0.22	0.27	2.17	0.4	0.31	92.81	0.04	0.02	0.11
		3	2.75	0.63	0.2	0.22	2.16	0.41	0.32	93	0.03	0.01	0.26
		3.5	2.89	0.67	0.21	0.26	2.15	0.42	0.34	92.87	0.05	0.02	0.11
		4	2.7	0.6	0.15	0.25	2.15	0.41	0.33	93.24	0.05	0.01	0.12
		4.5	3.53	0.99	0.25	0.28	2.14	0.42	0.34	91.88	0.06	0.01	0.11
		5	3.47	0.91	0.28	0.29	2.08	0.32	0.32	92.14	0.03	0.03	0.13
		5.5	3.27	0.96	0.3	0.32	2.13	0.46	0.32	92.07	0.05	0.02	0.1
		6	3.07	0.88	0.17	0.23	2.14	0.4	0.32	92.6	0.05	0.01	0.12
6.5	2.51	0.55	0.13	0.21	2.16	0.41	0.35	93.5	0.07	0.01	0.11		
7	3.28	0.95	0.25	0.24	2.16	0.4	0.32	92.23	0.04	0.01	0.12		
7.5	3.36	0.88	0.16	0.2	2.15	0.41	0.33	92.34	0.05	0.02	0.11		
C14	9.20	<i>Coating</i>	<i>1.22</i>	<i>0.27</i>	<i>1.58</i>	<i>4.58</i>	<i>1.89</i>	<i>0.51</i>	<i>0.57</i>	<i>88.46</i>	<i>0.19</i>	<i>0.04</i>	<i>0.69</i>
		0.1	1.74	0.71	0.12	0.33	2.11	0.32	0.34	94.08	0.06	0.03	0.17
		0.5	2.1	0.8	0.17	0.05	1.76	0.27	0.32	94.24	0.08	0.02	0.19
		1	2.16	0.74	0.16</								



a)

	Sr	Y	Zr	Nb	Cs	Ba	La	Ce	Pr	Nd	Sm	Eu	Gd	Tb	Dy	Ho	Er	Yb	Lu	Pb	Th	U	Al <sub>2</sub> O <sub>3</sub>	CaO	MgO	Na <sub>2</sub> O	K <sub>2</sub> O	Fe <sub>2</sub> O <sub>3</sub>	MnO	TiO <sub>2</sub>	P <sub>2</sub> O <sub>5</sub>
Sr	0.0	9.6	11.6	0.0	4.7	2.6	14.7	2.8	1.6	1.1	0.1	1.2	1.2	1.2	0.4	0.2	0.3	0.4	1.9	8.6	3.5	23.1	1.2	21.1	0.2	0.5	1.1	0.1	13.8	9.0	7.1
Y	0.0	41.9	23.6	47.3	3.7	95.2	67.4	91.3	94.8	93.9	95.9	93.3	94.3	98.7	97.5	96.3	97.0	88.1	88.1	44.6	64.5	0.6	2.4	17.1	11.2	3.1	53.2	45.2	19.8	47.9	35.3
Zr	9.6	41.9	73.2	36.2	5.0	52.5	83.7	64.1	56.9	60.2	43.3	59.1	58.0	48.5	47.1	54.9	47.1	43.8	40.3	14.2	81.0	1.3	49.7	12.1	5.9	0.2	76.0	63.4	20.4	93.8	14.5
Nb	11.6	23.6	73.2	22.9	12.9	29.5	53.8	33.4	27.8	33.1	29.3	31.2	22.9	25.0	26.0	28.1	22.4	30.3	1.6	57.9	0.5	61.0	0.2	10.5	0.1	48.5	23.4	5.7	49.4	14.0	
Cs	0.0	47.3	36.2	22.9	27.3	39.7	43.3	40.6	39.4	43.1	36.6	40.1	47.3	45.8	44.7	45.6	37.1	43.6	17.8	47.3	0.4	32.8	8.7	27.7	15.3	24.4	40.5	19.1	47.5	56.1	
Ba	4.7	3.7	5.0	12.9	27.3	6.7	78.4	98.0	98.7	96.3	95.0	97.1	97.3	96.9	96.9	95.8	96.7	83.7	36.2	71.2	1.8	3.7	15.8	13.6	3.0	56.6	41.6	23.7	56.9	28.7	
La	2.6	67.4	83.7	53.8	43.3	4.7	78.4	86.2	80.1	82.3	68.9	81.8	80.1	73.9	70.7	75.2	73.6	70.7	18.2	85.9	0.6	22.9	13.9	6.5	0.5	76.9	61.2	17.0	87.4	16.7	
Ce	2.8	91.3	64.1	33.4	40.6	6.1	98.0	86.2	98.5	98.0	92.4	98.3	97.5	94.8	94.5	95.0	93.9	80.8	33.4	81.4	1.1	8.4	19.1	13.6	2.5	68.0	48.3	21.3	68.4	30.1	
Pr	1.6	94.8	56.9	27.8	39.4	5.9	98.7	80.1	98.5	97.4	95.4	98.0	98.0	97.2	96.5	95.0	95.5	83.7	38.1	74.6	0.9	1.7	17.7	14.2	3.5	61.8	46.3	21.0	62.0	31.4	
Nd	1.1	93.9	60.2	33.1	43.1	3.7	96.3	82.3	98.0	97.4	93.2	97.3	96.2	95.9	94.0	95.0	93.7	85.8	33.0	82.7	0.4	12.9	19.6	13.4	2.8	69.8	48.6	18.2	63.6	33.4	
Sm	0.1	95.9	43.3	29.3	36.6	5.7	95.0	68.9	92.4	95.4	93.2	94.8	92.3	97.1	97.5	92.3	95.9	85.3	42.8	66.3	0.3	0.8	17.8	13.7	4.8	59.2	46.0	10.9	47.7	32.2	
Eu	1.2	93.3	59.1	31.2	40.1	8.0	97.1	81.8	98.3	98.0	97.3	94.8	96.9	96.9	96.1	94.3	95.4	80.7	42.1	76.8	0.0	2.3	22.6	14.6	3.5	67.4	49.7	16.7	64.3	36.8	
Gd	1.2	94.3	58.0	22.9	47.3	7.7	97.3	80.1	97.5	98.0	96.2	92.3	96.9	96.9	96.7	96.4	95.8	80.7	40.4	74.8	1.1	2.4	19.9	8.5	0.9	60.0	51.1	23.1	64.1	33.4	
Tb	0.4	98.7	48.5	25.0	45.8	5.2	96.9	73.9	94.8	97.2	95.9	97.1	95.9	96.7	96.7	96.4	98.7	89.2	42.8	70.0	0.0	3.4	17.3	10.9	2.5	60.0	49.6	18.3	54.4	33.0	
Dy	0.2	97.5	47.1	26.0	44.7	8.3	96.9	70.7	94.5	96.5	94.0	97.5	96.1	96.4	98.7	96.5	97.5	82.9	46.2	67.8	0.1	3.0	17.5	12.0	2.3	57.8	47.5	19.1	52.5	35.6	
Ho	0.3	96.3	54.9	28.1	45.6	2.8	95.8	75.2	95.0	95.0	95.3	92.3	94.3	95.8	96.4	96.5	96.1	82.5	42.1	73.6	0.9	13.3	18.3	12.6	2.2	60.5	49.8	26.6	59.4	31.5	
Er	0.3	96.3	54.9	28.1	45.6	2.8	95.8	75.2	95.0	95.0	95.3	92.3	94.3	95.8	96.4	96.5	96.1	82.5	42.1	73.6	0.9	13.3	18.3	12.6	2.2	60.5	49.8	26.6	59.4	31.5	
Yb	0.4	97.0	47.1	22.4	37.1	3.6	96.7	73.6	93.9	95.5	93.7	95.9	95.4	95.2	98.2	97.5	96.1	86.9	47.1	65.4	0.1	1.9	18.8	12.0	0.7	57.3	47.9	17.7	53.4	26.6	
Lu	1.9	88.1	43.8	30.3	43.6	2.5	83.7	70.7	80.8	83.7	85.8	80.7	81.6	88.2	82.9	82.5	86.9	26.2	20.2	63.4	0.1	11.2	11.2	4.1	0.7	55.7	52.0	12.0	45.6	19.6	
Pb	8.6	44.6	14.2	1.6	17.8	1.9	36.2	18.2	33.4	38.1	33.0	42.8	42.1	40.4	42.8	46.2	42.1	47.1	26.2	16.5	6.1	19.8	13.0	16.8	8.0	18.9	33.8	7.1	20.4	41.2	
Th	3.5	64.5	81.0	57.9	47.3	3.4	71.2	85.9	81.4	74.6	82.7	66.3	76.8	74.8	70.0	73.6	65.1	63.4	16.5	0.0	71.0	17.0	9.8	0.2	86.0	57.1	13.7	77.7	29.3		
U	23.1	0.1	1.3	0.5	0.0	4.4	1.8	0.6	1.1	0.9	0.4	0.3	0.0	1.1	0.0	0.1	0.9	0.1	6.1	0.0	0.0	6.5	13.9	1.4	0.9	4.4	9.9	58.8	0.7	5.1	29.3
Al <sub>2</sub> O <sub>3</sub>	1.2	2.4	49.7	61.0	32.8	0.5	3.7	22.9	8.4	1.7	12.9	0.8	2.3	2.4	3.4	3.0	13.3	1.9	11.2	19.8	71.0	6.5	0.1	5.2	14.2	48.7	0.8	11.3	20.2	1.9	
CaO	21.1	17.1	12.1	0.2	8.7	9.9	15.8	13.9	19.1	17.7	19.6	17.8	22.6	19.9	17.3	17.5	18.3	18.8	11.2	13.0	17.0	13.9	0.1	0.6	0.6	16.1	13.2	1.7	12.7	37.7	
MgO	0.2	11.2	5.9	10.5	2.7	15.7	13.6	6.5	13.6	14.2	13.4	13.7	14.6	8.5	10.9	12.0	12.6	12.0	4.1	16.8	9.8	1.4	5.2	0.6	72.1	12.2	0.1	7.3	6.4	1.4	
Na <sub>2</sub> O	0.5	3.1	0.2	11.5	0.3	43.6	3.0	0.5	2.5	3.5	2.8	4.8	3.5	0.9	2.5	2.3	2.2	3.7	0.7	8.0	0.2	0.9	14.2	0.2	0.9	2.3	2.4	0.6	0.0	0.0	0.0
K <sub>2</sub> O	1.1	53.2	76.0	48.5	24.4	1.4	56.6	76.9	68.0	61.8	69.8	59.2	67.4	60.0	60.0	57.8	60.5	57.3	55.7	18.9	86.0	4.4	48.7	16.1	12.2	2.3	65.3	2.4	76.0	22.2	
Fe <sub>2</sub> O <sub>3</sub>	0.1	45.2	63.4	23.4	40.5	4.1	41.6	61.2	48.3	46.3	48.6	46.0	49.7	51.1	49.6	47.5	49.8	47.9	52.0	33.8	57.1	9.9	0.8	13.2	0.1	65.3	2.8	66.6	16.6	6.6	
MnO	13.8	19.8	20.4	5.7	19.1	0.9	23.7	17.0	21.3	21.0	18.2	10.9	16.7	23.1	18.3	19.1	26.6	17.7	12.0	7.1	13.7	58.8	11.3	1.7	7.3	2.8	19.1	3.7	3.7	15.4	
TiO <sub>2</sub>	9.0	47.9	93.8	49.4	47.5	2.3	56.9	87.4	68.4	62.0	63.6	47.7	64.3	64.1	54.4	52.5	59.4	53.4	45.6	20.4	77.7	0.7	20.2	12.7	6.4	2.6	76.7	66.7	19.0	15.4	
P <sub>2</sub> O <sub>5</sub>	7.1	35.3	14.5	14.0	56.1	32.4	28.7	16.7	30.1	31.4	33.4	32.2	36.8	33.4	33.0	35.6	31.5	26.6	19.6	41.2	29.3	5.1	1.9	37.7	1.4	22.2	16.6	3.7	15.1	13.7	

b)

	Sr	Y	Zr	Nb	Cs	Ba	La	Ce	Pr	Nd	Sm	Eu	Gd	Tb	Dy	Ho	Er	Yb	Lu	Pb	Th	U	Al <sub>2</sub> O <sub>3</sub>	CaO	MgO	Na <sub>2</sub> O	K <sub>2</sub> O	Fe <sub>2</sub> O <sub>3</sub>	MnO	TiO <sub>2</sub>	P <sub>2</sub> O <sub>5</sub>
Sr	0.0	9.5	12.2	0.0	4.7	2.2	14.4	2.5	1.2	0.8	0.0	0.9	0.8	0.1	0.0	0.1	0.2	1.6	11.9	3.3	23.0	3.5	21.1	0.0	1.0	1.0	0.1	13.7	8.8	8.6	
Y	0.0	49.7	35.0	44.4	3.7	94.3	71.2	90.6	93.5	93.6	94.9	92.6	93.1	98.4	96.9	96.9	97.0	87.8	35.3	68.9	0.6	46.7	36.5	4.2	0.1	59.2	52.2	13.2	50.3	30.2	
Zr	9.5	49.7	77.5	34.3	5.0	60.9	85.3	71.4	68.1	67.0	51.1	66.5	67.4	57.3	58.8	60.5	53.0	47.4	16.2	82.1	1.2	58.1	13.6	6.2	0.3	76.2	63.4	21.4	94.3	14.8	
Nb	12.2	35.0	77.5	26.2	12.9	42.7	63.5	46.3	43.6	45.2	42.4	43.7	34.9	37.4	38.5	37.1	31.0	38.4	3.1	64.5	0.3	55.6	0.1	13.9	0.3	52.4	24.8	7.5	64.6	26.1	
Cs	0.0	44.4	34.3	26.2	27.3	35.6	39.5	36.4	35.8	39.1	32.2	35.9	44.6	42.8	41.4	41.7	32.5	39.5	11.6	44.3	0.6	62.8	14.5	8.5	27.9	21.9	39.4	14.8	43.3	53.2	
Ba	4.7	3.7	5.0	12.9	27.3	6.7	4.7	6.7	5.0	3.7	5.7	8.0	7.7	5.2	8.1	4.8	4.8	3.5	2.5	1.9	4.4	0.5	9.9	15.7	43.6	1.4	4.1	0.9	2.2	32.4	
La	2.2	94.3	60.9	42.7	35.6	6.7	82.5	98.0	98.5	96.1	94.0	96.8	96.7	96.3	96.3	95.7	96.2	82.1	27.0	75.3	0.2	44.7	32.3	6.5	0.1	61.7	46.6	17.4	59.2	23.4	
Ce	14.4	71.2	85.3	63.5	39.5	4.7	82.5	88.9	85.9	84.6	72.6	84.6	84.4	78.3	74.7	76.4	75.7	71.0	16.0	85.6	0.2	51.8	19.0	4.4	0.0	77.0	61.7	14.7	87.1	14.5	
Pr	2.5	90.6	71.4	46.3	36.4	6.1	98.0	88.9	99.2	99.2	97.7	91.8	98.0	97.5	94.3	94.4	93.0	78.5	25.7	84.4	0.1	52.7	34.0	7.3	0.2	72.2	52.3	15.7	70.1	25.3	
Nd	1.2	93.5	68.1	43.6	35.8	5.9	98.5	85.9	99.2	97.9	94.3	98.3	97.7	96.6	95.6	95.3	95.2	82.5	28.4	80.8	0.0	40.2	37.4	6.5	0.2	69.3	53.5	14.3	66.2	26.1	
Sm	0.8	93.6	67.0	45.2	39.1	3.7	96.1	84.6	97.7	97.9	92.7	96.9	95.9	96.0	93.7	94.6	92.7	84.1	25.3	85.9	0.0	61.6	34.8	7.2	0.2	74.3	52.7	12.8			

Element	Uncertainty		Coefficient of regression R		Slope s (%/m)		Total concentration change $\Delta t$ (%)	
	Average	Maximum	a	b	a	b	a	b
Al <sub>2</sub> O <sub>3</sub>	-	22	0.59	-137.08	-13.64	0.51	-113.72	
Th	10	19	0.51	-66.54	-6.62	0.63	-91.29	
Nb	6	8	0.56	-55.16	-5.49	0.59	-63.91	
Zr	3	6	0.67	-53.25	-5.30	0.75	-65.51	
TiO <sub>2</sub>	-	10	0.55	-40.82	-4.06	0.68	-54.80	
La	3	6	0.41	-49.35	-4.91	0.70	-87.33	
Ce	3	6	0.57	-61.75	-6.14	0.72	-84.59	
Pr	3	11	0.45	-58.79	-5.85	0.70	-97.68	
Nd	2	6	0.37	-48.90	-4.87	0.68	-91.88	
Sm	7	14	0.38	-50.81	-5.06	0.62	-88.03	
Eu	9	14	0.30	-38.46	-3.83	0.59	-76.98	
Gd	4	12	0.39	-44.16	-4.39	0.64	-76.90	
Tb	6	11	0.40	-51.00	-5.08	0.69	-90.79	
Dy	5	14	0.34	-35.98	-3.58	0.63	-68.57	
Ho	5	12	0.36	-36.98	-3.68	0.66	-69.07	
Er	5	11	0.43	-40.85	-4.07	0.67	-67.04	
Yb	8	12	0.37	-35.61	-3.54	0.63	-62.45	
Lu	7	16	0.35	-29.56	-2.94	0.56	-50.31	
Y	2	6	0.30	-26.25	-2.61	0.59	-52.53	
Ba	3	4	0.46	-39.03	-3.88	0.46	-39.03	
Sr	2	3	0.42	-9.33	-0.93	0.49	-11.88	
CaO	-	4.5	0.19	0.80	0.08	0.10	0.43	
MgO	-	4	0.10	2.35	0.23	0.06	-1.33	
Cs	5	10	0.29	-33.62	-3.35	0.46	-57.74	
K <sub>2</sub> O	-	5	0.42	-29.78	-2.96	0.51	-39.41	
Li	5	11	0.09	6.18	0.62	0.40	0.56	
Na <sub>2</sub> O	-	9	0.43	26.21	2.61	0.33	21.50	
Fe <sub>2</sub> O <sub>3</sub>	-	3	0.46	-25.77	-2.56	0.53	-32.69	
P <sub>2</sub> O <sub>3</sub>	-	7	0.04	-3.05	-0.30	0.17	-14.87	
Pb	3	6	0.02	-0.61	-0.06	0.23	-8.18	
U	3	7	0.29	21.57	2.15	0.43	33.56	
MnO	-	7	0.49	42.61	4.24	0.70	62.80	

Regression parameters for the concentrations versus scarp height plots.

Mean and maximum uncertainties on concentrations are reported. R is coefficient of regression, s the slope of the regression in %/m, and t the total concentration change over the whole scarp height H, in % (t = s x H). Negative (positive) slopes indicate a concentration decrease (increase) as one goes upscarp.

a) Regressions are calculated from all measurements listed in Table 1.a) Regressions are calculated from all measurements listed in Table 1.b) Regressions are calculated from all measurements listed in Table 1, but those for the C15 sample.

	Na <sub>2</sub> O	MgO	Al <sub>2</sub> O <sub>3</sub>	SiO <sub>2</sub>	P <sub>2</sub> O <sub>5</sub>	SO <sub>3</sub>	K <sub>2</sub> O	TiO <sub>2</sub>	MnO	Fe <sub>2</sub> O <sub>3</sub>
Na <sub>2</sub> O	35.9	49.7	58.3	5.1	20.4	28.9	40.5	40.7	44.8	
MgO	0.4	0.5	16.4	1.0	0.2	1.0	22.7	2.2		
Al <sub>2</sub> O <sub>3</sub>	88.1	4.1	67.6	60.6	48.4	29.0	54.1			
SiO <sub>2</sub>	88.1	0.1	54.0	65.6	74.3	39.3	79.7			
P <sub>2</sub> O <sub>5</sub>	4.1	0.1	16.3	0.6	7.4	1.2	10.2			
SO <sub>3</sub>	67.6	54.0	16.3	76.1	22.7	32.0	27.7			
K <sub>2</sub> O	0.2	60.6	0.6	76.1	46.4	47.5	51.9			
TiO <sub>2</sub>	40.5	48.4	74.3	7.4	22.7	29.2	89.5			
MnO	40.7	29.0	39.3	1.2	32.0	29.2	32.7			
Fe <sub>2</sub> O <sub>3</sub>	44.8	2.2	54.1	10.2	27.7	51.9	89.5	32.7		

Correlation coefficients between major elements in the coating (from Table A). The coefficients were calculated using simple linear regressions. When the coefficient is larger than 70, the field is shown in grey, and the intensity of the grey is correlated to the R value.

	Wave 5		Wave 4		Wave 3		Wave 2	
	Base (-0.35 m) T <sub>r</sub> = 4.85 ka	Top (1.25 m) T <sub>r</sub> = 1.85 ka	Base (1.95 m) T <sub>r</sub> = 1.85 ka	Top (3.85 m) T <sub>r</sub> = 0.70 ka	Base (4.45 m) T <sub>r</sub> = 0.70 ka	Top (6.65 m) T <sub>r</sub> = 3.10 ka	Base (7.15 m) T <sub>r</sub> = 3.10 ka	Top (9.2 m) T <sub>r</sub> = 1.50 ka
	Concentration peak (%)	Concentration peak (%)	Concentration peak (%)	Concentration peak (%)	Concentration peak (%)	Concentration peak (%)	Concentration peak (%)	Concentration peak (%)
Sr	-	-	-	-	-	-	-	-
Y	46.1	37.8	64.1	45.5	27.2	38.4	52.1	46.3
Zr	69.3	11.1	19.1	27.5	9.3	4.5	30.4	10.9
Nb	-	-	24.5	56.5	34.9	11.2	73.2	46.8
Cs	-	46.3	46.8	52.7	-	-	85.3	42.8
Ba	-	-	-	-	-	-	-	-
La	60.0	44.4	88.8	57.7	41.6	47.2	45.2	55.6
Ce	85.0	28.5	52.9	39.9	30.2	31.2	45.1	53.9
Pr	85.3	55.6	90.6	60.3	39.2	46.2	49.5	59.5
Nd	80.0	50.3	86.4	66.8	41.0	51.0	50.4	61.8
Sm	91.1	67.7	100.3	79.6	43.8	48.1	60.7	58.7
Eu	54.0	48.3	81.6	58.4	51.3	61.7	69.3	84.0
Gd	70.8	41.4	85.5	58.4	28.6	33.9	54.4	67.2
Tb	79.2	45.1	90.8	57.6	19.7	40.7	48.6	52.3
Dy	61.1	47.3	68.2	46.9	31.8	40.2	58.3	59.2
Ho	49.1	39.9	68.8	38.8	26.4	35.1	55.8	56.0
Er	56.7	41.5	66.7	39.4	28.7	41.6	51.8	35.2
Yb	52.4	35.5	67.1	36.3	34.1	40.9	48.7	55.5
Lu	42.5	34.1	43.3	51.3	42.5	38.5	53.3	52.1
Pb	-	-	16.9	2.9	-	-	29.8	21.6
Th	107.3	84.1	65.1	66.1	15.6	28.1	67.3	41.4
U	5.5	5.9	-	-	-12.5	13.1	61.7	90.8
Li	30.0	15.4	-	-	49.7	59.1	4.7	25.9
Al <sub>2</sub> O <sub>3</sub>	-	-	-	-	-	-	-	-
CaO	-	-	-	-	-	-	-	-
MgO	5.8	9.0	5.8	2.6	13.0	4.2	-	-
Na <sub>2</sub> O	17.7	16.1	-	-	44.5	41.9	-	-
K <sub>2</sub> O	67.8	36.0	20.9	17.9	24.4	11.1	50.0	41.0
Fe <sub>2</sub> O <sub>3</sub>	-	-	4.6	7.2	-	-	54.2	35.1
MnO	-	-	-	-	8.0	14.0	-3.9	62.6
TiO <sub>2</sub>	76.4	10.3	20.1	12.7	4.7	10.4	30.0	18.5
P <sub>2</sub> O <sub>5</sub>	34.1	41.9	51.3	64.8	-	-	51.3	48.1

Amplitudes of the concentration peaks at transition zones between earthquake slip zones. Data are presented and 'waves' defined in Figure 6b. 'Base' and 'top' are for the two peak-concentrations that bound each 'wave'. The residence times (Tr) over which the peak-concentrations have built up are indicated. Elements having concentrations higher than 46% (average overall value) are highlighted in grey.

Elements	R coefficient for model of Fig. 9 (%)	R coefficient for linear regression (%)
Y	90.61	58.51
La	89.55	69.91
Ce	85.58	71.60
Pr	92.25	69.99
Nd	92.38	67.68
Sm	94.04	62.31
Eu	92.83	58.79
Gd	92.74	63.91
Tb	90.57	69.07
Dy	91.99	63.10
Ho	91.88	66.07
Er	92.11	67.15
Yb	89.14	62.68
Lu	84.28	55.96
<b>R<sub>min</sub></b>	<b>84.28</b>	<b>55.96</b>
<b>R<sub>average</sub></b>	<b>90.71</b>	<b>64.77</b>
<b>R<sub>max</sub></b>	<b>94.04</b>	<b>71.60</b>

Best-fit coefficients (R) obtained (for REE-Y) using the model shown in figure 9 (left column) compared to values calculated with a simple linear regression (right column).

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